

Previous Lecture

Coulomb's Law (vector form)

$$\mathbf{F}_{2 \text{ on } 1} = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2 (\mathbf{r}_1 - \mathbf{r}_2)}{|\mathbf{r}_1 - \mathbf{r}_2|^3}$$

Field of a point charge q at origin

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^3} \mathbf{r}$$

Principle of Superposition

Announcement

Beginning next Tuesday, you should pick a seat and sit in it every time. We will take attendance by checking which seats are occupied.

Today

- Continuous charge distributions.
- Charge density
- Integrating to get the electric field
- Field lines

Tricks of the Trade

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- Make approximations

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- Use symmetry

Tricks of the Trade

- Make approximations
- Use symmetry
- Solve problems algebraically

From the Equation Sheet

$$\mathbf{F}_{2 \text{ on } 1} = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2 (\mathbf{r}_1 - \mathbf{r}_2)}{|\mathbf{r}_1 - \mathbf{r}_2|^3}$$

$$\mathbf{F} = q\mathbf{E}$$

$$d\mathbf{E} (\text{at } \mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{dQ (\mathbf{r} - \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3}$$

Limiting Case I

Close to a long wire: $x \ll a \Rightarrow \sqrt{x^2 + a^2} \approx a$

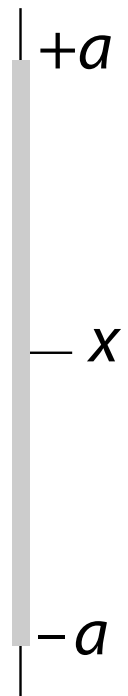
$$\mathbf{E} = \frac{1}{2\pi\epsilon_0} \frac{\lambda}{x} \frac{a}{\sqrt{x^2 + a^2}} \hat{\mathbf{i}} \approx \frac{1}{2\pi\epsilon_0} \frac{\lambda a}{x a} \hat{\mathbf{i}}$$

$$\mathbf{E} \approx \frac{1}{2\pi\epsilon_0} \frac{\lambda}{x} \hat{\mathbf{i}}$$

Near a long wire, \mathbf{E} depends only on the \perp distance to the wire and is directed \perp to the wire. On the equation sheet you will find

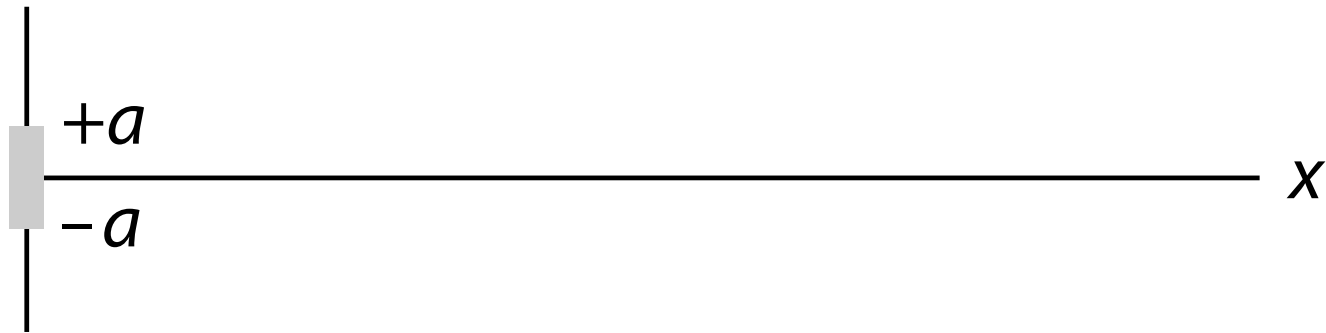
$$E_{\text{line}} = \frac{1}{2\pi\epsilon_0} \frac{\lambda}{r}$$

(You need to know the direction!)



