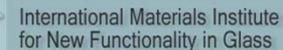


Advanced Vitreous State – The Physical Properties of Glass



Dielectric Properties of Glass

Lecture 2: Dielectric in an AC Field

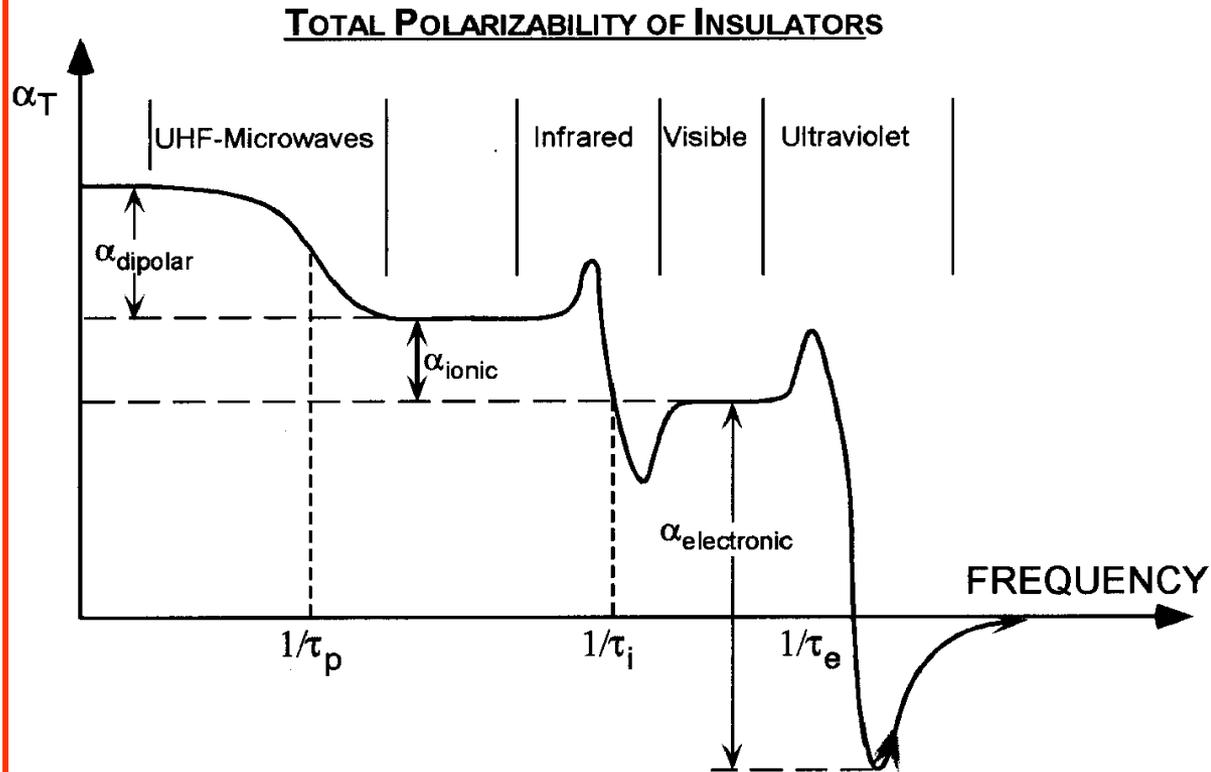
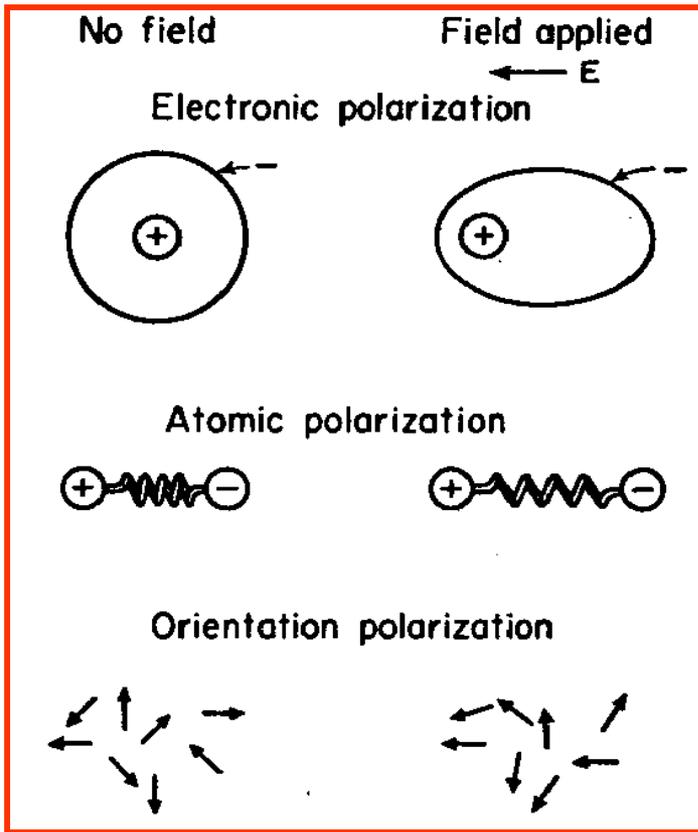
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Classic sources of polarizability in glass vs. frequency

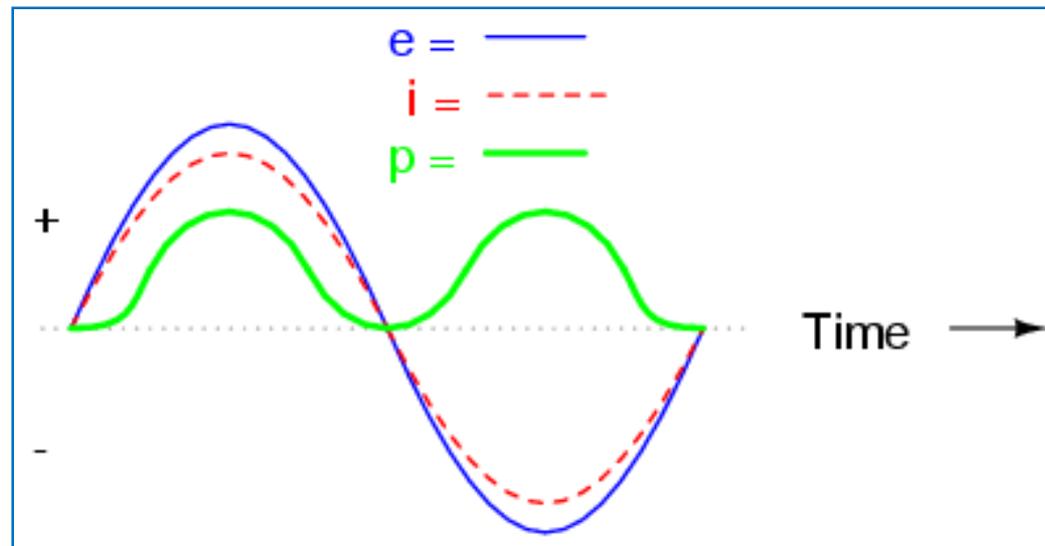
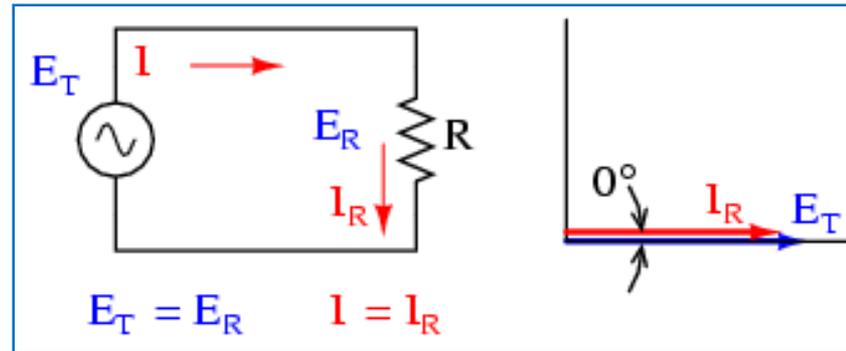


