

Dynamic heterogeneity of glassy ionics: Results from nuclear magnetic resonance and low-frequency spectral hole burning

Roland Böhmer

Video Module 1: Introduction

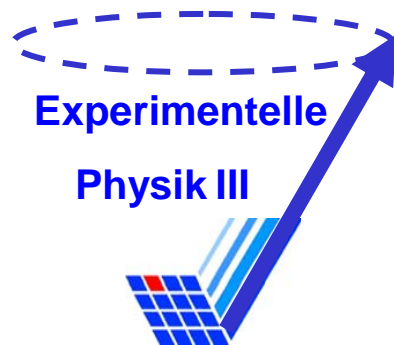
1. Introduction
2. Ion dynamics studied by NMR
3. Nonresonant spectral hole burning in CKN
4. Conclusions

Glass Lecture Series: prepared for and produced by the
International Material Institute for New Functionality in Glass
An NSF sponsored program – material herein not for sale
Available at www.lehigh.edu/imi

Delivered at Lehigh University Dec. 7, 2006



Univ. Dortmund



S. Berndt

R. Küchler

F. Qi

C. Rier

DFG

Ag K. R. Jeffrey
Guelph / Canada

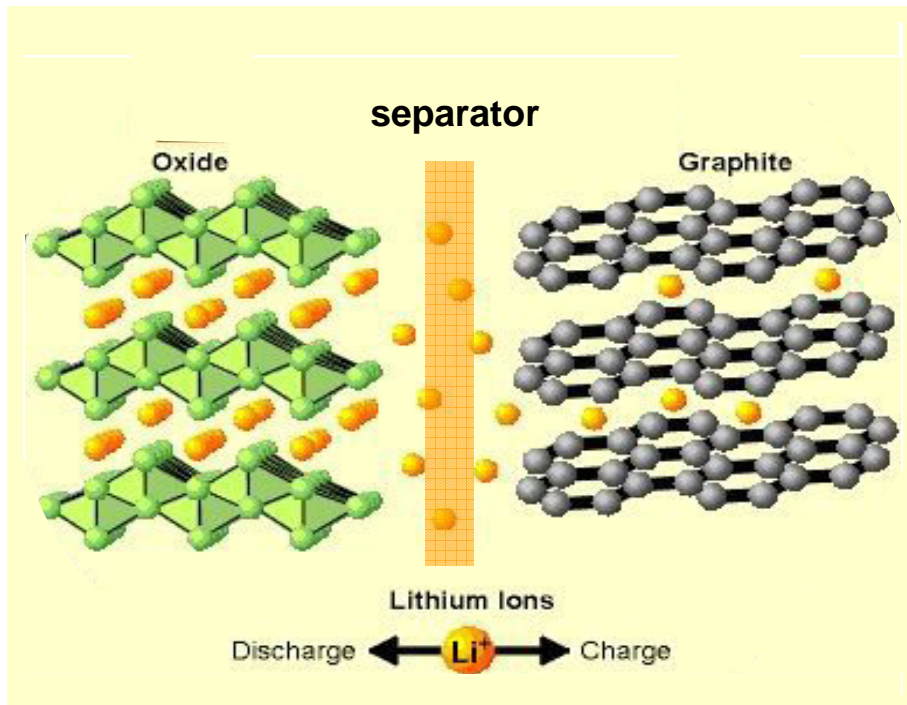
CKN R. Richert
Arizona State Univ.



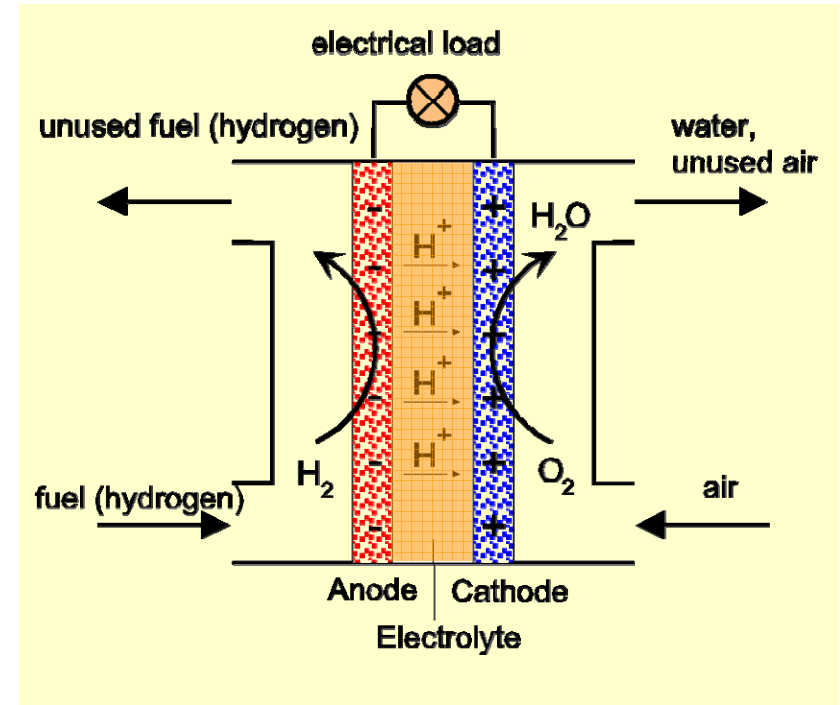
**Graduierten-
kolleg 298**

Transport in solid electrolytes

Li ion battery

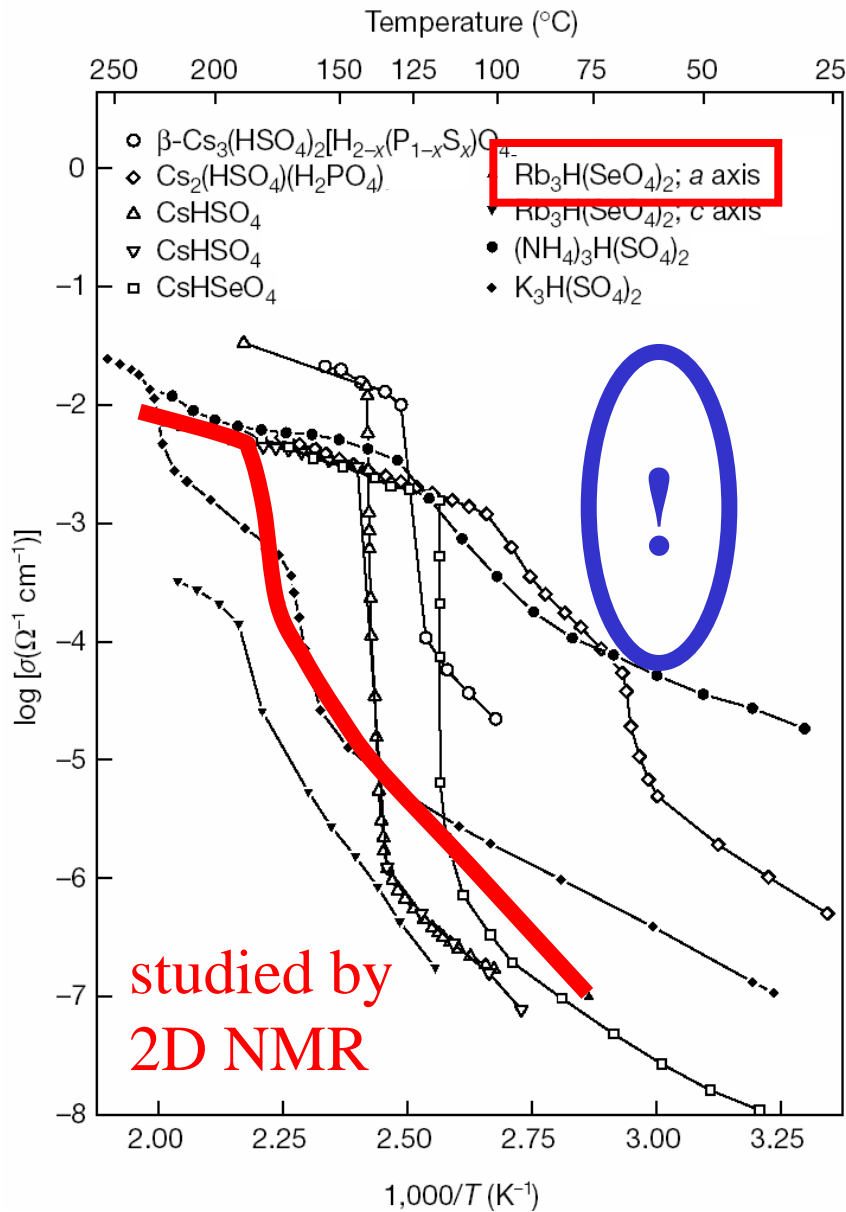


"fuel" cell

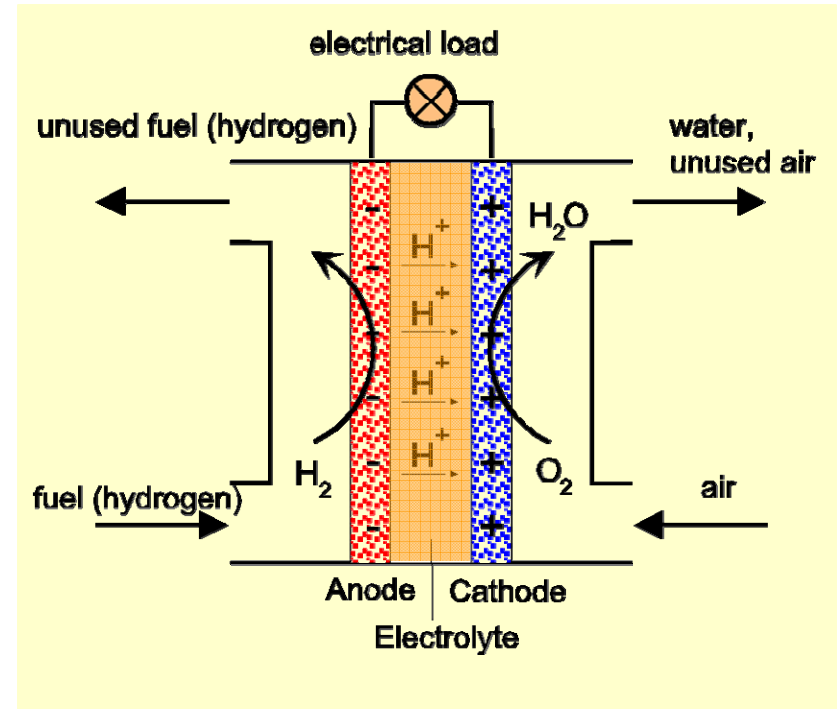


Membranes with high ionic conductivities required

electrical conductivity of solid electrolytes



Haile et al., Nature 410, 910 (2001)



"superionic"
anhydrous fuel cell

Insight into transport mechanisms?

applications

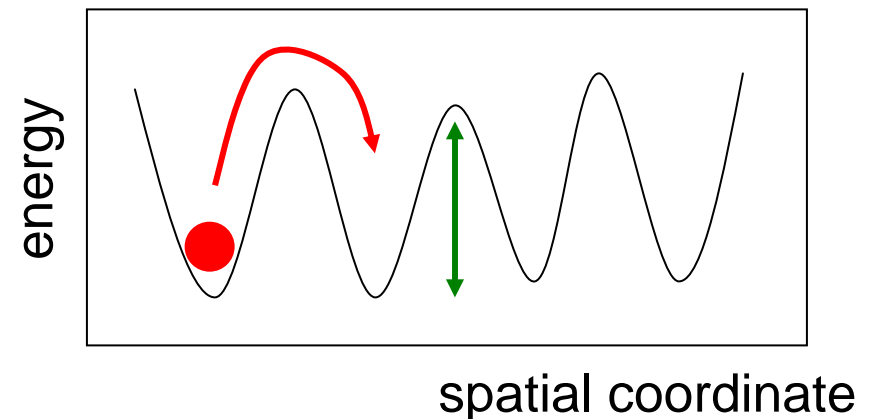
sensors, fuel cells,
rechargeable batteries, ...

requirements

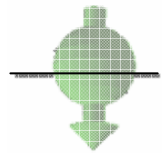
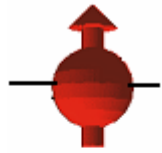
stable, light, solid, inexpensive,
high ionic & low electronic conductivity,
suitable operating temperature, ...

goal of group at Dortmund U

new experimental methods
for better understanding
of transport mechanisms
in solid ion conductors



Principle of magnetic resonance

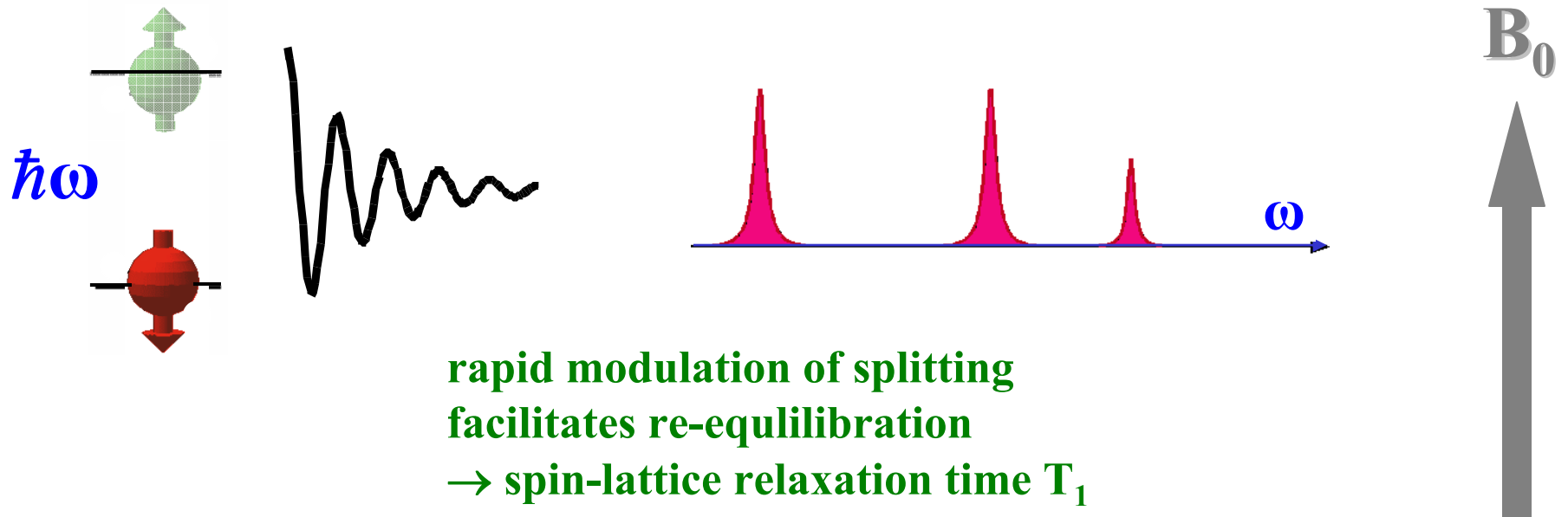


B_0



element specific
quantitative
locally selective
non-destructive
experimentally versatile