Limiting Case II

Far from a short wire. $x \gg a \Rightarrow \sqrt{x^2 + a^2} \approx x$



$$\mathbf{E} = \frac{1}{2\pi\epsilon_0} \frac{\lambda}{x} \frac{a}{\sqrt{x^2 + a^2}} \hat{\mathbf{i}} \approx \frac{1}{2\pi\epsilon_0} \frac{\lambda a}{x x} \hat{\mathbf{i}} = \frac{1}{4\pi\epsilon_0} \frac{2a\lambda}{x^2} \hat{\mathbf{i}}$$

Since the total charge on the wire is $Q = 2a\lambda$,

$$\mathbf{E} \approx \frac{1}{4\pi\epsilon_0} \frac{Q}{x^2} \hat{\mathbf{i}}.$$

From a distant point, the charge appears to be concentrated at the origin, and the result becomes Coulomb's law for a point charge.

Charge Density

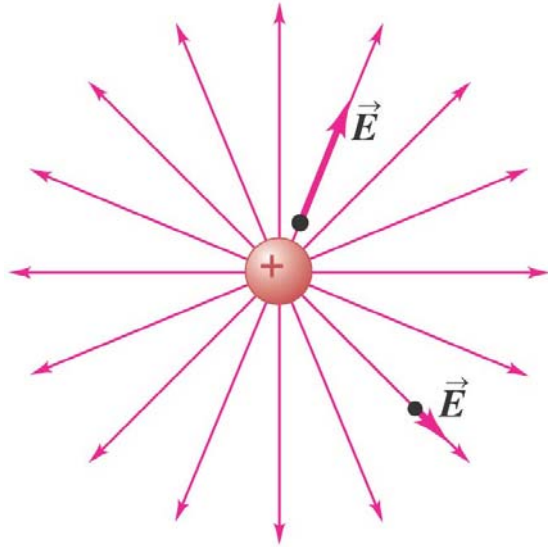
λ $\frac{\text{charge}}{\text{length}}$ **line**