$$\alpha = \alpha_e + \alpha_a + \alpha_d$$

At optical frequencies $(\epsilon_r - 1)/(\epsilon_r + 2) = \frac{n^2 - 1}{n^2 + 2} = \frac{Nq_e^2}{3\epsilon_0 m} \sum_j \frac{f_j}{\omega_{oj}^2 - \omega^2 + i\Gamma_j\omega}$

Dielectric in AC Field: Macroview i.e. a bit of EE

$$E = E_0 \sin \omega t$$
$$i = i_0 \sin \omega t$$

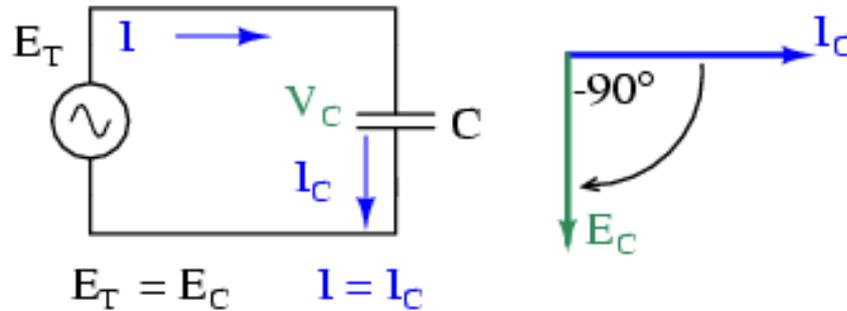


Voltage and current are “in phase” for resistive circuit. Power or energy loss, p , $\propto Ri^2$

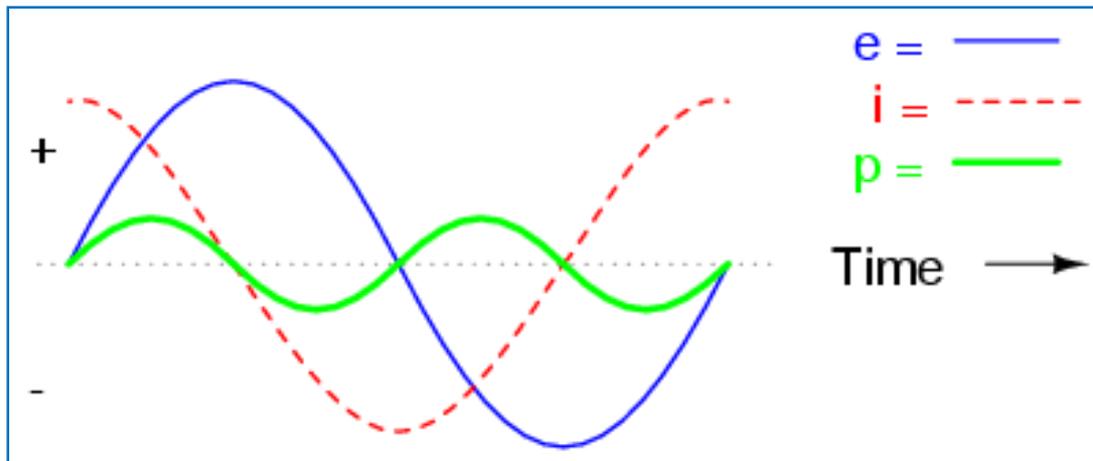
http://www.ibiblio.org/kuphaldt/electricCircuits/AC/AC_6.html

Ideal vs. Real Dielectric

$$E = E_0 \sin \omega t$$
$$i = i_0 \cos \omega t$$



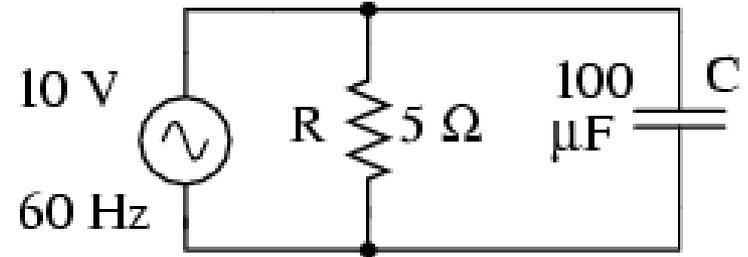
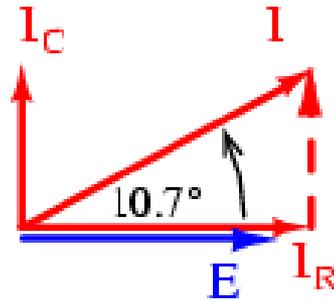
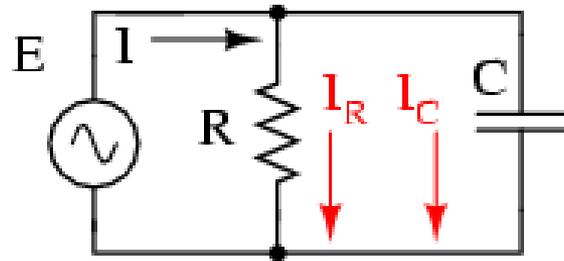
C is frequency independent



In an ideal dielectric, current is ahead of voltage (or voltage lags behind the current) by 90°.

The power is positive or negative, average being zero i.e. there is no energy loss in a perfect dielectric.