Principle of magnetic resonance



frequency perturbation $\Delta\omega$ encoded by

spatial coordinate

electronic environment

distances and angles

orientation

imaging

chemical analysis

structural elucidation

fiber texture testing

element specific

quantitative

locally selective

non-destructive

experimentally versatile

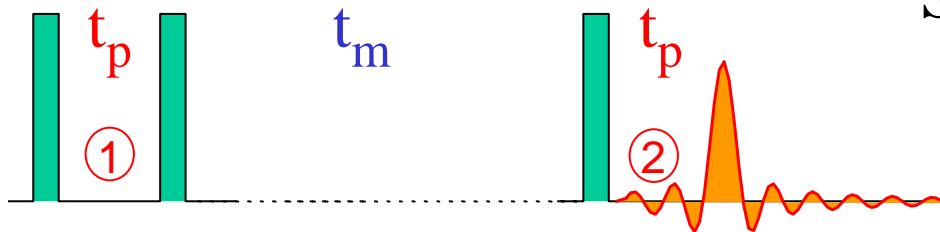
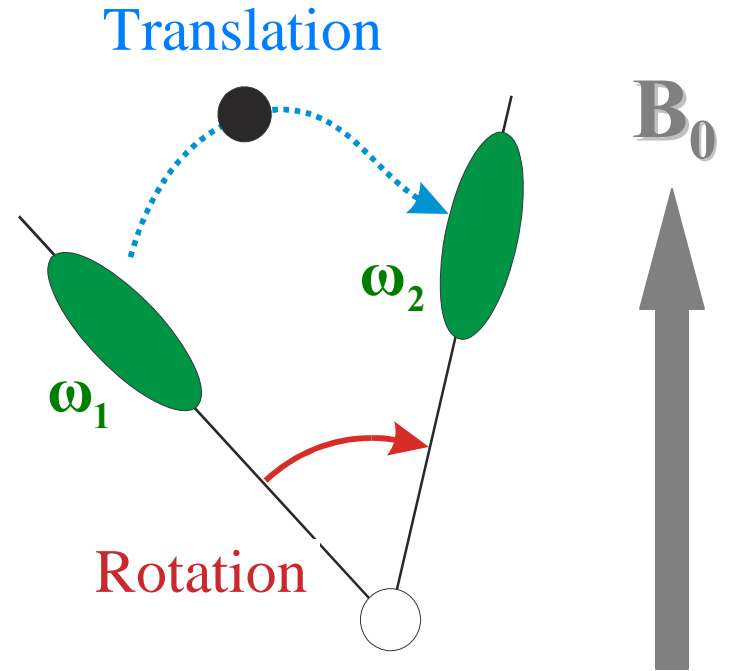
Motional processes

detection of NMR frequencies
and their time evolution



NMR frequency
encoded by
spatial coordinate
chemical environment
orientation
...

diffusion, flow
exchange, **translation**
reorientational motion



t_p evolution time
 t_m mixing time

$$S_2(t_p, t_m) = \langle \exp[it_p \omega(0)] \exp[-it_p \omega(t_m)] \rangle$$

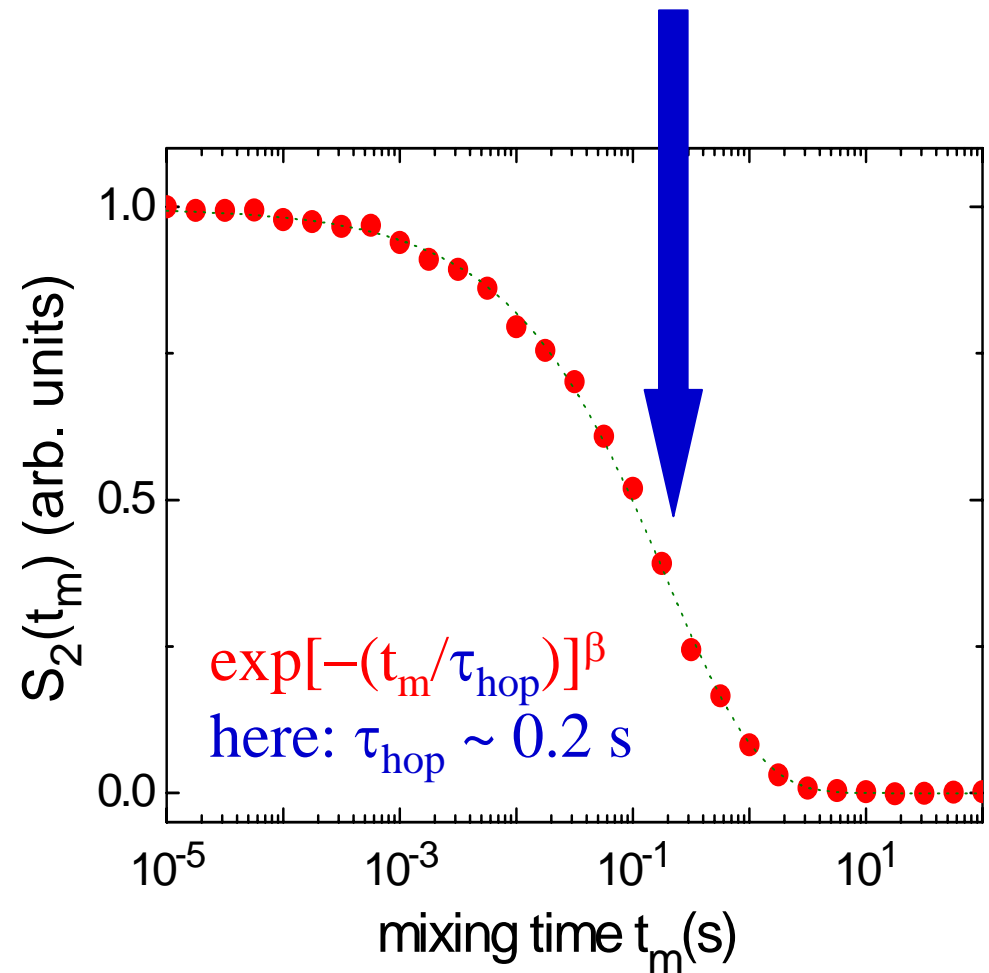
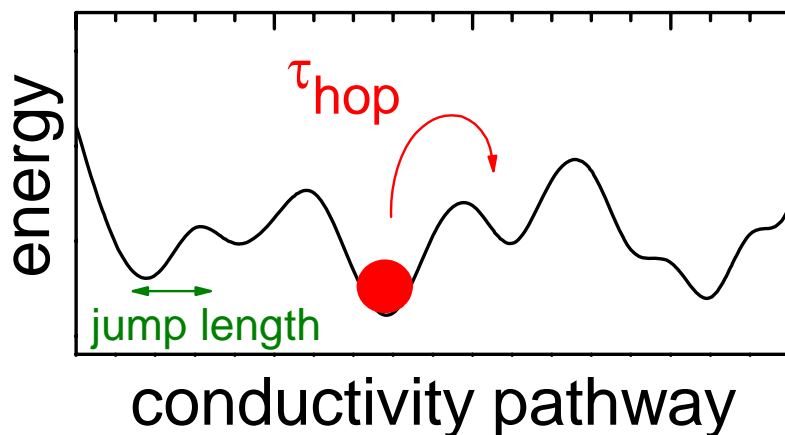
$$S(\vec{q}, t) \propto \langle \exp[i\vec{q}\vec{r}(0)] \exp[-i\vec{q}\vec{r}(t)] \rangle$$

correlation of ω at
2, 3, 4 ... points in time

Echo height measures the fraction of ions which did NOT hop during t_m

Direct determination of correlation time on which an ion hops

Scattering „vector“ q , i.e. spatial sensitivity is determined by the inverse mean jump length



This lecture continues on a 2nd module -

Video Module 2 : Ion Dynamics by NMR (Part 2)

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1. Introduction

Ion conductors

Nuclear magnetic resonance

2. Ion dynamics studied by NMR

Li hopping in aluminosilicates

Heterogeneity in silver borate glasses

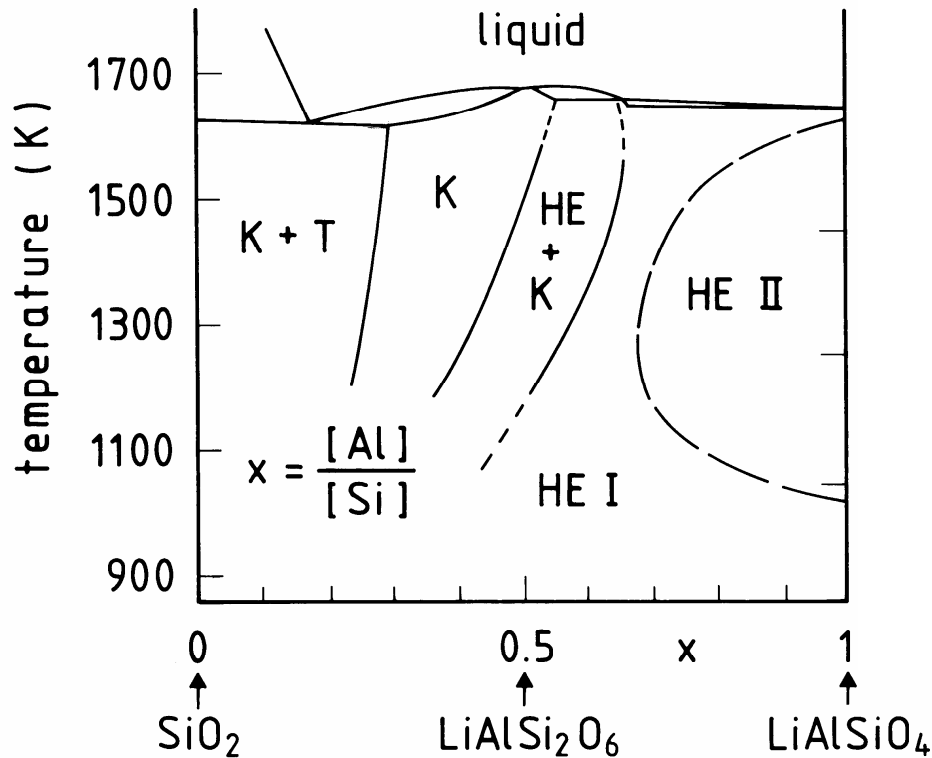
3. Nonresonant spectral hole burning

Material and method

Experiments on Ca-K-NO₃ glass

4. Conclusions

Lithium aluminosilicates



Mesocrystalline texture

(Schott, Mainz)

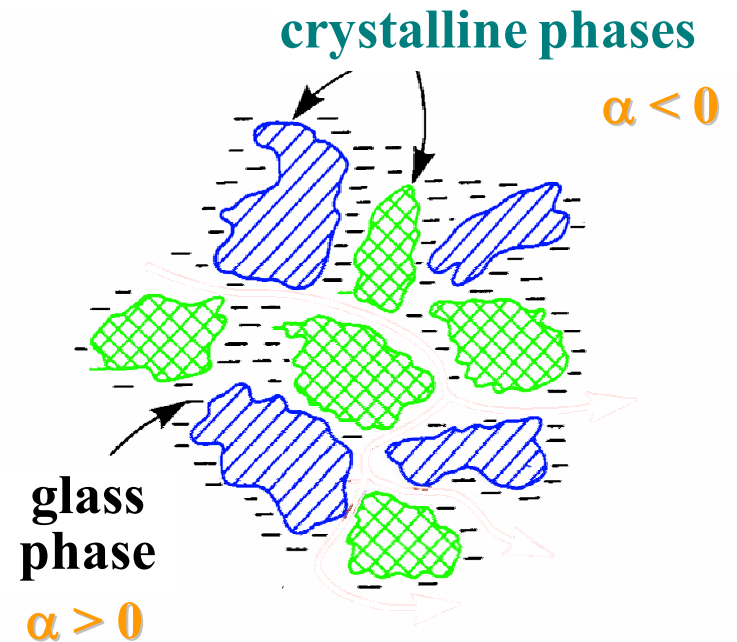
Zerodur[®] $x = 0.52$ 70% 62 nm

Zerodur M[®] $x = 0.44$ 50% 45 nm

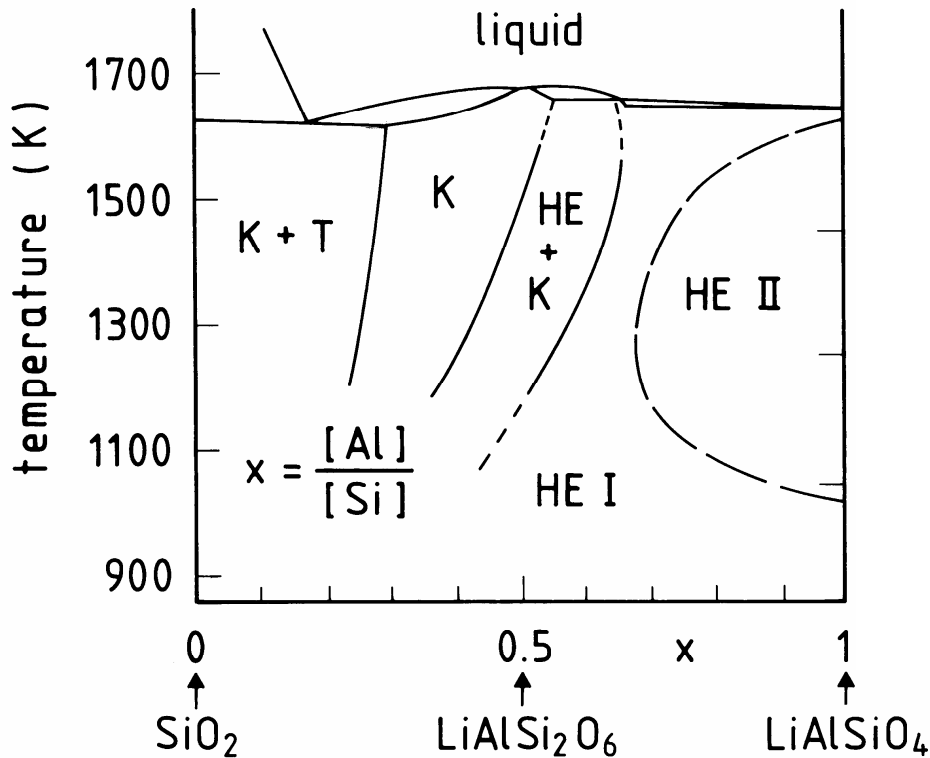
crystalline model systems

β -spodumene **LiAlSi₂O₆**
 β -eucryptite **LiAlSiO₄**

$x = [Al]/[Si]$
0.5
1



Lithium aluminosilicates

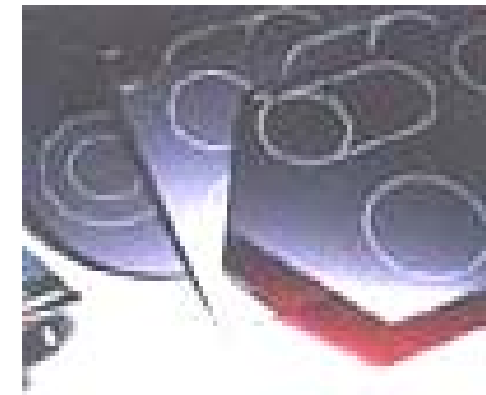


Mesocrystalline texture

(Schott, Mainz)