σ $\frac{\text{charge}}{\text{area}}$ **surface**

ρ $\frac{\text{charge}}{\text{volume}}$ **volume**

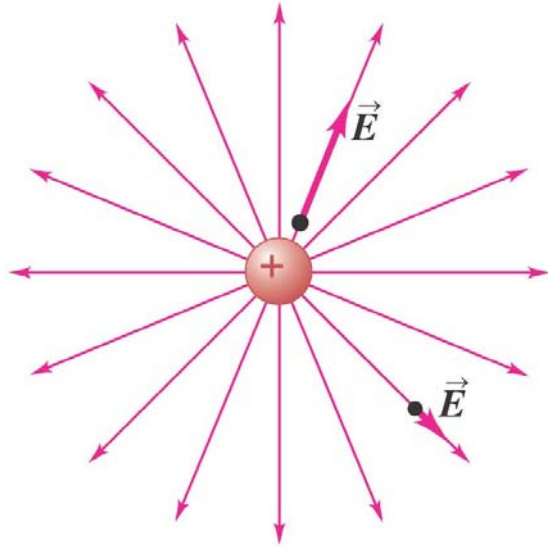
Field Lines

Field lines of a point charge



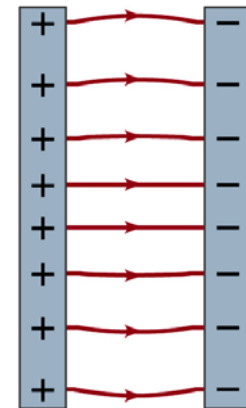
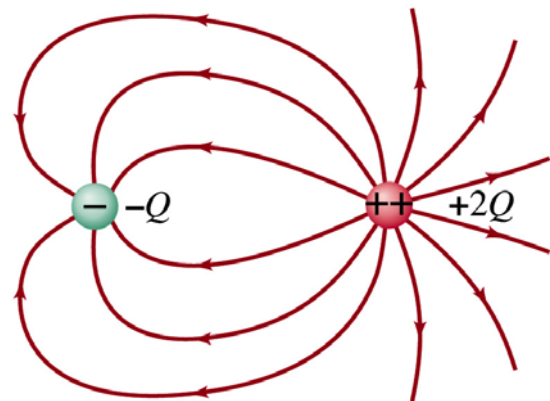
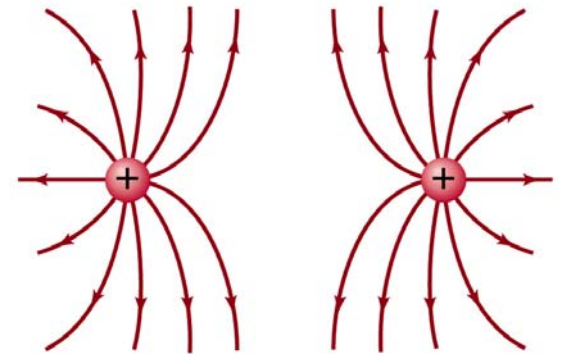
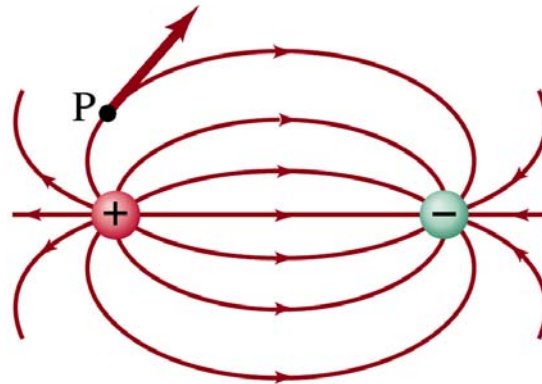
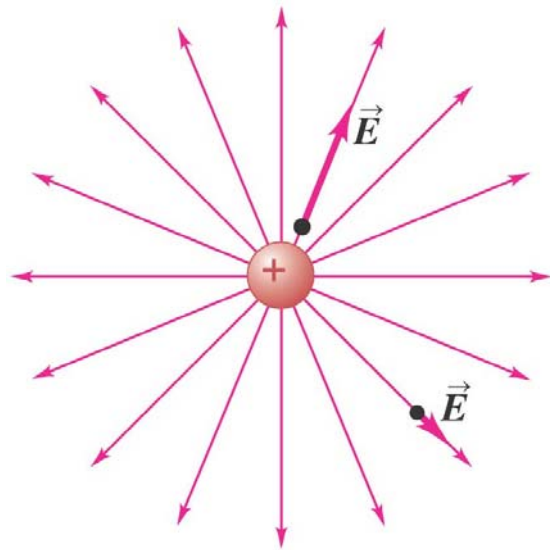
Field Lines

Field lines of a point charge



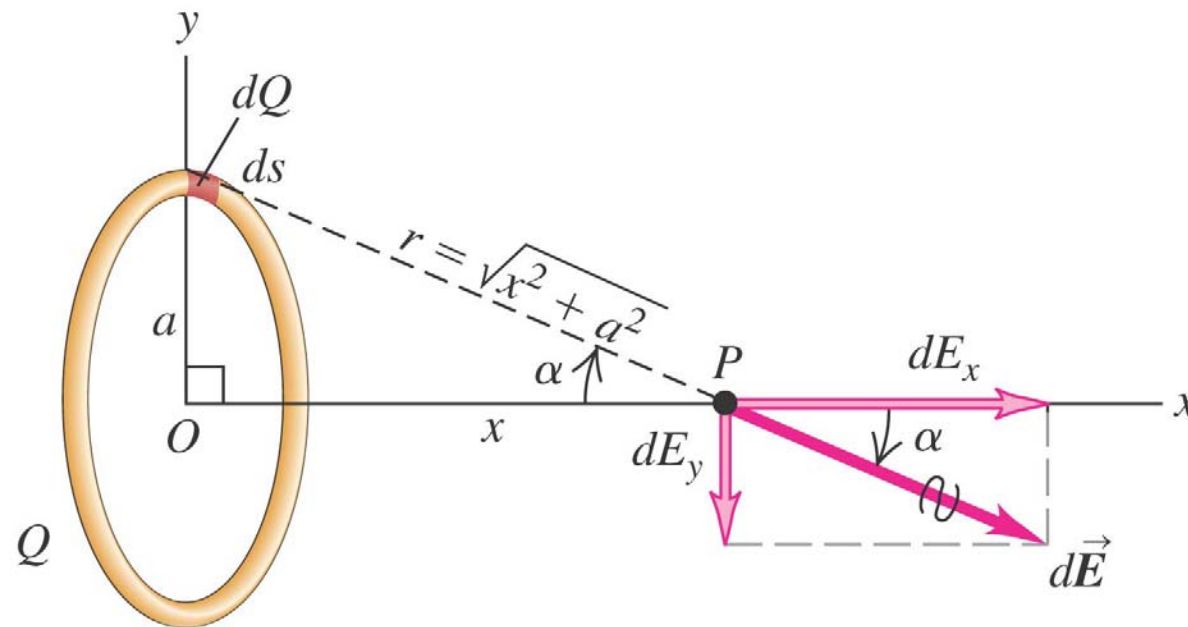
What are the field lines of a long line of charge (uniform density λ)?