Real dielectric: A parallel circuit of R and C



$$I = I_R + I_C$$

$$E = E_R = E_C$$

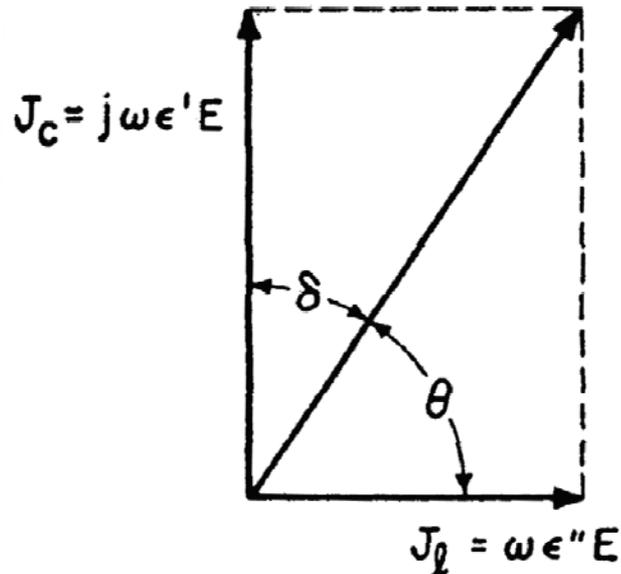
The total current can be considered as made of a lossy resistive component, I_L (or I_R) that is in-phase with voltage, and a capacitive current, I_C , that is 90° out-of-phase.

Unlike ideal dielectric, real dielectric has finite conductivity that causes loss of energy per cycle.

In this case, the current is ahead of voltage by $<90^\circ$.

Complex Relative Permittivity

$$\epsilon_r = \epsilon'_r - j\epsilon''_r$$



There are many parameters to represent the dielectric response (permittivity (ϵ^), susceptibility (χ^*), conductivity (σ^*), modulus (M^*), impedance (Z^*), admittance (Y^*), etc.) emphasizing different aspects of the response. However, they are all interrelated mathematically. One needs to know only the real and imaginary parts of any one parameter.*

Charging and loss current density.

ϵ_r = dielectric constant

ϵ'_r = real part of the complex dielectric constant

ϵ''_r = imaginary part of the complex dielectric constant

j = imaginary constant $\sqrt{-1}$

$$\epsilon^*(\omega, T) = \epsilon' - j[\sigma(\omega, T)/\omega]$$

$$\epsilon_0 \epsilon_r''(\omega, T) = \sigma'(\omega, T)/\omega$$

Energy loss in a dielectric

Loss tangent or loss factor $\tan \delta = \frac{\epsilon_r''}{\epsilon_r'}$

Describes the losses in relation to dielectric's ability to store charge.

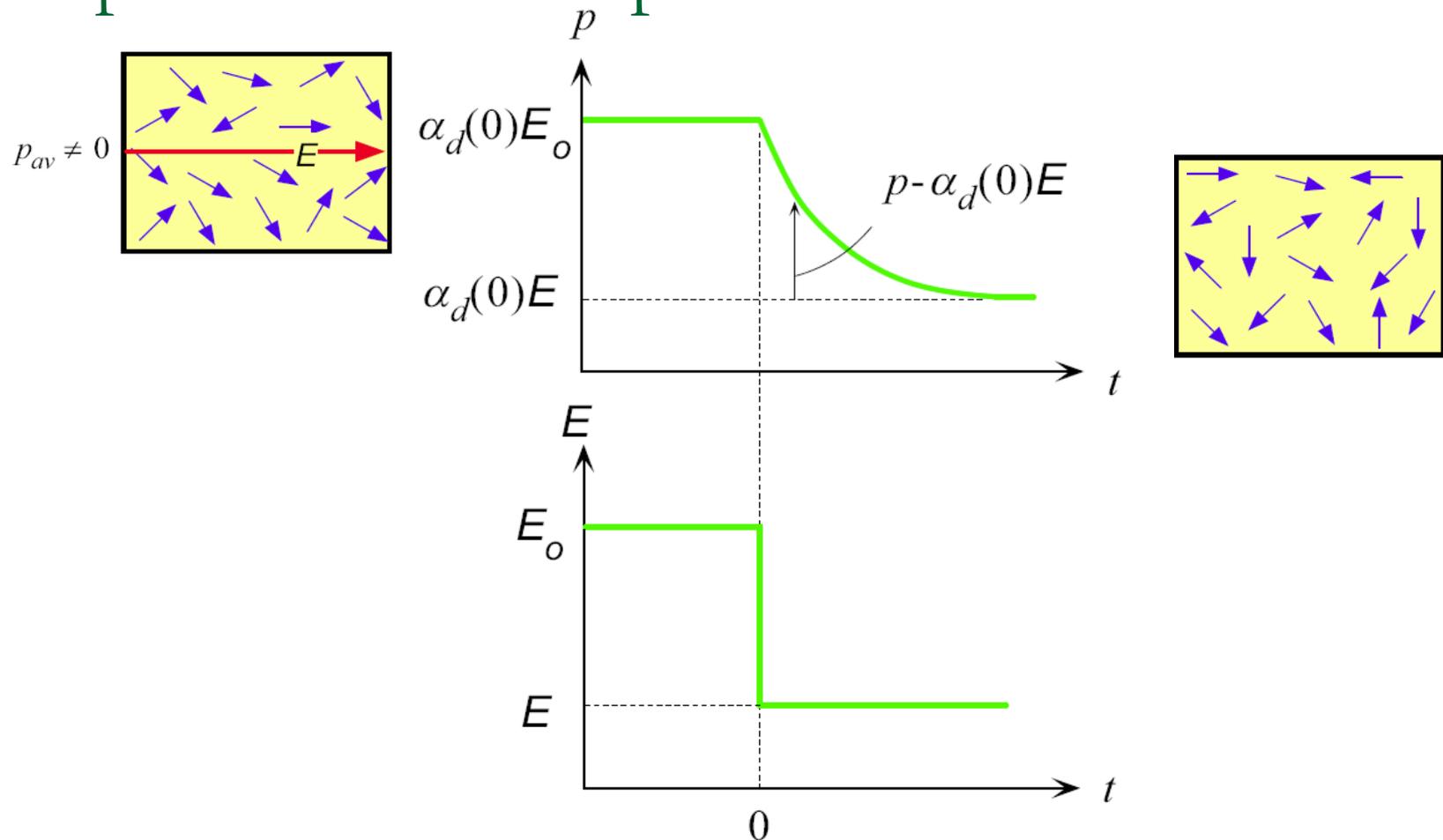
Energy absorbed or loss/volume-sec

$$W_{\text{vol}} = \omega E^2 \epsilon_o \epsilon_r'' = \omega E^2 \epsilon_o \epsilon_r' \tan \delta$$

Loss tangent of silica is 1×10^{-4} at 1 GHz, but can be orders of magnitude higher for silicate glass (Corning 7059) = 0.0036 @ 10 GHz.

Depends on ω and T.

Depolarization of dipolar dielectric



The dc field is suddenly changed from E_0 to E at time $t = 0$. The induced dipole moment p has to decrease from $\alpha_d(0)E_0$ to a final value of $\alpha_d(0)E$. The decrease is achieved by random collisions of molecules in the gas.