Zerodur[®]	x = 0.52	70%	62 nm
Zerodur M[®]	x = 0.44	50%	45 nm

precision optics,
ceramic cooking ware



crystalline model systems

β-spodumene



0.5

β-eucryptite

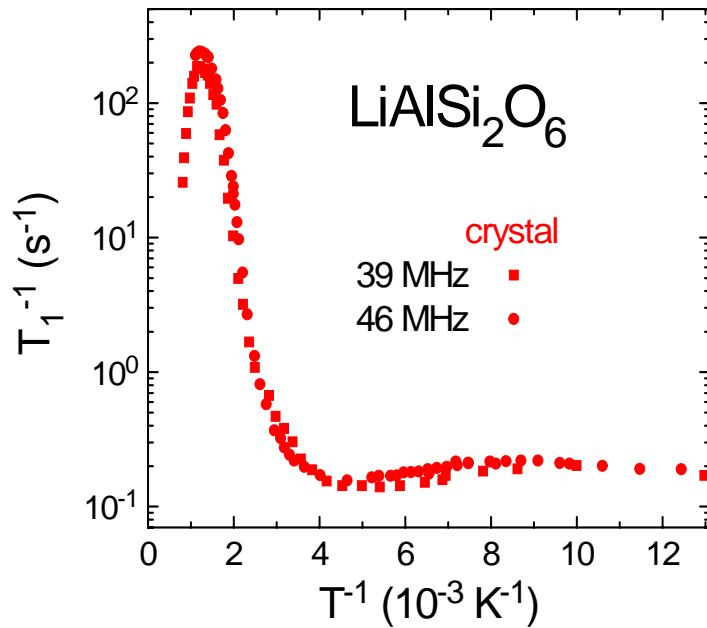


1

x = [Al]/[Si]

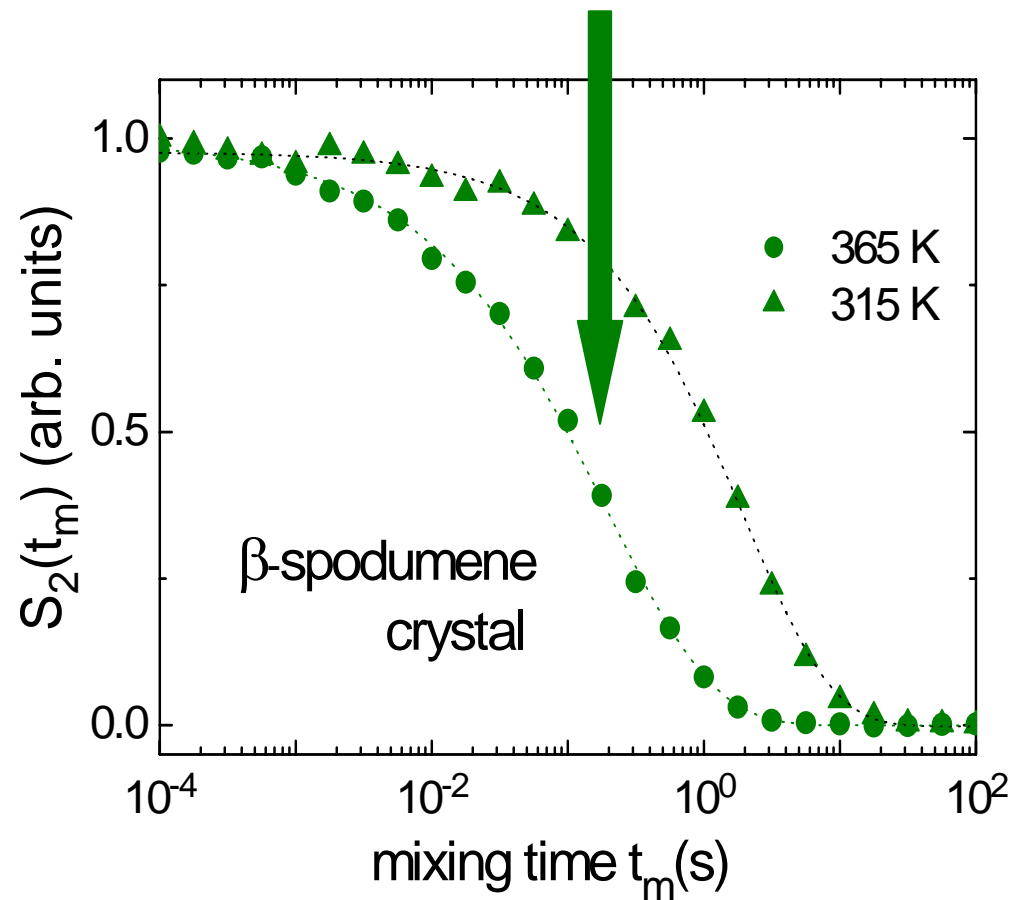
^7Li ion conductor β -spodumene

**Spin-lattice relaxation
indirect information
on fast motion**

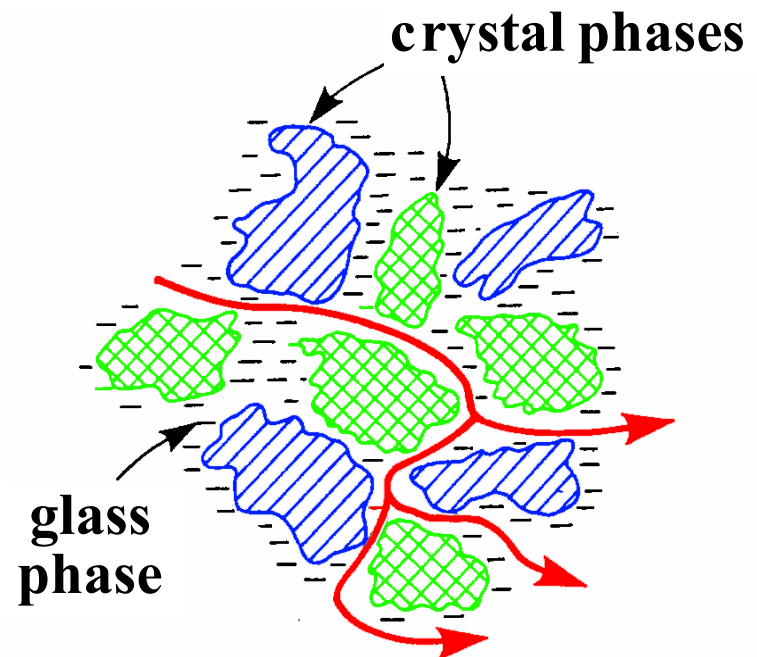
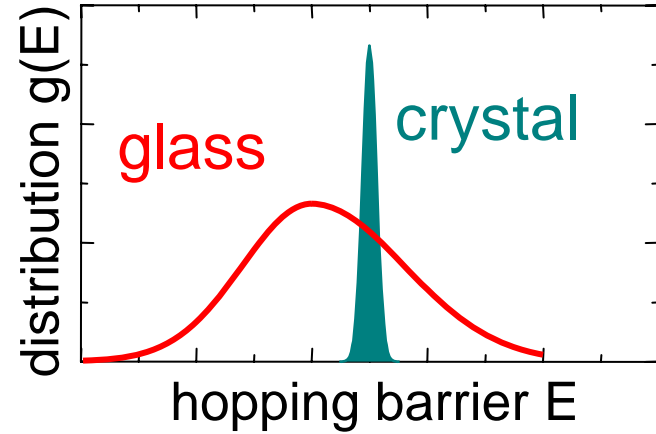
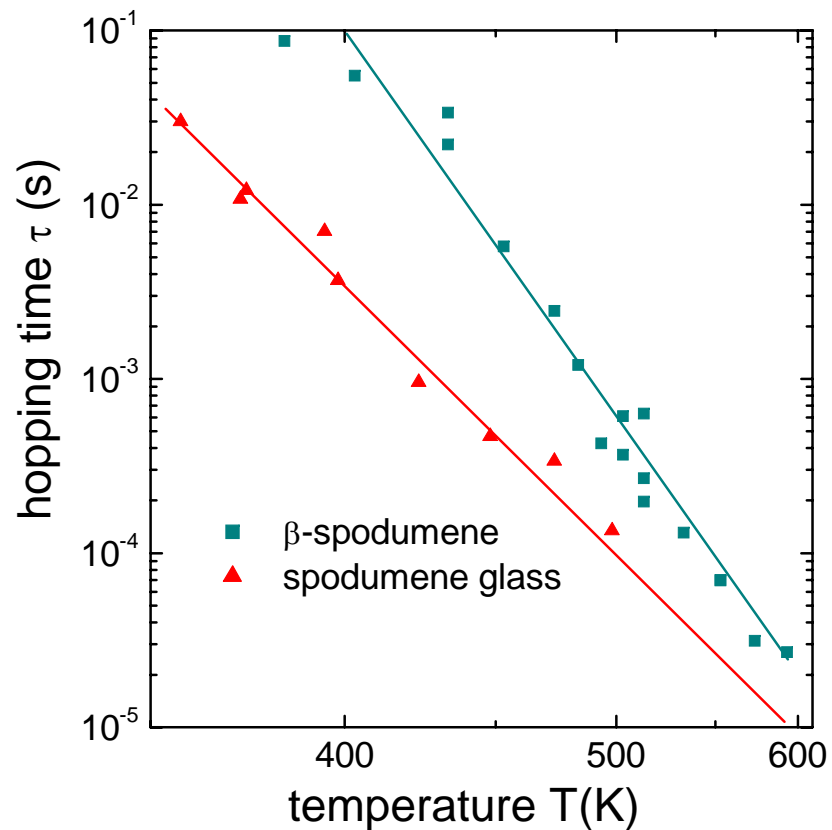


**rate maxima
for $\omega\tau \sim 0.61$**

**T-dependent time scale
and non-exponentiality
of ion hopping**



Motion of Li^+ in crystalline and in glassy environments



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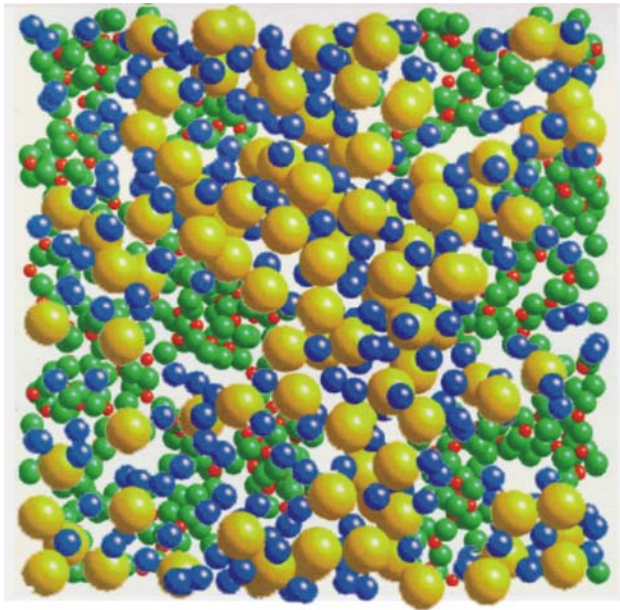
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Material and method

Experiments on Ca-K-NO₃ glass

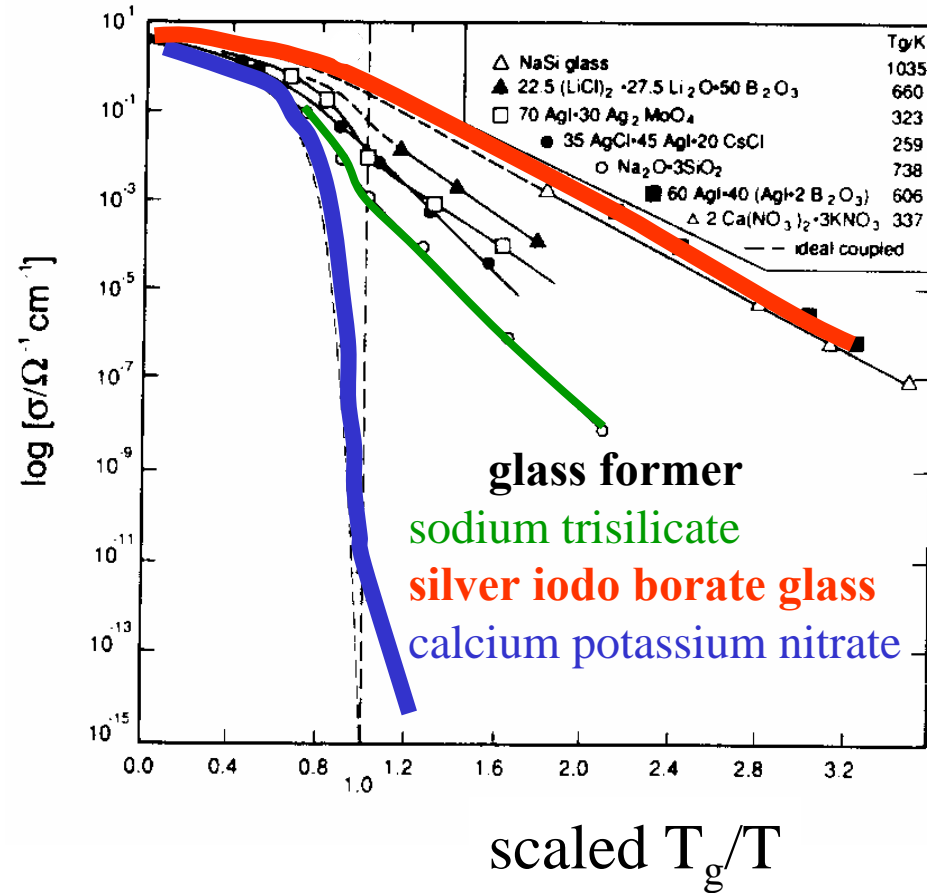
4. Conclusions

Conductivity of glassy ionics



Ag: 1.0 I: 1.8 B: 0.5 O: 0.9 Å

Howells et al., JP-CM 11, 9275 (1999)