Field Lines



Field of a Ring of Charge

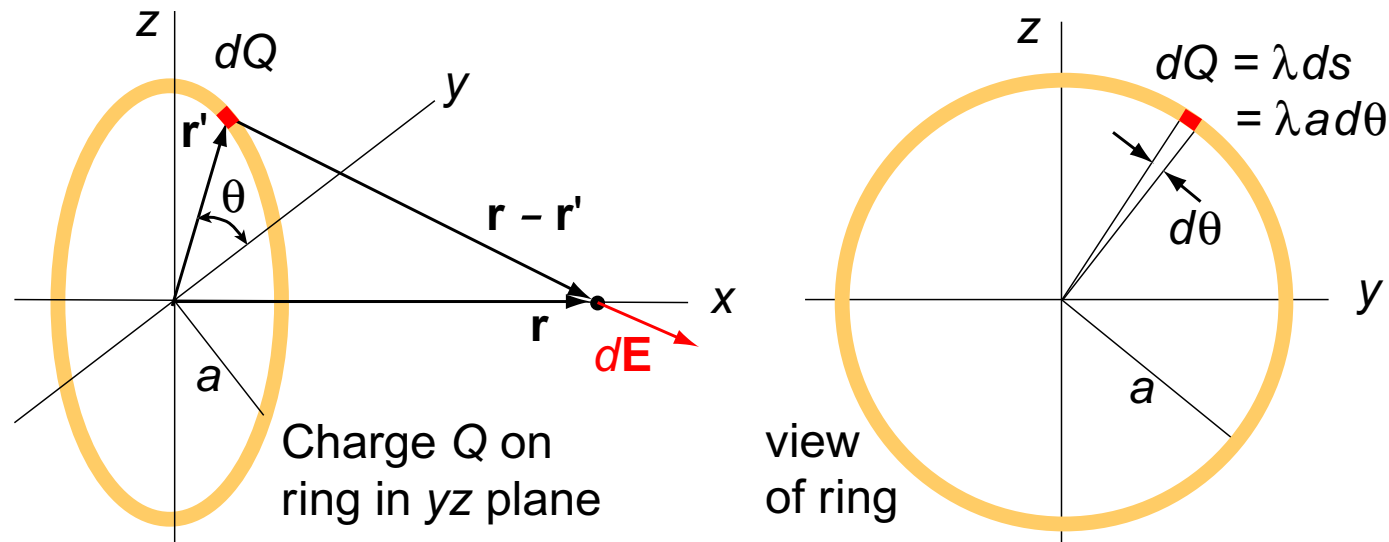
Here's Fig. 21.24 from the text:



We can relate this diagram to our general formula

$$d\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{dQ (\mathbf{r} - \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3}.$$