Dipolar Relaxation Equation

$$\frac{dp}{dt} = -\frac{p - \alpha_d(0)\mathcal{E}}{\tau}$$

p = instantaneous dipole moment = $\alpha_d \mathbf{E}$,

dp/dt = rate at which p changes, α_d = dipolar orientational polarizability,

\mathcal{E} = electric field, τ = relaxation time

When AC field $\mathbf{E} = \mathbf{E}_0 \exp(j\omega t)$, the solution for p or α_d vs. ω :

$$\alpha_d(\omega) = \frac{\alpha_d(0)}{1 + j\omega\tau}$$

ω = angular frequency of the applied field, j is $\sqrt{-1}$.

Debye Equations

$$\epsilon'_r = 1 + \frac{[\epsilon_r(0) - 1]}{1 + (\omega\tau)^2} \quad \epsilon''_r = \frac{[\epsilon_r(0) - 1]\omega\tau}{1 + (\omega\tau)^2}$$

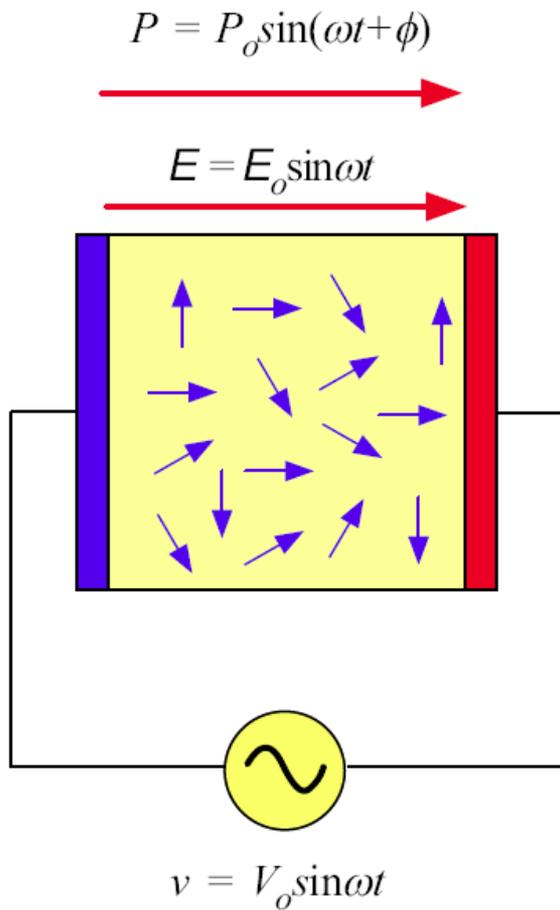
ϵ_r = dielectric constant (complex)

ϵ'_r = real part of the complex dielectric constant

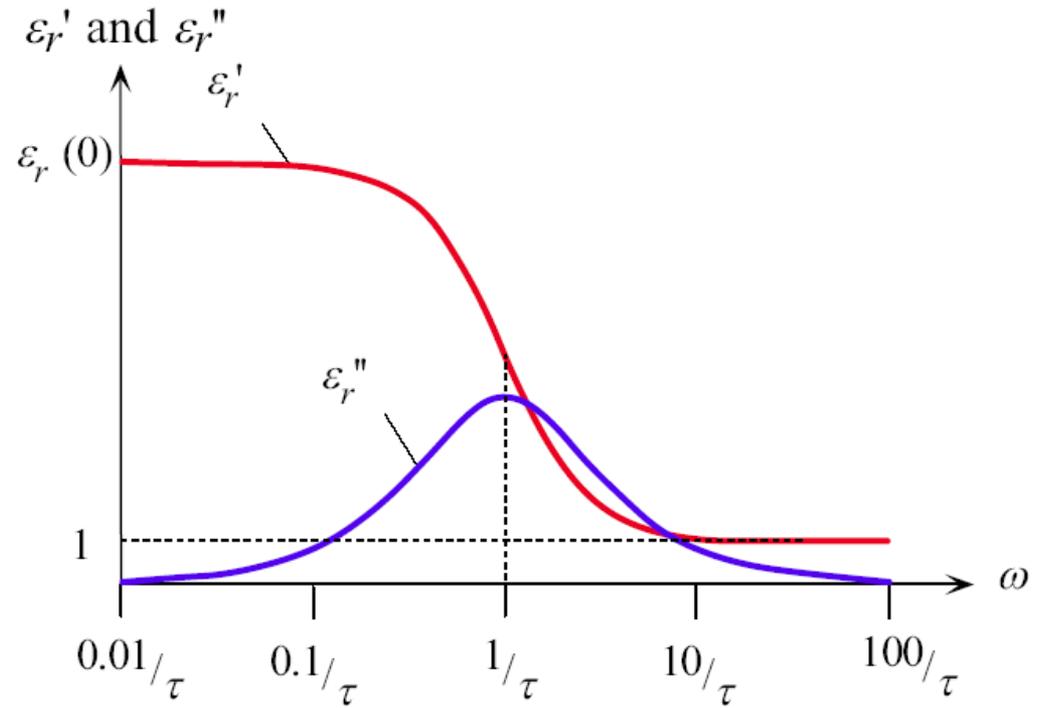
ϵ''_r = imaginary part of the complex dielectric constant

ω = angular frequency of the applied field

τ = relaxation time



(a)



(b)

(a) An ac field is applied to a dipolar medium. The polarization $P(P = Np)$ is out of phase with the ac field.