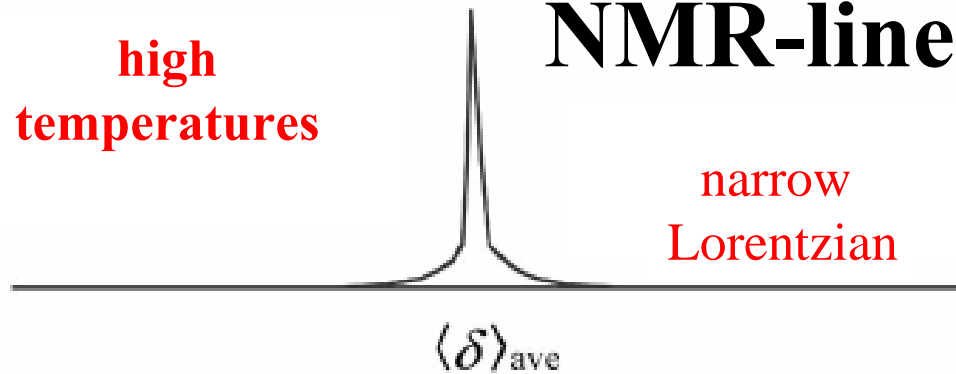
Angell, Chem. Rev.
90, 523 (1990)

T_{sc}
738 K
610 K
333 K

large decoupling and T ~ T_g are prerequisites for battery application

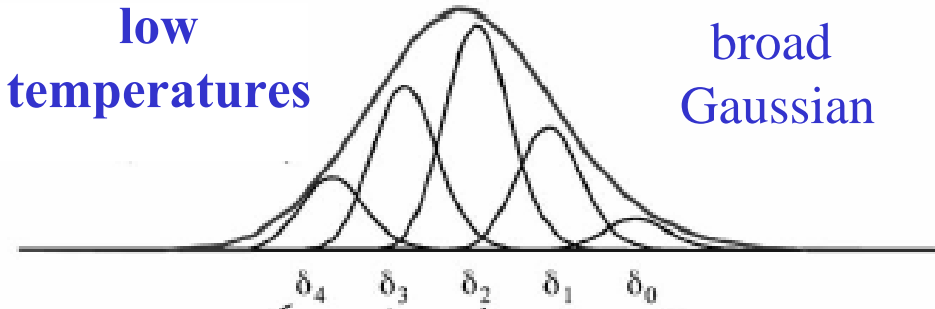
NMR-line shape of AgI-glasses

high temperatures



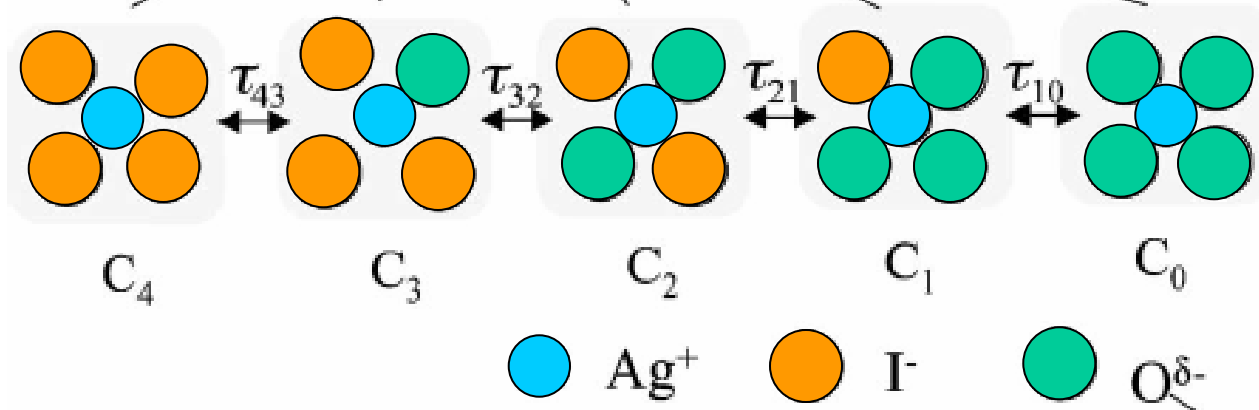
motional narrowing

low temperatures



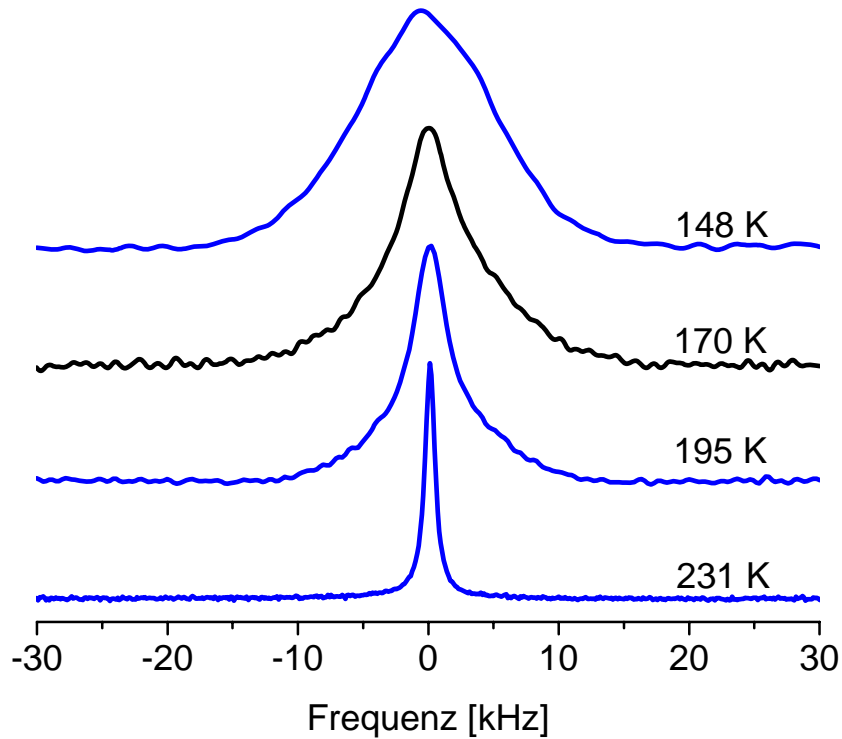
intermediate temperatures ?

?



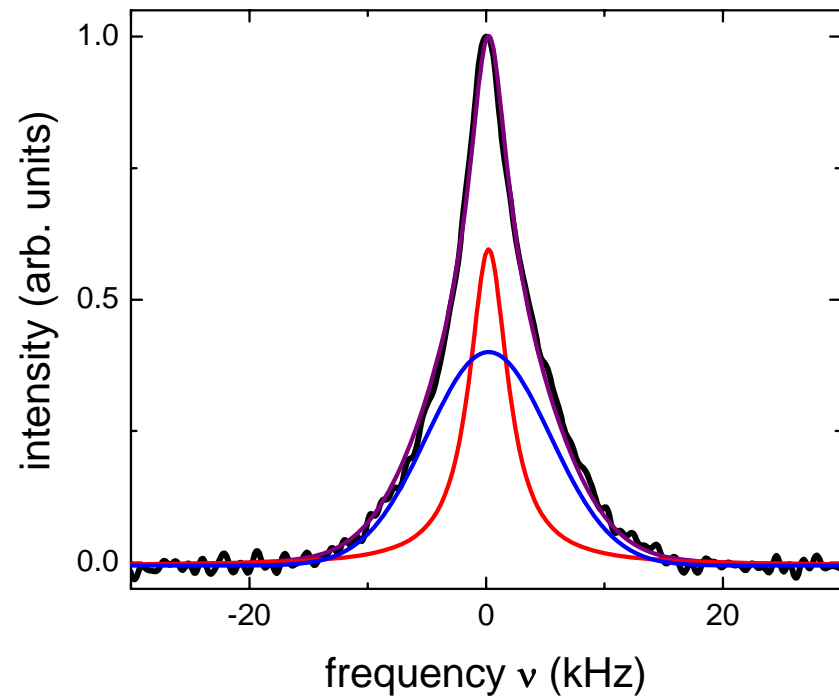
hopping rate τ ?

Line shape of $(^{109}\text{AgI})_x-(\text{Ag}_2\text{O}-\text{B}_2\text{O}_3)_{1-x}$



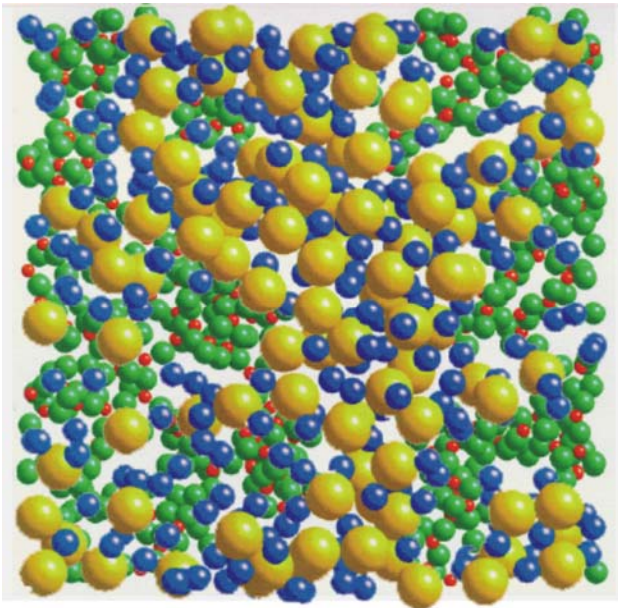
**broadening due to
distribution of isotropic
chemical shifts**

superposition of sub-spectra



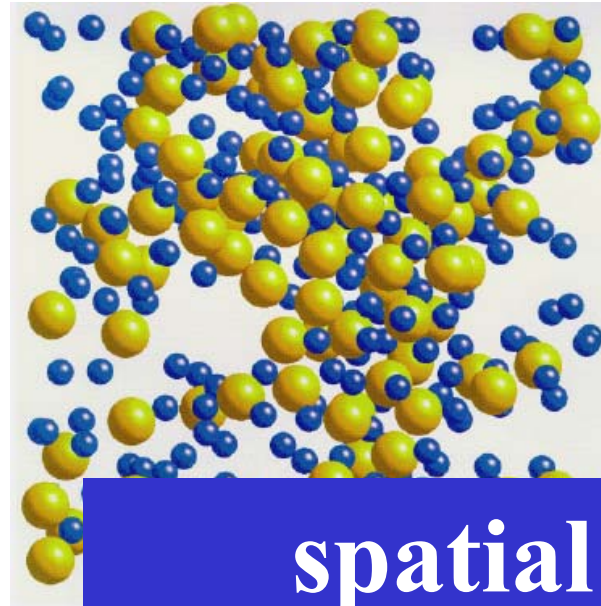
hint: **slow** and **fast**
Ag⁺ ions could coexist

Heterogeneity of ion conductors

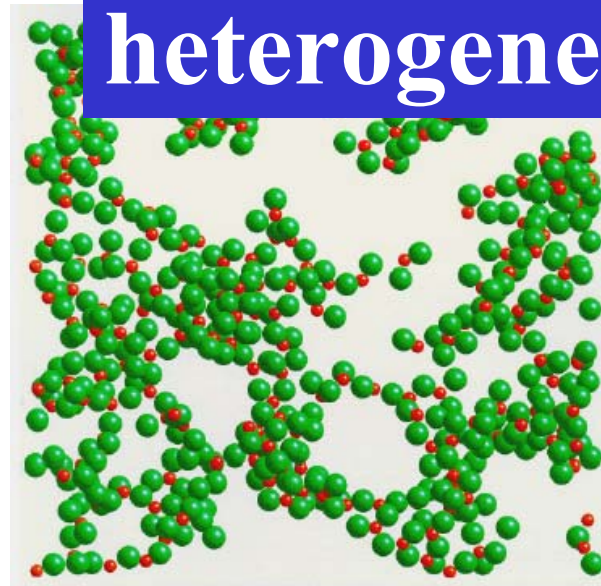


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Howells et al., JP-CM 11, 9275 (1999)



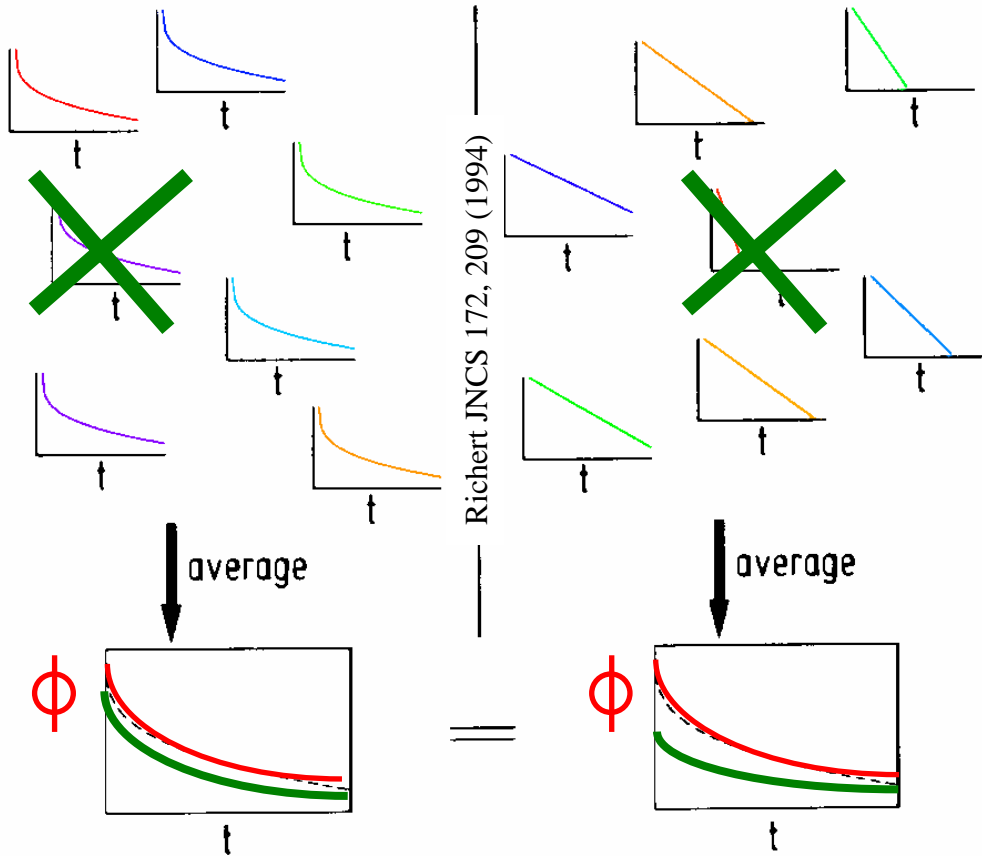
spatial
heterogeneity !