Field of a Ring of Charge



Using $\lambda = Q/(2\pi a)$, $\mathbf{r} = x \hat{\mathbf{i}}$, and $\mathbf{r}' = a \cos \theta \hat{\mathbf{j}} + a \sin \theta \hat{\mathbf{k}}$,

$$dQ = \lambda a d\theta = \frac{Q d\theta}{2\pi}$$

$$d\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{dQ (\mathbf{r} - \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} = \frac{1}{4\pi\epsilon_0} \frac{Q d\theta}{2\pi} \left(\frac{x \hat{\mathbf{i}} - a \cos \theta \hat{\mathbf{j}} - a \sin \theta \hat{\mathbf{k}}}{(x^2 + a^2)^{3/2}} \right)$$

$$\mathbf{E} = \int_{\theta=0}^{\theta=2\pi} d\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{Qx}{(x^2 + a^2)^{3/2}} \hat{\mathbf{i}} \quad (\hat{\mathbf{j}}, \hat{\mathbf{k}} \text{ components zero})$$

Limiting Case (Field of Ring)

The field, on axis, of a ring of total charge Q is

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{Qx}{(x^2 + a^2)^{3/2}} \hat{\mathbf{i}}$$

What happens if we are a long way from the ring ($x \rightarrow \infty$)?

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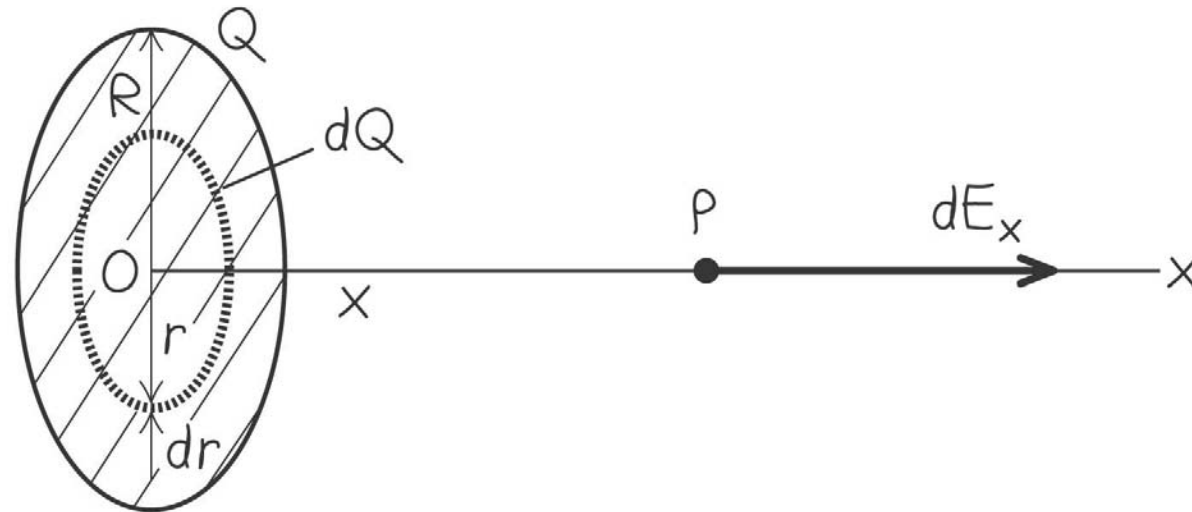
What happens if we are a long way from the ring ($x \rightarrow \infty$)?

All the charge distributed on the ring appears to be concentrated at one point, and we have the field of a point charge Q :

$$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{Q}{x^2} \hat{\mathbf{i}}$$

Field of a Charged Disk

Example 21.12:



We can find the field of a uniformly charged disk by integrating over a series of concentric rings of width dr from $r = 0$ to $r = R$. The result is

$$\mathbf{E} = \frac{\sigma}{2\epsilon_0} \left[1 - \frac{1}{\sqrt{(R^2/x^2) + 1}} \right] \hat{\mathbf{i}}$$