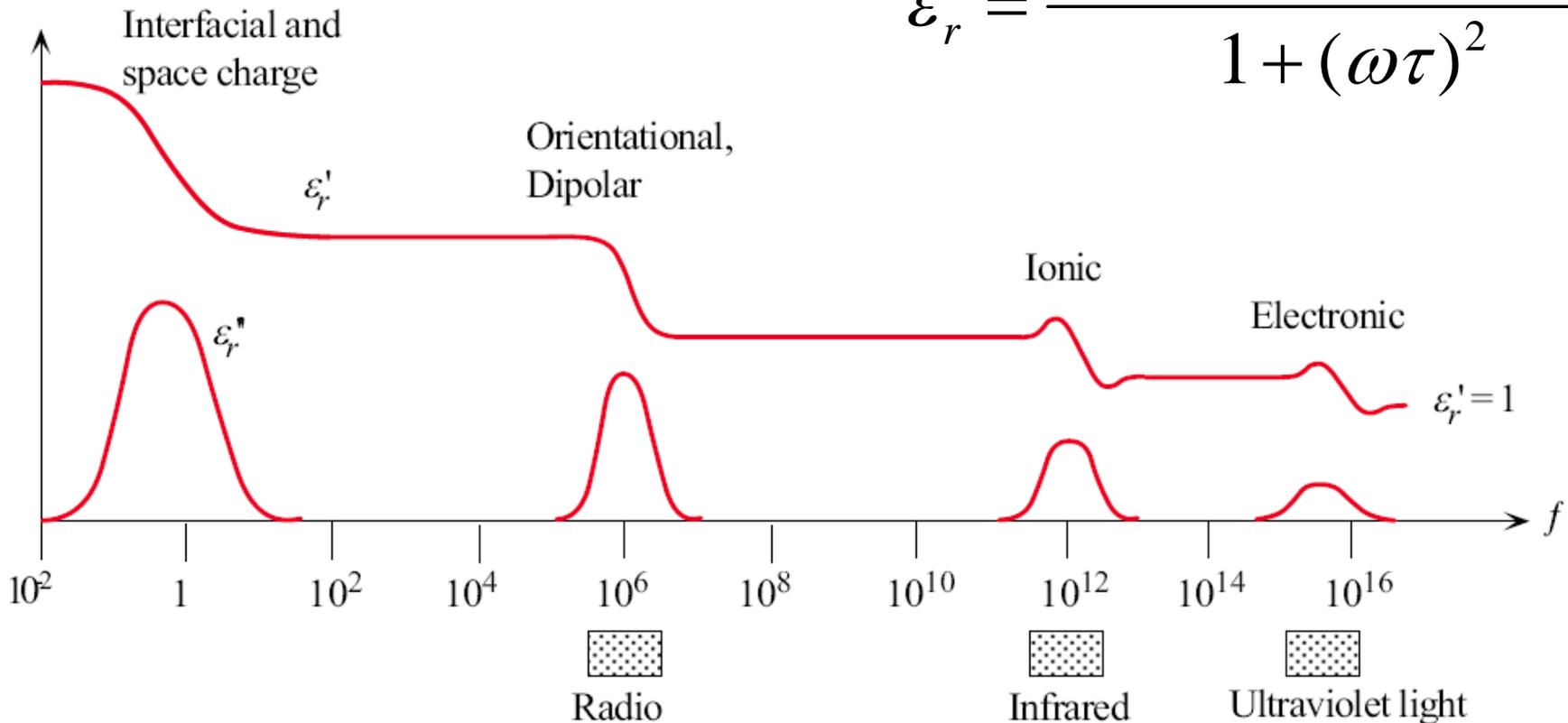
(b) The relative permittivity is a complex number with real (ϵ_r') and imaginary (ϵ_r'') parts that exhibit frequency dependence.

Dielectric constant over broad frequency range

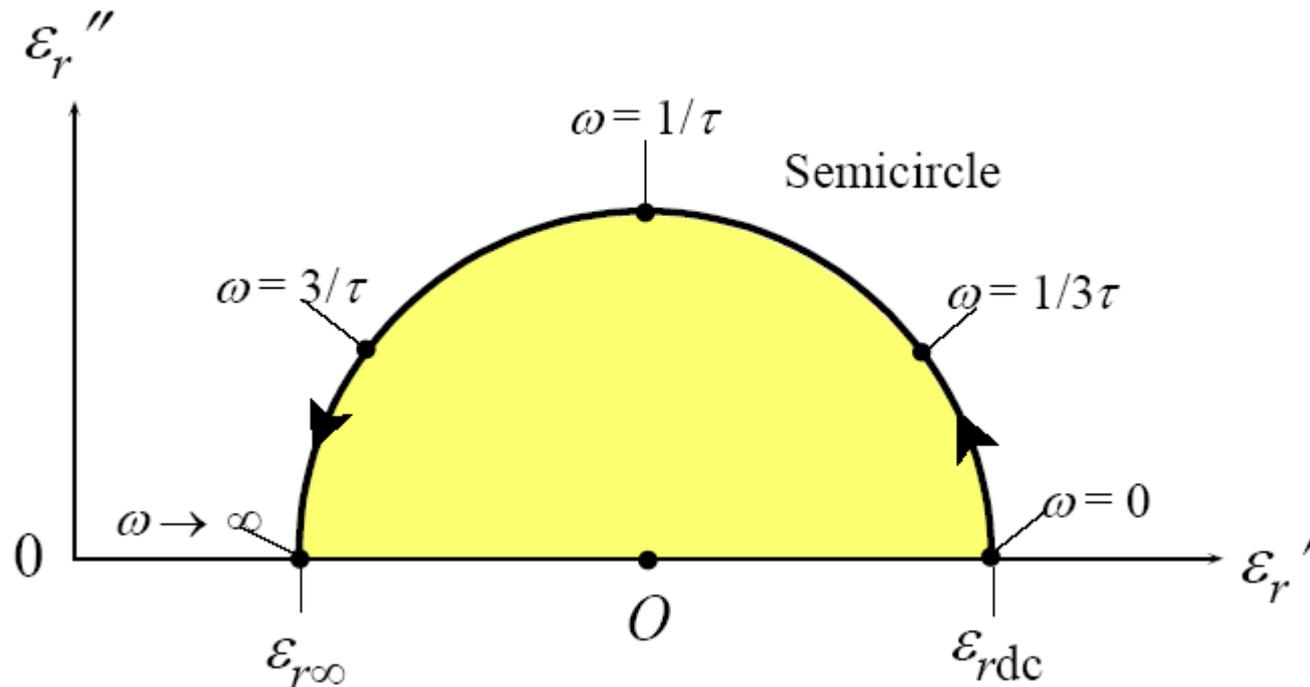
Dipolar contribution, typically below GHz range

$$\epsilon'_r(\omega) = \epsilon_r(opt) + \frac{[\epsilon_r(dc) - \epsilon_r(opt)]}{1 + (\omega\tau)^2}$$

$$\epsilon''_r = \frac{[\epsilon_r(dc) - \epsilon_r(opt)]\omega\tau}{1 + (\omega\tau)^2}$$

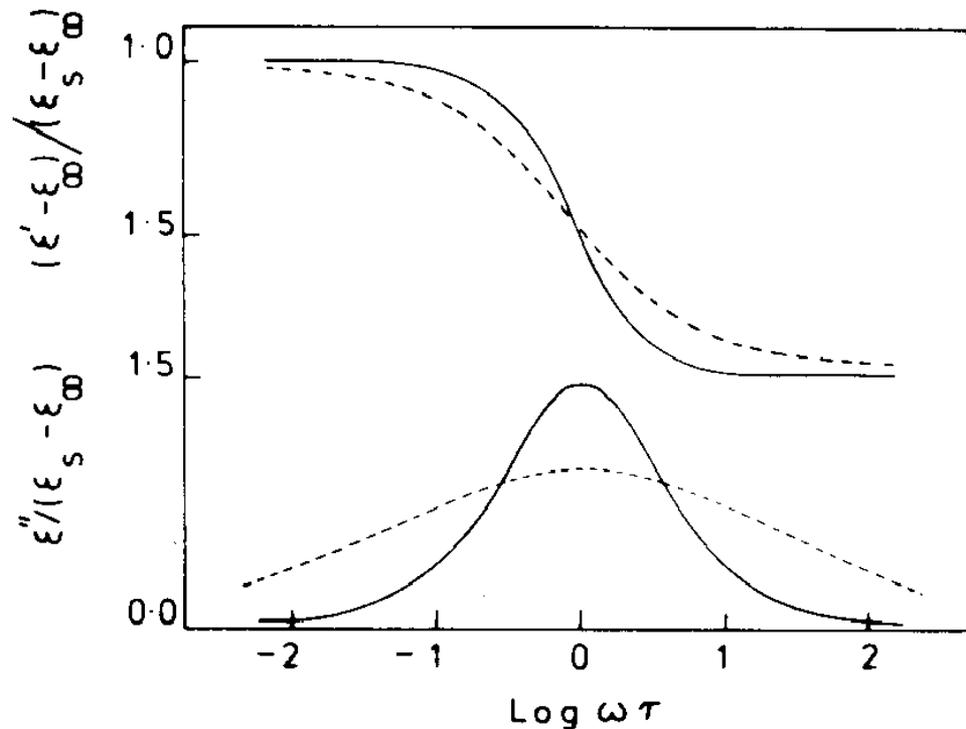


Cole-Cole plots



Cole-Cole plot is a plot of ϵ_r'' vs. ϵ_r' as a function of frequency, ω . As the frequency is changed from low to high frequencies, the plot traces out a circle if Debye equations are obeyed.