Nature of the non-exponential relaxation ?

homogeneous

heterogeneous



amplitude
modification

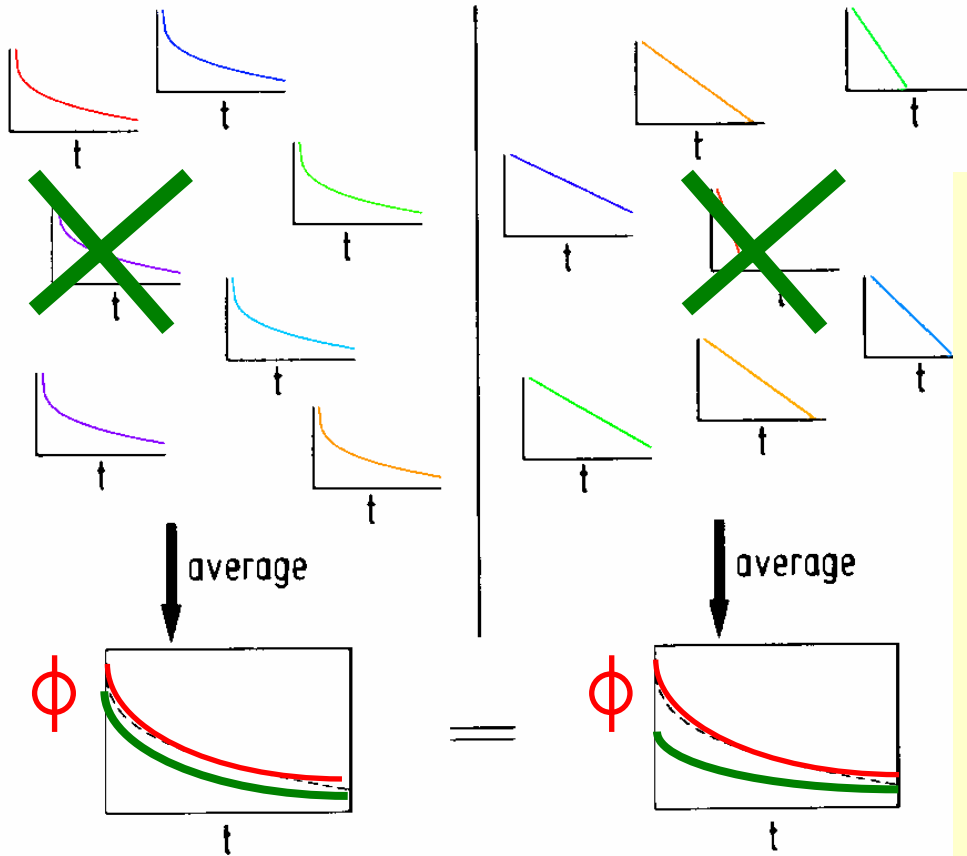
line shape
modification

dynamic
heterogeneity ?

Nature of the non-exponential relaxation

homogeneous

heterogeneous



amplitude
modification

line shape
modification

experimentally distinguishable
via sub-ensemble selection

multi-dimensional NMR
single-molecule detection
nonresonant hole burning
optical probe spectroscopy
computer simulation

resonance frequencies:

^1H : 43 MHz/Tesla

^{109}Ag : 2 MHz/Tesla !!

Interdisziplinäres Zentrum
für magnetische Resonanz
(IZMR)

Universität Dortmund

problem:

Ag: very low
sensitivity

solution:

availability of very
high magnetic fields

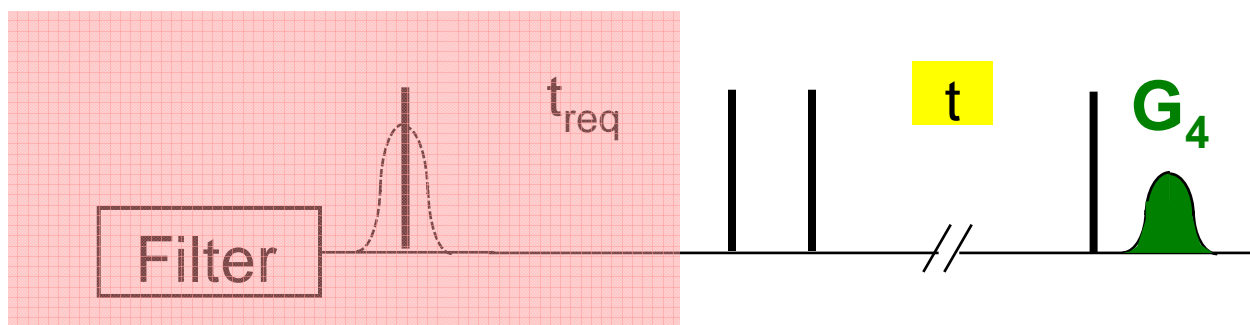


four-time NMR

jump correlation function



measure for the probability that no jump occurred during t_{filter}



jump correlation function of the selected subensemble

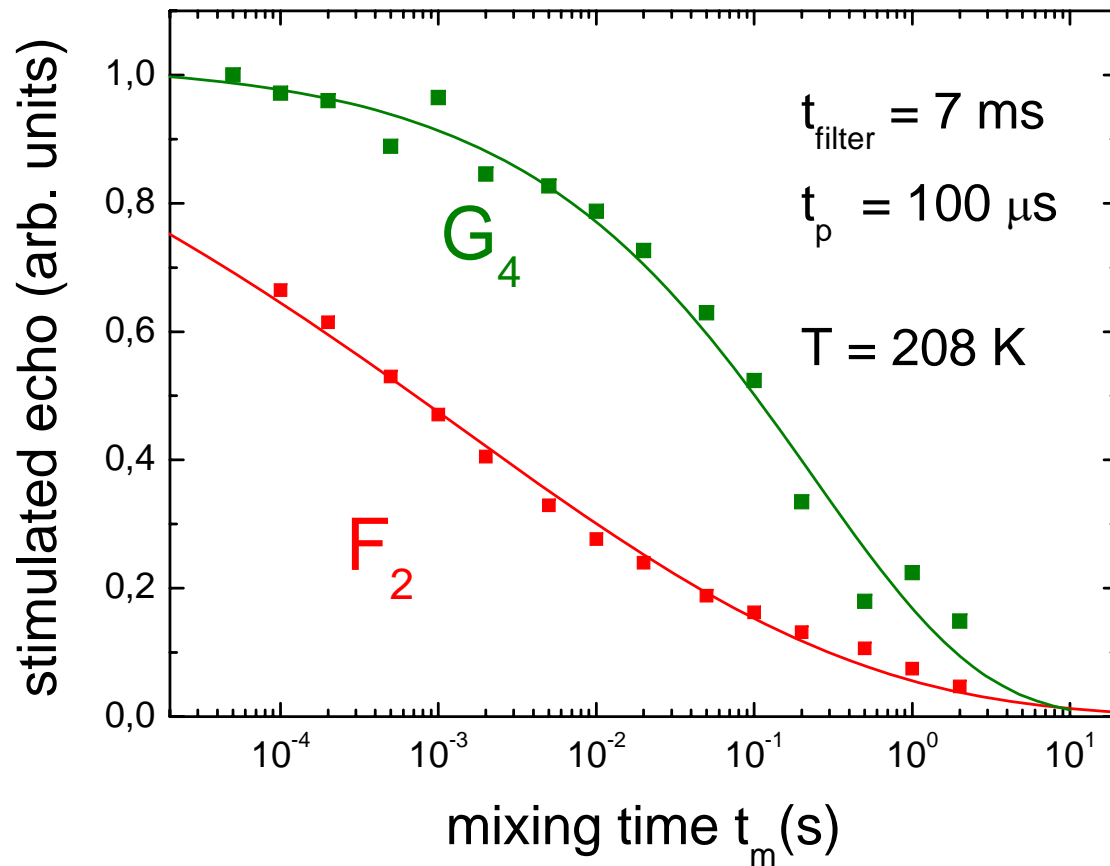
extensively used for polymers and supercooled organic liquids:

Schmidt-Rohr, Spiess, Phys. Rev. Lett. 66, 3020 (1991)

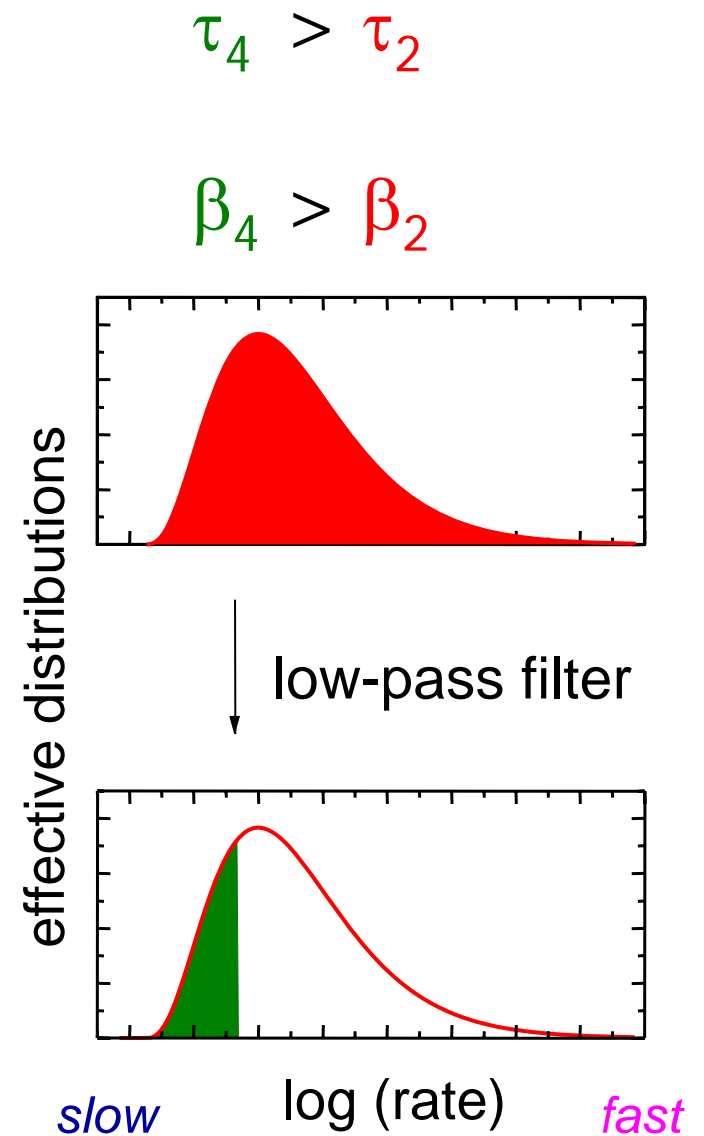
Böhmer, Diezemann, Hinze, Rössler, Prog. NMR Spectrosc. 39, 191 (2001)

Ag-phosphates: Vogel, Brinkmann, Eckert, Heuer, PRB 69, 094302 (2004)

$(\text{AgI})_{0.5}-(\text{Ag}_2\text{O}-\text{B}_2\text{O}_3)_{0.5}$



successful low-pass filtering
confirms
dynamic heterogeneity
in AgI-borate glasses



This lecture continues on a 3rd module -

Video Module 3 : Nonresonant spectral hole burning

Dynamic heterogeneity of glassy ionics: Results from nuclear magnetic resonance and low-frequency spectral hole burning

Roland Böhmer

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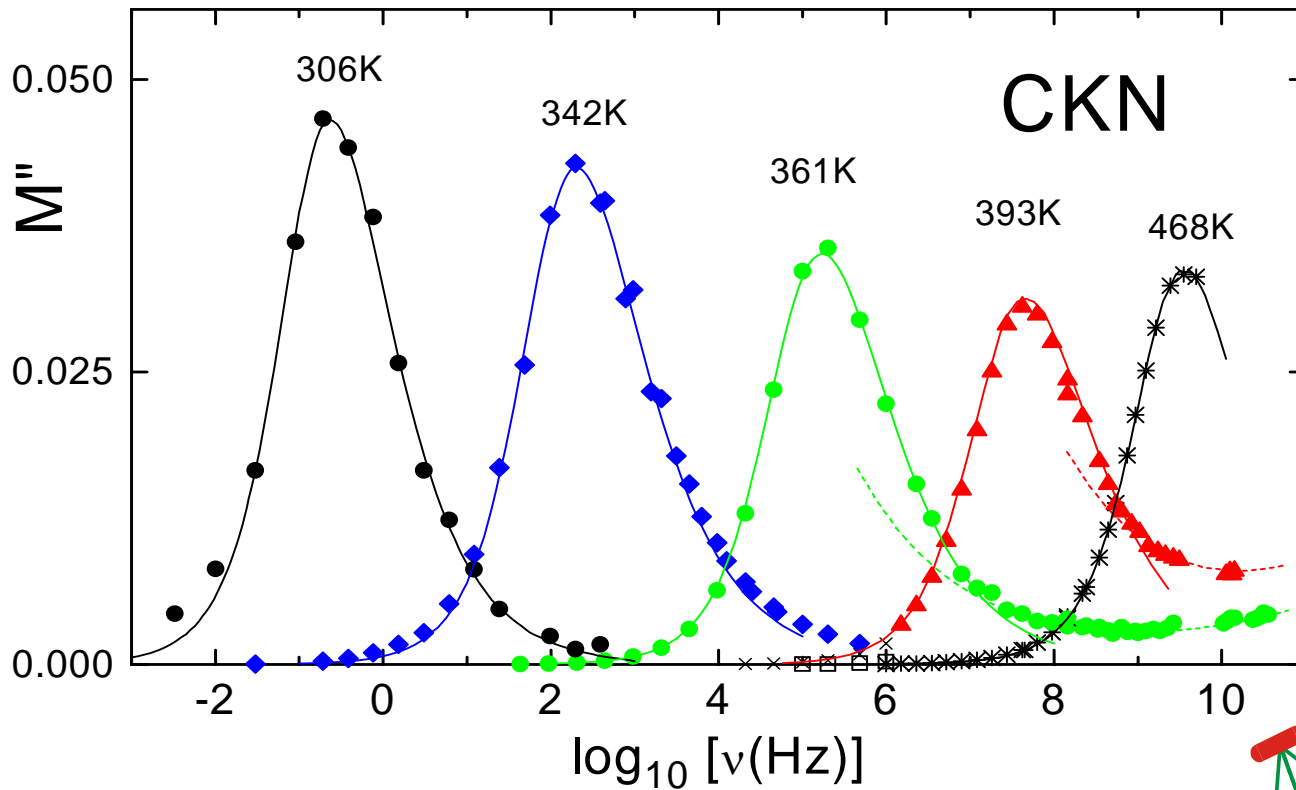
3. Nonresonant spectral hole burning