Dipolar dielectric loss in complex systems



$\tau = \tau_0 \exp(Q/RT)$
 where Q is activation energy for the reorientation of a dipole.

How would the loss peak change with increasing T ?

Debye Eqs are valid when the dipole (ion) conc is small i.e. non-interacting dipoles, and ϵ'' vs $\log \omega$ shows symmetric Debye peak at $\omega\tau = 1$

For high x , the dipoles interact causing distribution of $\tau \Rightarrow$ the loss peak is smeared.

$$\epsilon^* = \epsilon_\infty + (\epsilon_s - \epsilon_\infty) \int_0^\infty \frac{G(\tau) d\tau}{1 + i\omega\tau}$$

where $G(t)$ is an appropriate distribution function.

An example: $18\text{Na}_2\text{O}-10\text{CaO}-72\text{SiO}_2$ glass

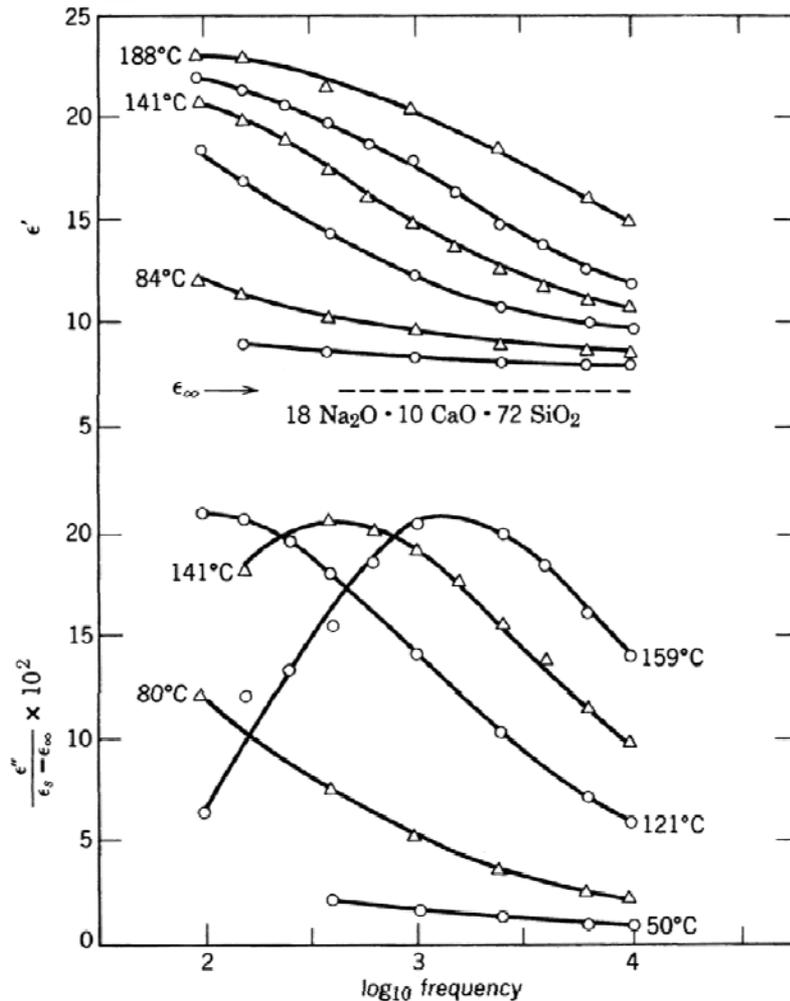


Fig. 18.18. Dielectric dispersion and absorption curves, corrected for dc conductivity, for a typical soda-lime-silicate glass. From H. E. Taylor, *J. Soc. Glass Technol.*, **43**, 124 (1959).

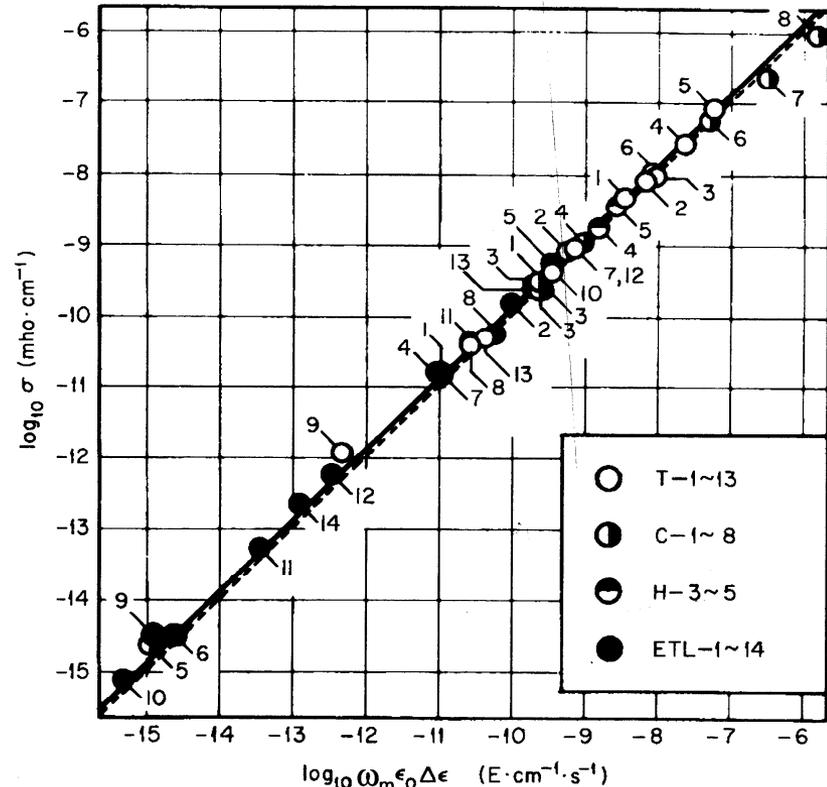
Barton-Nakajima-Namikawa (BNN) relation

$$\sigma_{dc} = p \epsilon_0 \Delta \epsilon \omega_m$$

where p is a constant ~ 1 . $\Delta \epsilon$ is the step in ϵ' across the peak, ω_m is freq of ϵ'' maximum.

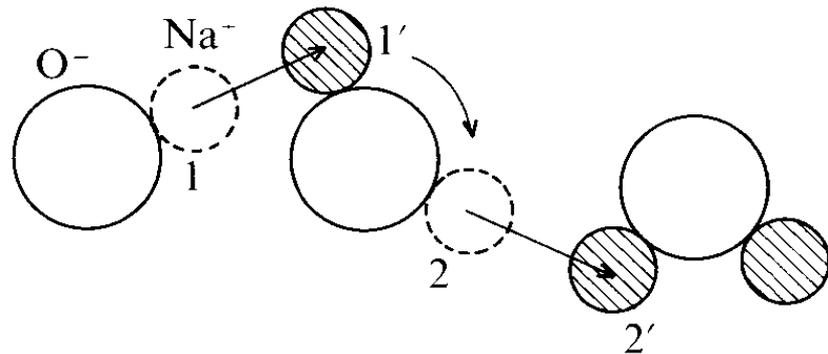
$$\omega_m = 1/\tau = \tau_0^{-1} \exp[-\Delta E_m/RT]$$

Dc conductivity and ϵ'' maximum have same activation energy \Rightarrow common origin.



Correlation between conductivity and dielectric relaxation. \circ = Taylor (1957, 9); \bullet = Charles (1962; 1963); \ominus Heroux (1958); \bullet = measured at Electrotechnical Laboratory, MITI, Japan. (After Nakajima, 1972.)

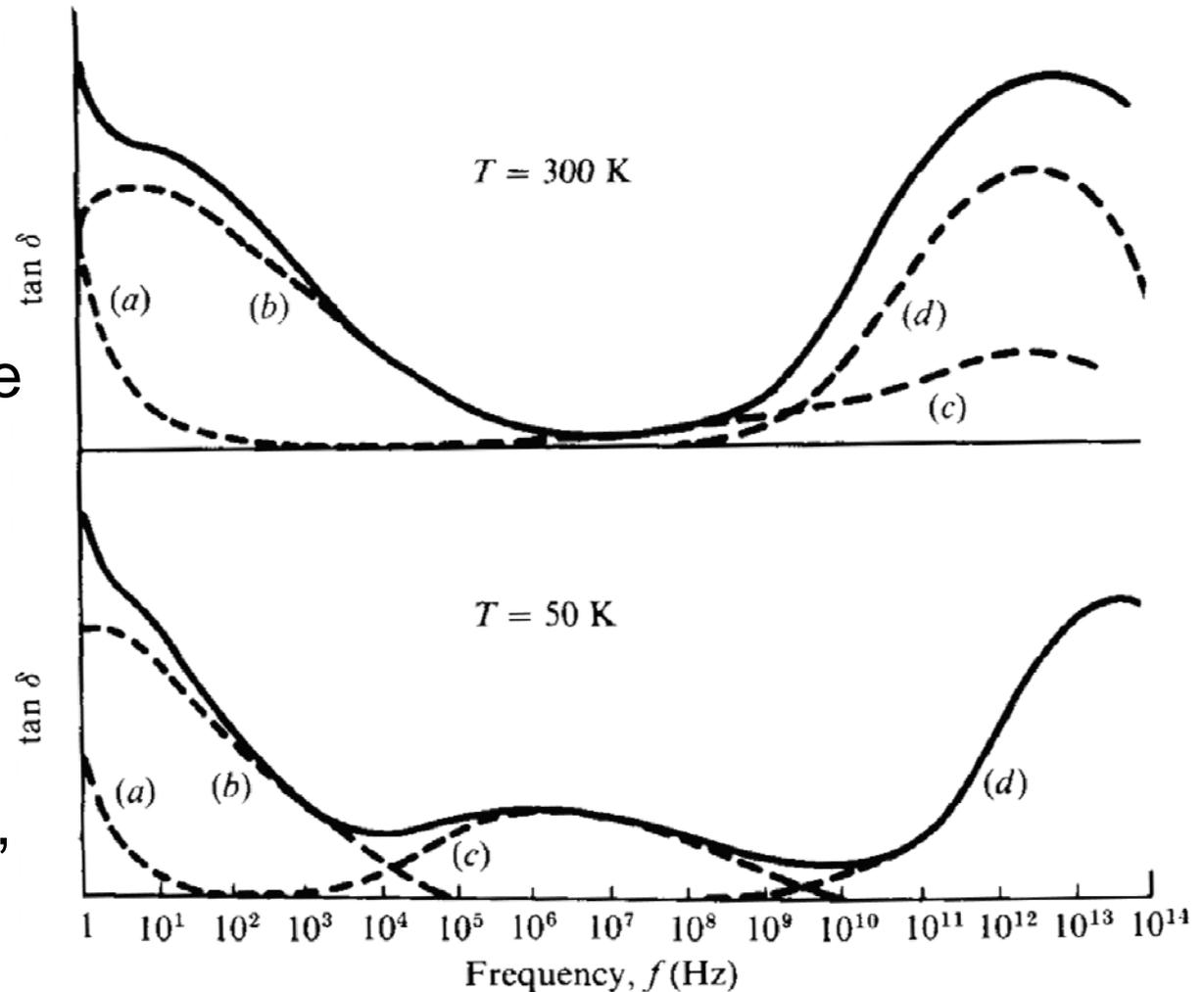
re 15-8. Confirmation of the BNN relationship. The symbols T, C, H, ETL correspond to measurements at four different laboratories cited by Tomozawa⁽¹¹⁾.