Material and method

Experiments on Ca-K-NO₃ glass

4. Conclusions

Non-exponential conductivity and structural relaxation

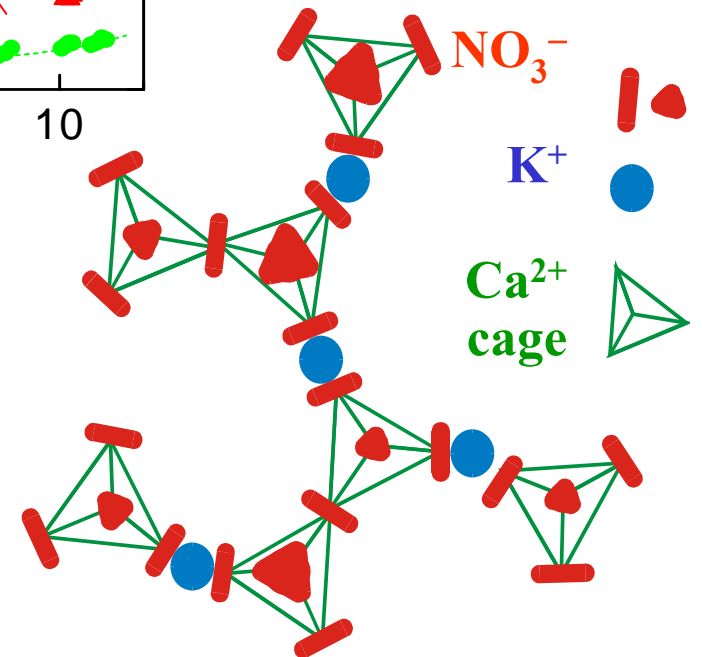


$$M = i\omega\epsilon_0/\sigma$$

$$T_g = 333 \text{ K}$$

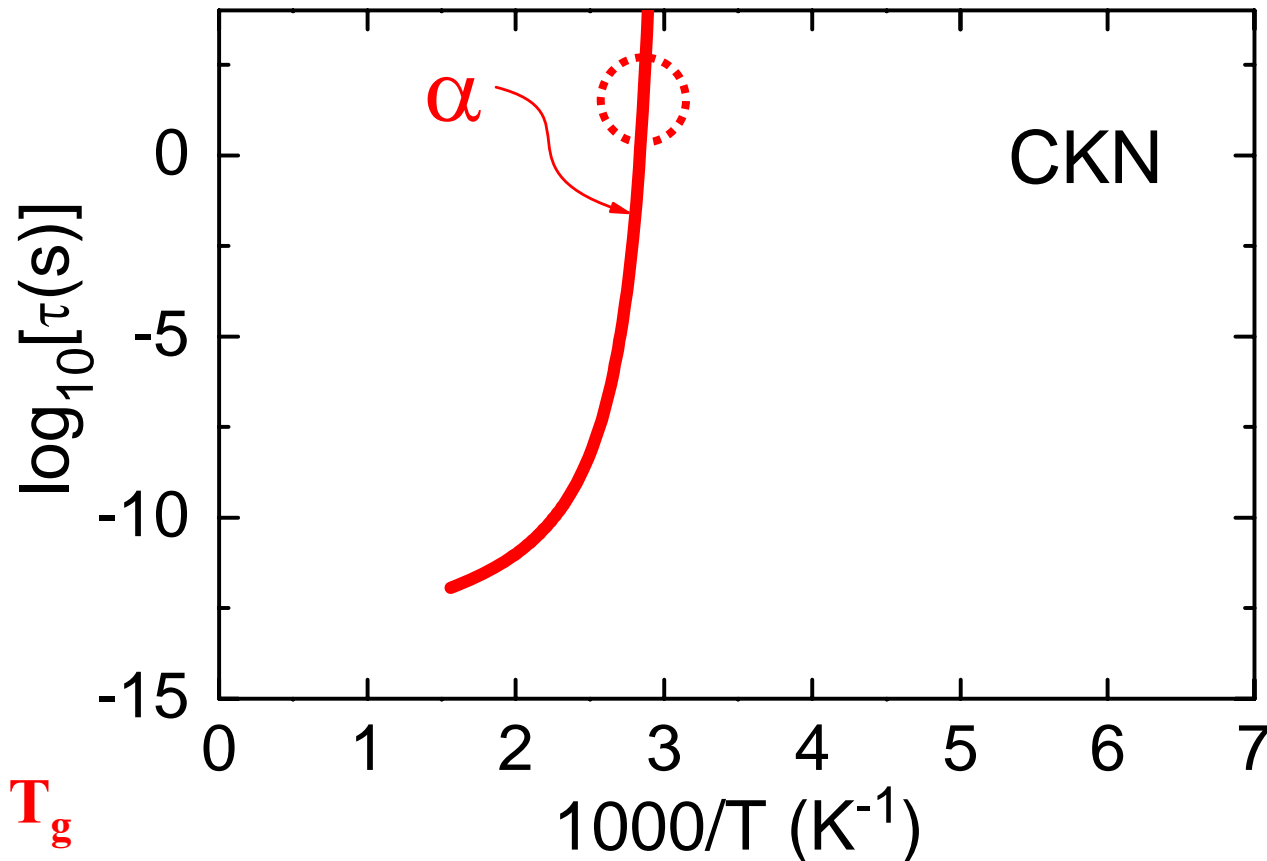


Pimenov, Lunkenheimer, Rall, Kohlhaas, Loidl,
Böhmer, Phys. Rev. E **54**, 676 (1996)



Structural

relaxation



$$T \geq T_g$$

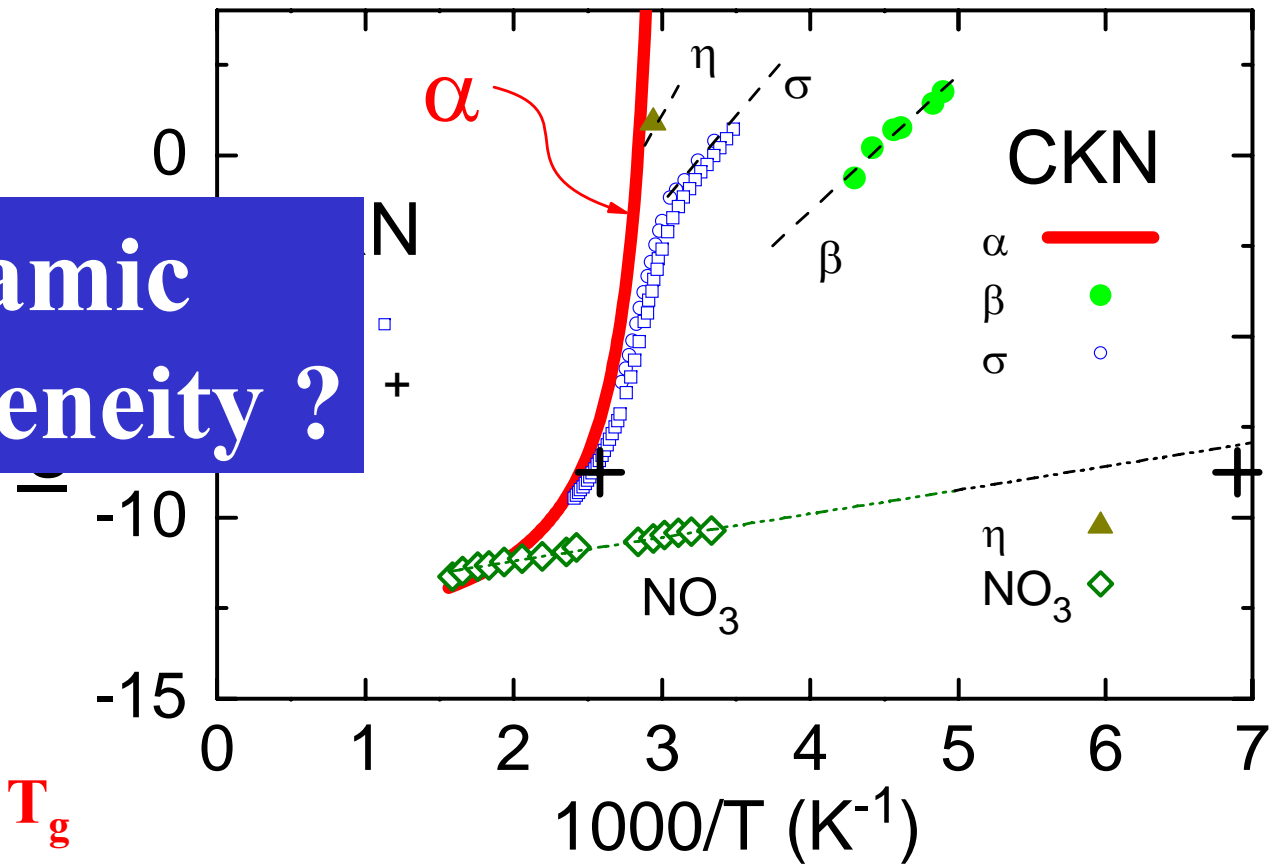
large slope in Angell plot

fragility $m = \left. \frac{d(\log_{10} \tau/s)}{d(T_g/T)} \right|_{T=T_g}$

Böhmer, Ngai, Angell, Plazek, JCP 99, 4201 (1993)

Structural and decoupled relaxations

dynamic heterogeneity ?



$T \geq T_g$

large slope in Angell plot

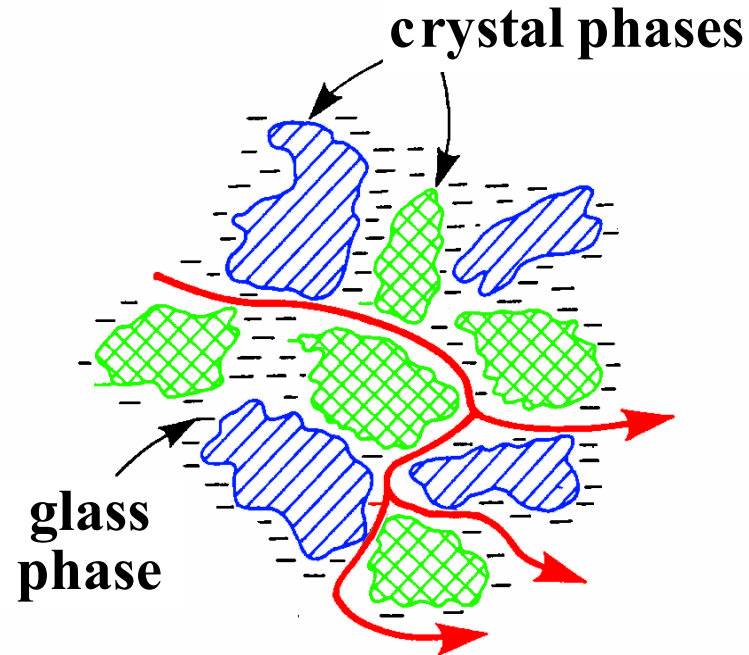
fragility $m = \left. \frac{d(\log_{10} \tau/s)}{d(T_g/T)} \right|_{T=T_g}$

$T < T_g$
thermally activated processes

Origin of non-exponentiality

Nanoscale heterogeneity in ionic conductors ?

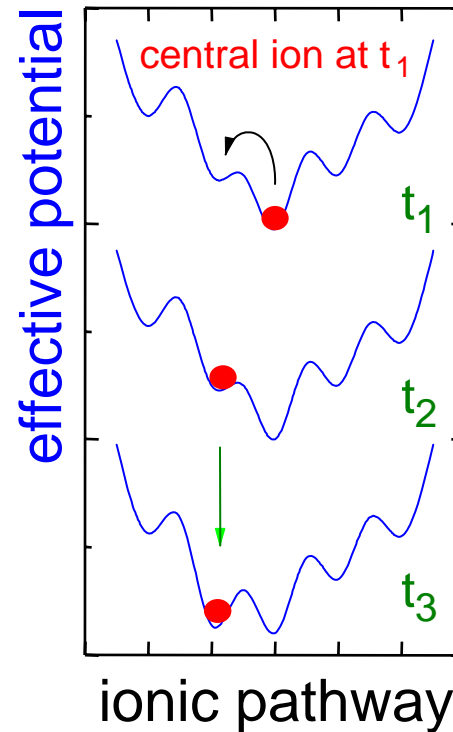
heterogeneous cluster model:



glass ceramics, nanocrystals,
ion conductors in general ??

or correlation effects ?

homogeneous jump-relaxation-model:
Coulomb interactions



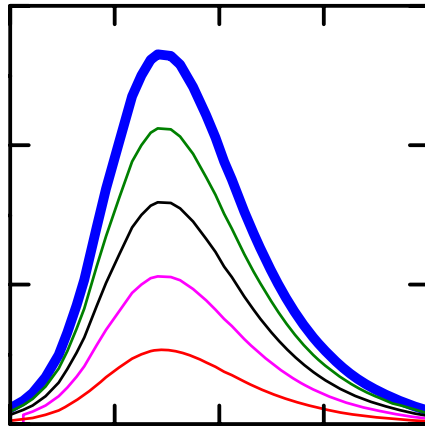
effectively time
dependent
rates $W(t)$



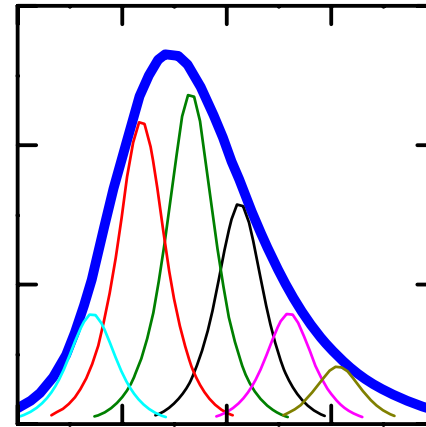
time evolution
 $\exp[-t W(t)]$
 $\propto \exp[-(t/\tau)^\beta]$

**frequency
domain**

HOMOGENEOUS



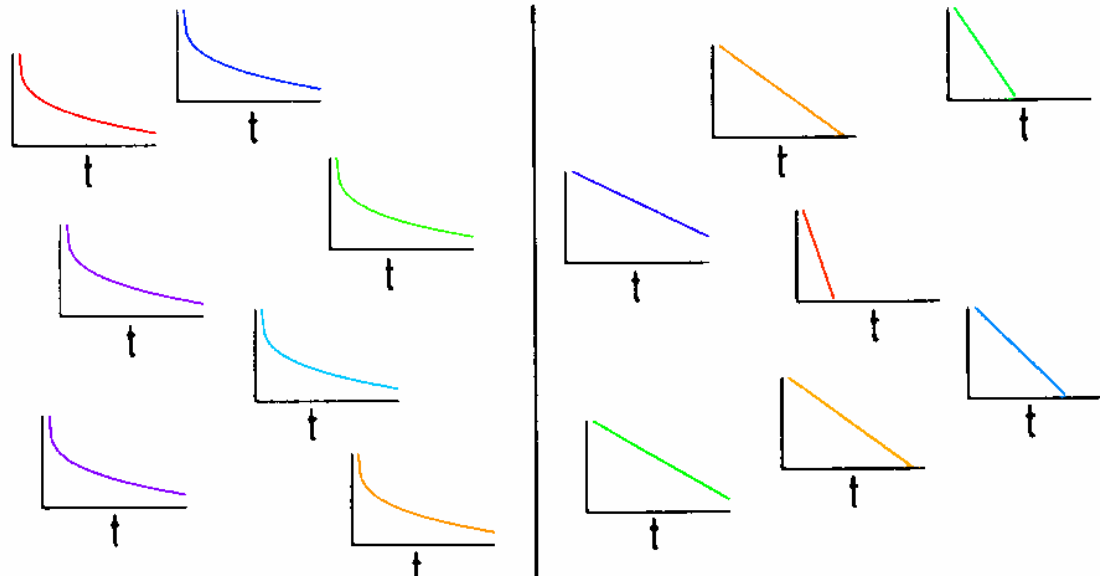
HETEROGENEOUS



χ''

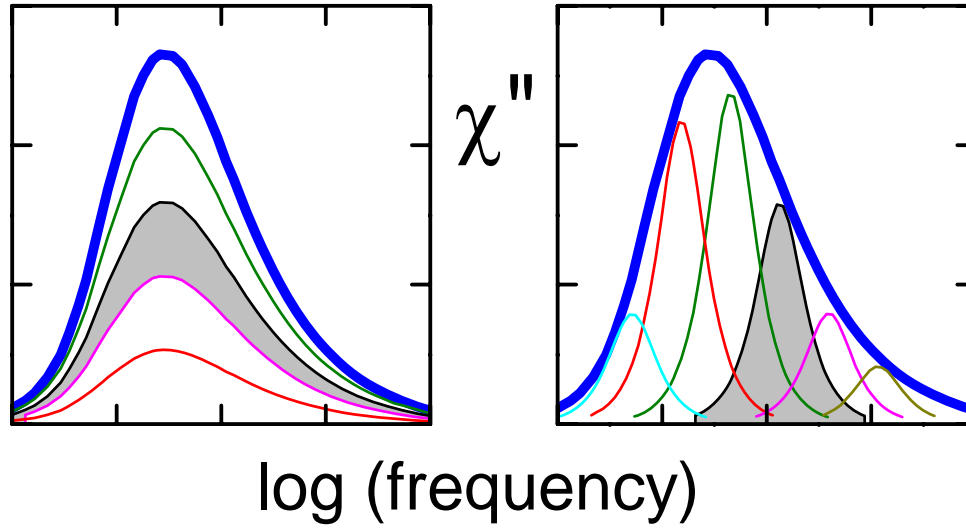
$\log(\text{frequency})$

**time
domain**

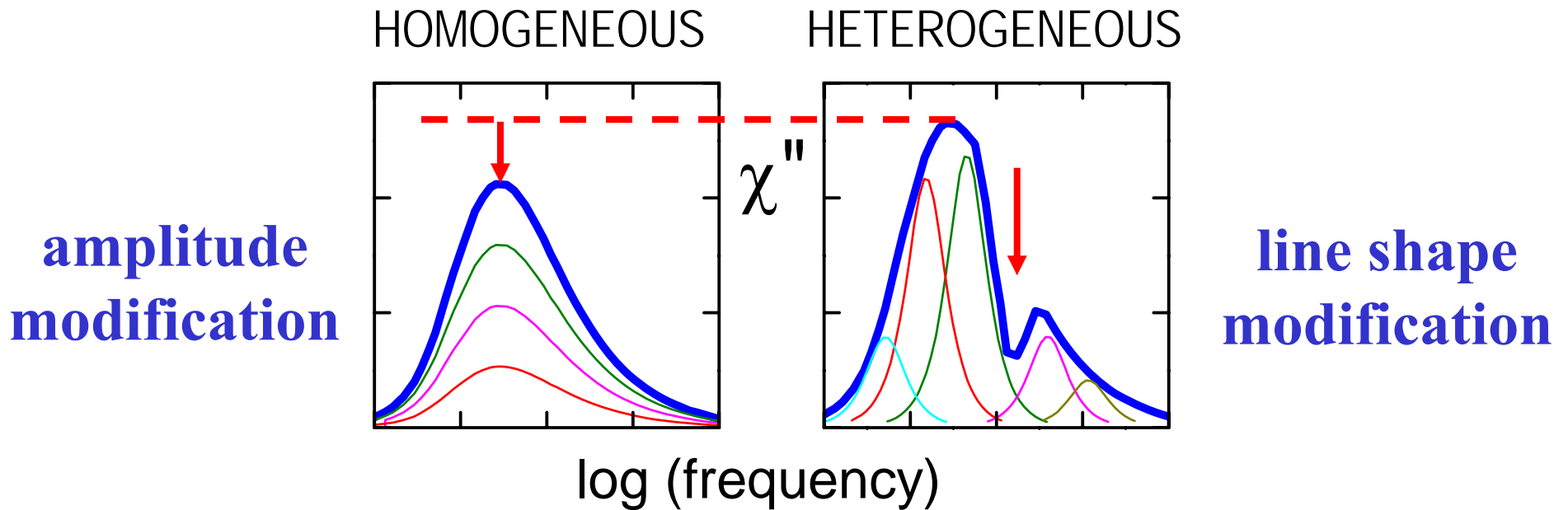


HOMOGENEOUS

HETEROGENEOUS



mark a subensemble



Selection experiments

1. **spectral or spatial selection
or excitation of sub-ensemble**

2. **detection of
sub-ensemble**

Nonresonant hole burning

supercooled liquids : Schiener, Loidl, Böhmer, [Chamberlin](#), Science 274, 752 (1996)



Nonresonant hole burning

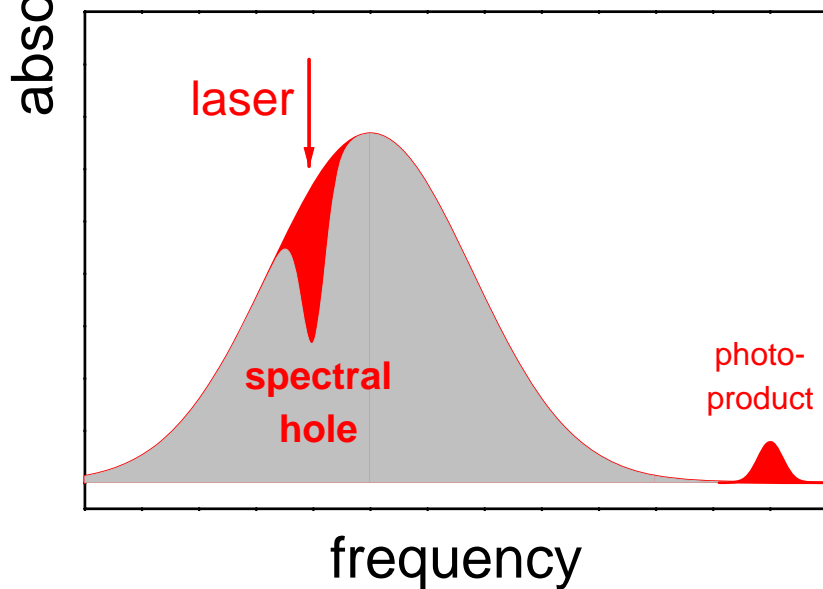
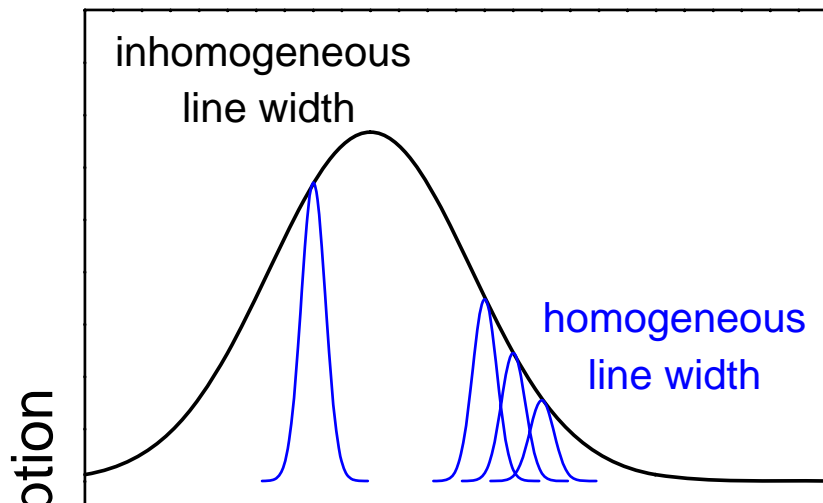
supercooled liquids	α : Schiener, Loidl, Böhmer, Chamberlin, Science 274, 752 (1996)
	β : Richert, EPL 54, 767 (2001)
	wing: Jeffrey, Richert, Duvvuri, JCP 119, 6150 (2003)
relaxor ferroelectrics	Kircher, Schiener, Böhmer, PRL 81, 4520 (1998)
plastic crystal	Wirsch, Kircher, Böhmer (1998, unpublished,*)
magnetic (spin glasses)	Chamberlin, PRL 83, 5134 (1999)
ion conductor CKN	Richert, Böhmer, PRL 83, 4337 (1999)
quantum paraelectrics	Kleemann et al., Ferroelectrics 261, 43 (2001)
stochastic models	Diezemann, EPL 53, 604 (2001)
electrical circuit analog	Richert, Physica A 322, 143 (2003)
binary glass-formers	Blochowicz, Rössler, JCP 122, 224511 (2005)
mechanical (polymers)	Shi, McKenna, PRL 94, 157801 (2005)
*review	Böhmer, Diezemann, in: Broadband Dielectric Spectroscopy edited by Kremer, Schönhals (Springer, 2002), p. 523-569

} solids

Nonresonant hole burning

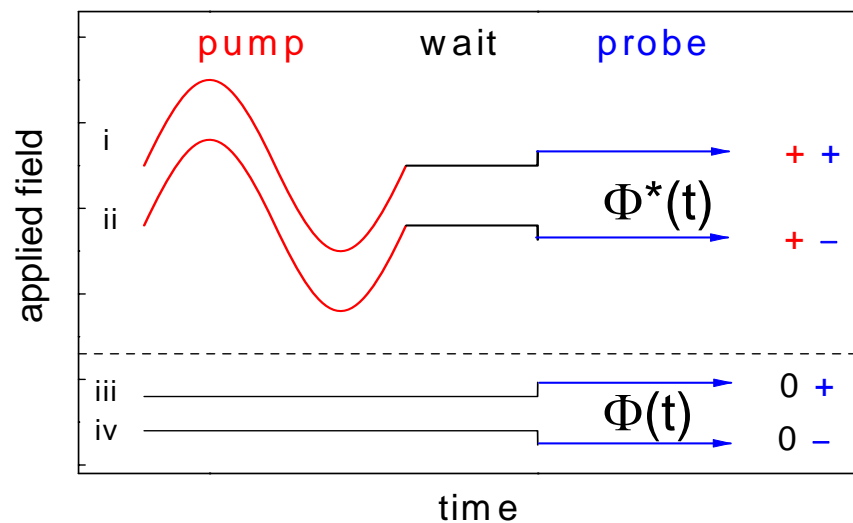
- supercooled liquids α : Schiener, Loidl, Böhmer, Chamberlin, Science 274, 752 (1996)
 β : Richert, EPL 54, 767 (2001)
wing: Jeffrey, Richert, Duvvuri, JCP 119, 150 (2003)
- relaxor ferroelectrics Kircher, Schiener, Böhmer, PRL 81, 4520 (1998)
- plastic crystal Werscher, Zimm, Böhmer (1998, unpublished)
- magnetic (spin glasses) Chamberlin, PRL 83, 5134 (1999)
- ion conductor C Richert, Böhmer, PRL 83, 33 (1999)
- quantum paraelectrics Kleemann, et al., Ferroelectrics 261, 43 (2001)
- stochastic models Diezemann, Richert, EPL 53, 604 (2001)
- electrical circuit analog Richert, Physica A 322, 143 (2003)
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- mechanical (polymers) Shi, McKenna, PRL 94, 157801 (2005)
- review Böhmer, Diezemann, in: Broadband Dielectric Spectroscopy edited by Kremer, Schönhals (Springer, 2002), p. 523-569

conventional SHB

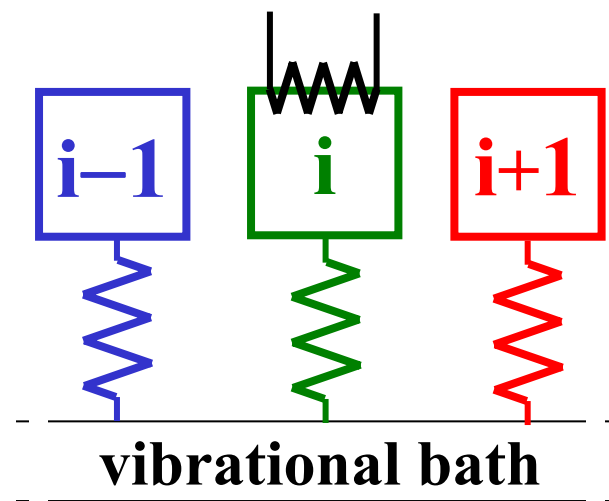


NMR (optical) spectroscopy

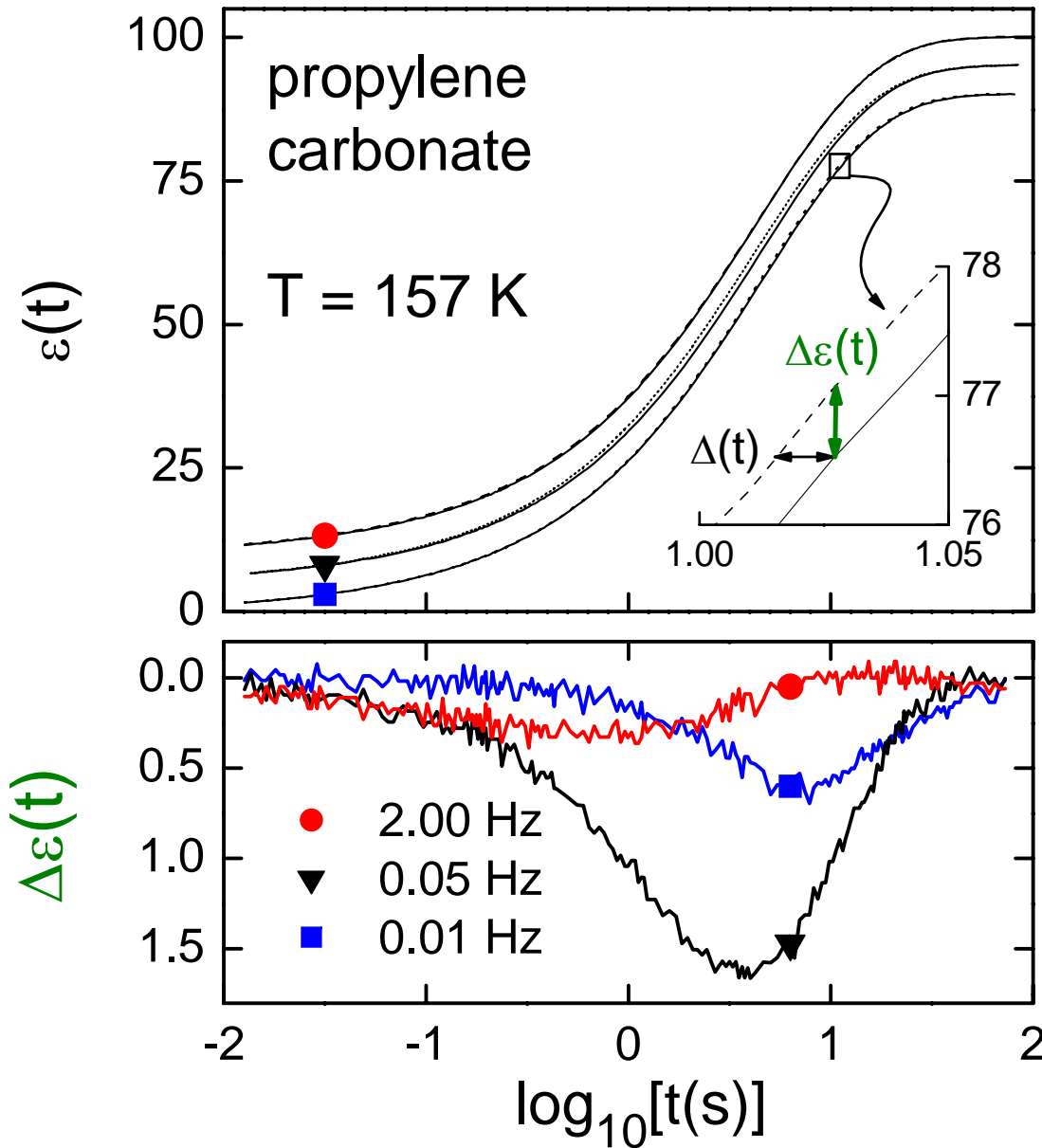
nonresonant (NHB)



phase cycle: $\Delta\Phi(t) = \Phi^*(t) - \Phi(t)$



10 years ago: First results on a viscous liquid



**high-frequency pump:
short-time modifications**

**low-frequency pump:
long-time modifications**



selection of different
sub-ensembles possible:
dynamic heterogeneity

Schiener, Loidl, Böhmer,
Chamberlin, Science 274,
752 (1996)