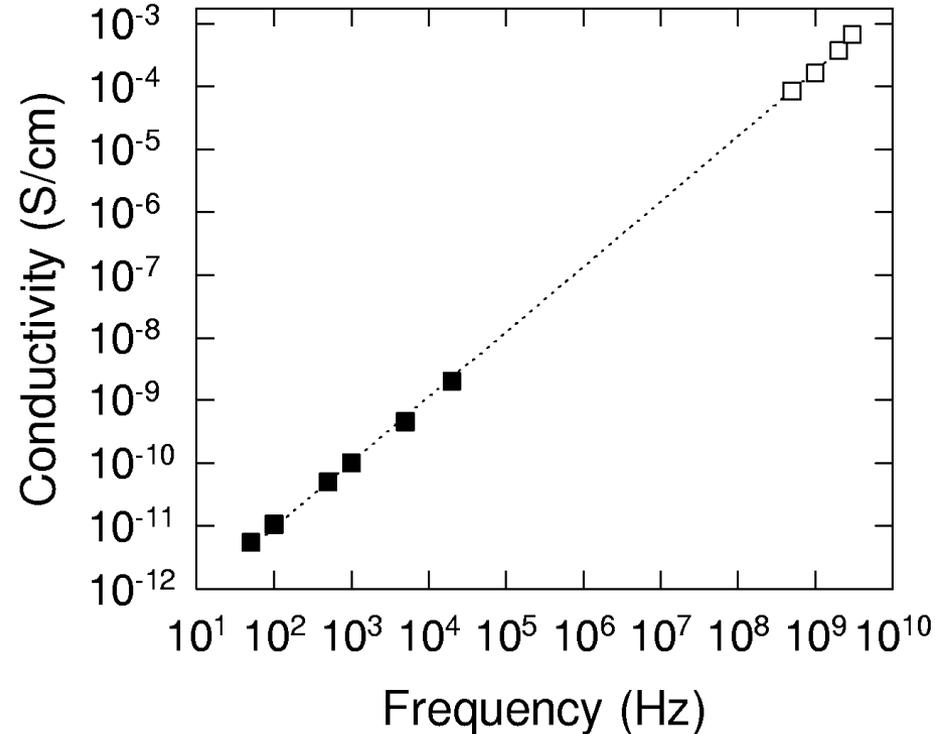
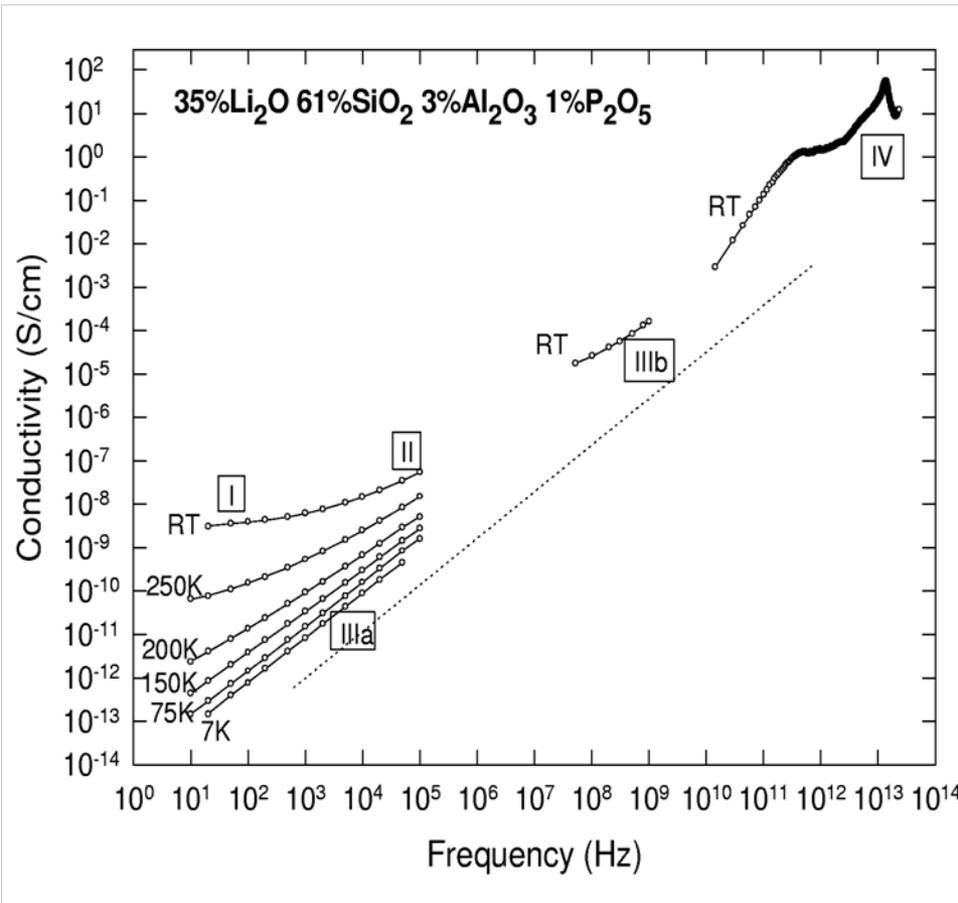
Loss tangent over a wide frequency range

Fig. 11.8. Various types of dielectric losses in a glass: (a) conduction, (b) dipole relaxation, (c) deformation, (d) vibration. After Ref. 404.

Below the visible frequencies, there are at least four different mechanisms that are responsible for dielectric loss in glass: (a) dc conduction, (b) dipole, (c) deformation/jellyfish, (d) vibration



Frequency-temperature interchange



The source of dielectric loss (ac conductivity) at low T – low ω and high T – MW ω has a common underlying origin.

The jellyfish mechanism

- ** It is a **group of atoms** which collectively move between different configurations, much like the wiggling of a *jellyfish* in *glassy ocean*.
- ** There is **no single atom** hopping involved.
- ** The fluctuations are much **slower than typical atom vibrations**.
- ** The exact nature of the 'jellyfish' (ADWPC) depends on the material.
- ** In the same material more than one 'jellyfish' might exist and be observed in different T and f ranges.



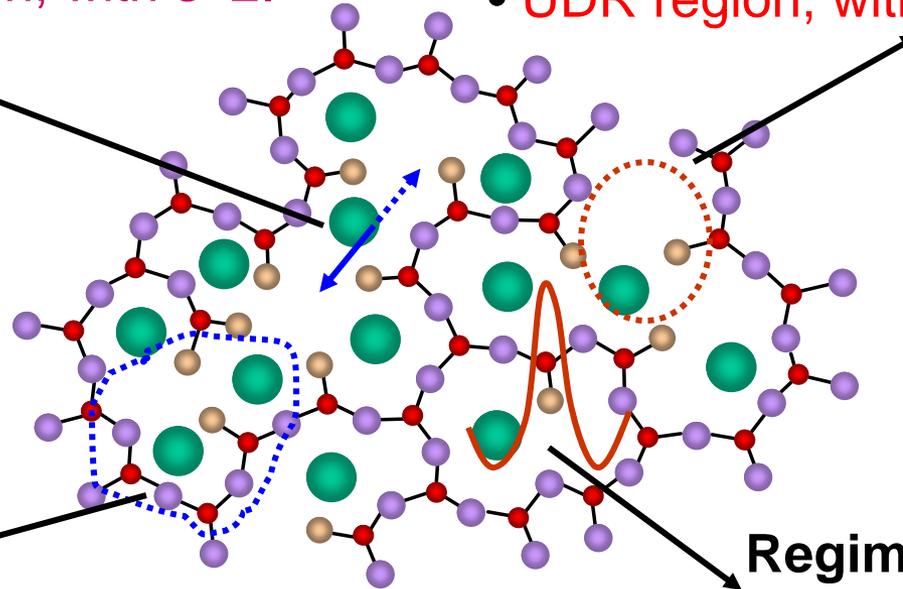
Broad view of the structural origin of conductivity

Regime IV: Very high f

- Vibrational loss region, with $s \approx 2$.

Regime II: High T - Intermediate f

- UDR region, with $s \approx 0.6$.



Regime III: MW f or Low T

- Jellyfish region, with $s \sim 1.0$.

Na  Si  BO  nBO 

Regime I: High T - low f

- DC conductivity region, with $s=0$.

Random network structure of a sodium silicate glass in two-dimension (after Warren and Bischoe)