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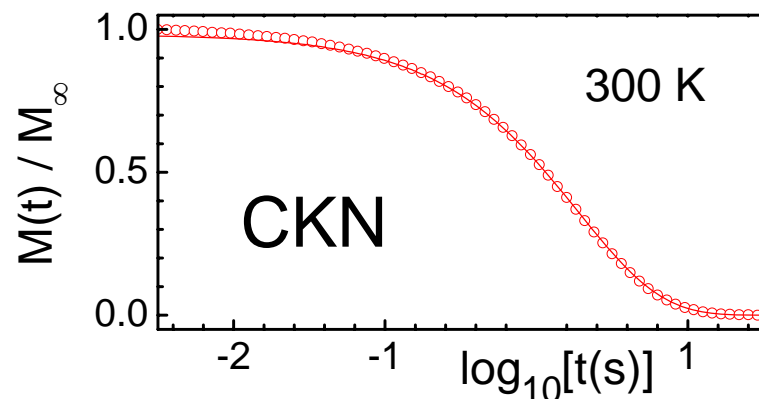
3. Nonresonant spectral hole burning

Material and method

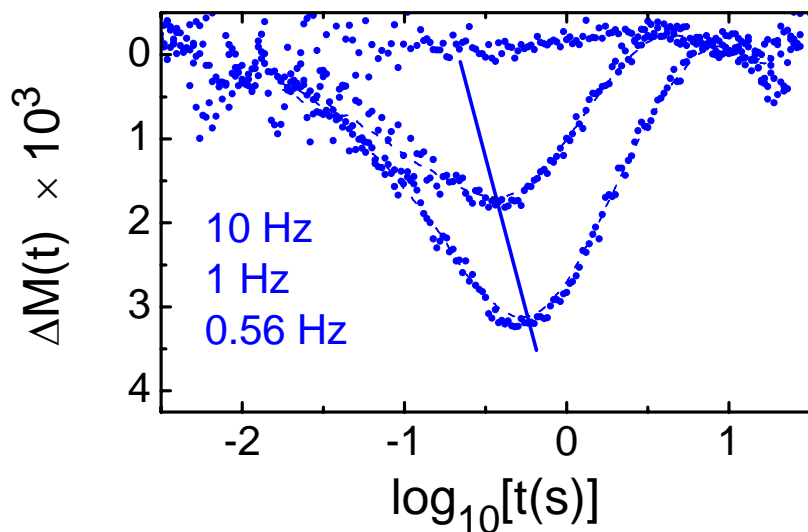
Experiments on Ca-K-NO_3 glass

4. Conclusions

high pump frequencies

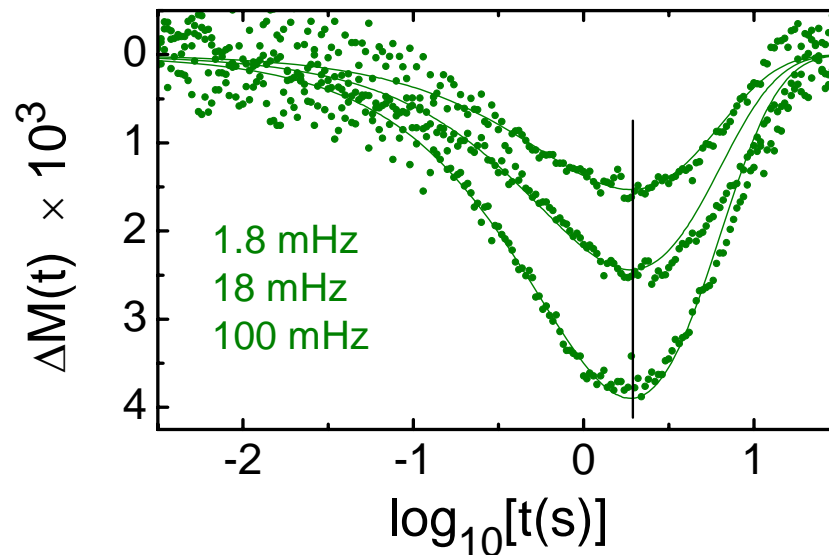


low pump frequencies



→ frequency selectivity

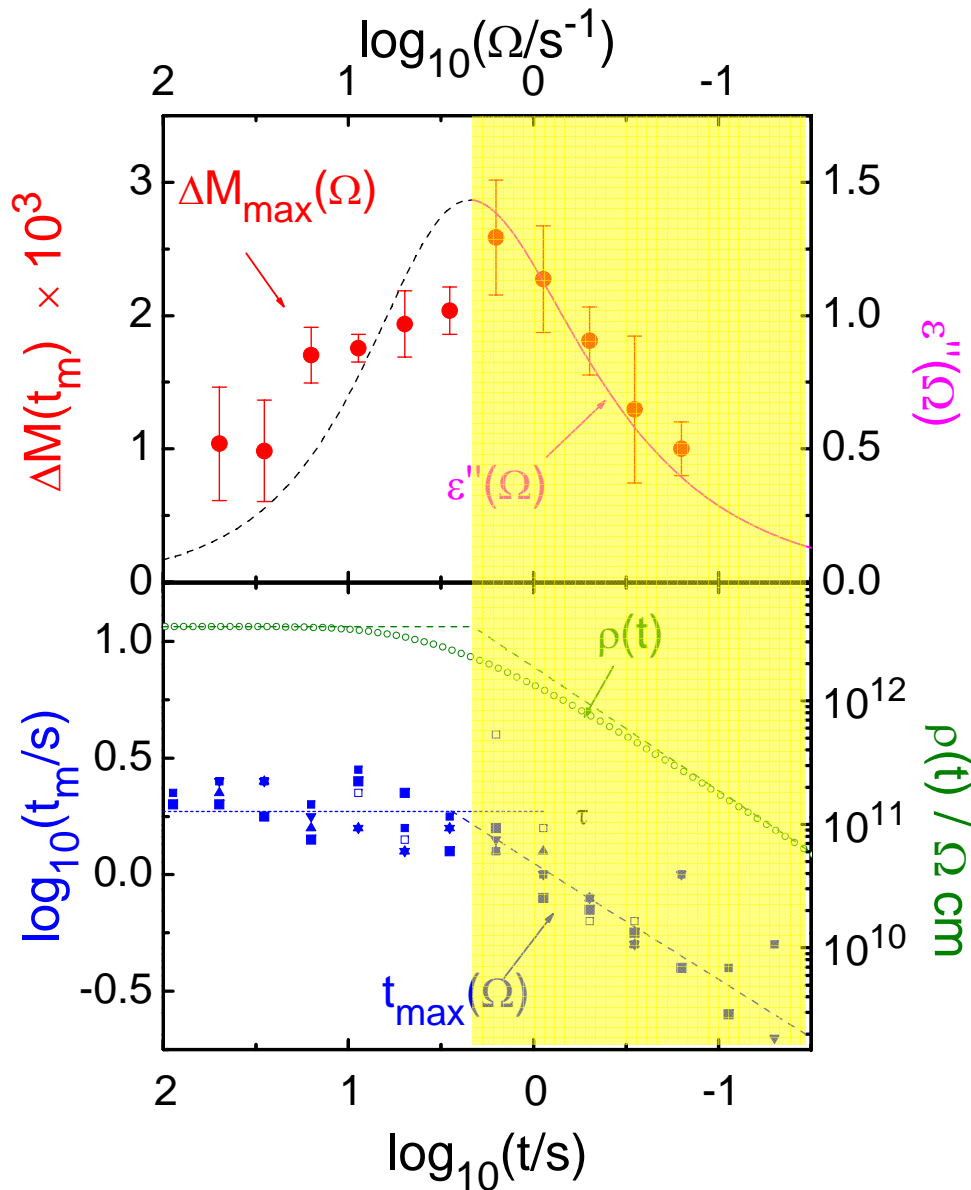
dynamic heterogeneity



→ no frequency selectivity

homogeneous for $\Omega_C/2\pi < 0.1$ Hz

CKN – pump frequency dependence



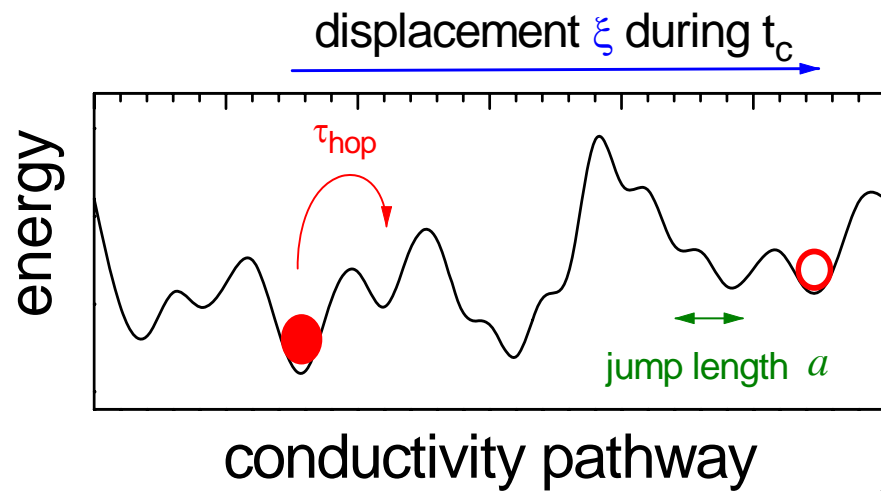
heterogeneous regime:
 hole depths \propto dielectric loss
 frequency selective
 energy absorption

OBSERVATION:
 t_c corresponds to transition from
 dc to ac resistivity ρ (or σ)



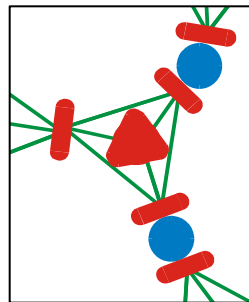
ion needs $t_c \sim 1/\Omega_c$
 to average over heterogeneity

heterogeneity length ξ ?



$$\left. \begin{aligned} \langle \xi^2 \rangle &= 6 D t_c \\ D &= a^2 / (6 \tau_{hop}) \end{aligned} \right\} \langle \xi^2 \rangle = a^2 t_c / \tau_{hop}$$

a hopping length \approx
mean separation of
sites for mobile ions



nonresonant
hole burning spectroscopy
at low frequencies

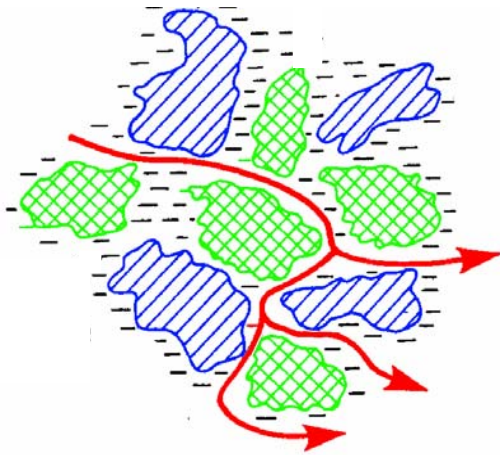


dynamic and spatial
heterogeneity in CKN

τ_{hop} hopping time necessary
to estimate ξ but
not known for CKN

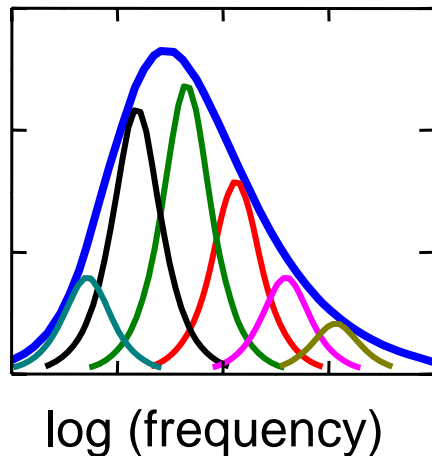
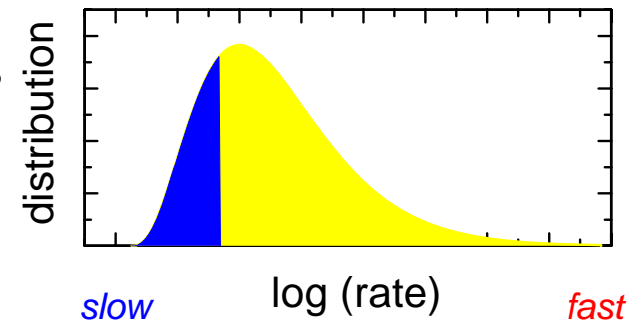
requires direct measurement
of τ_{hop} in the slow domain !

Conclusions



Structure dynamic relationship in crystalline and glassy lithium aluminosilicates

High magnetic fields facilitate investigations of dynamic heterogeneities in ion conductors also with less favorable nuclear probes (Ag)



Nonresonant spectral hole burning is a useful tool for different classes of materials.

Direct observation of transition from heterogeneous to homogeneous behavior in CKN