

Special Topics in Relaxation in Glass and Polymers

Lecture 6: Viscoelasticity II Bulk Viscoelasticity

Dr. Ulrich Fotheringham

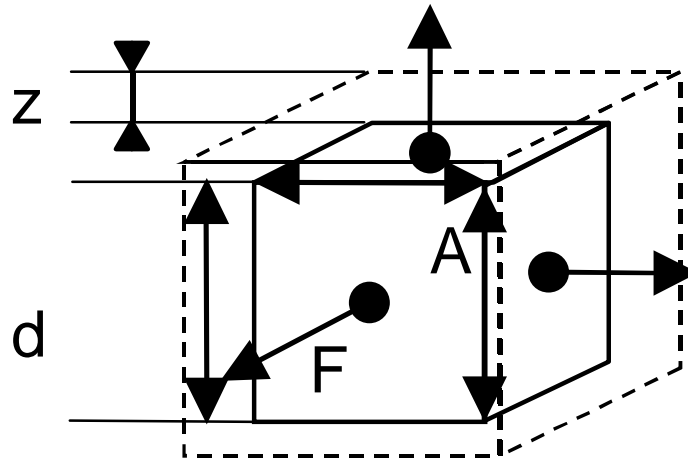
Research and Technology Development

SCHOTT AG

Shear Viscoelasticity and Bulk Viscoelasticity

As it has been said in the previous lecture, one has to distinguish between shear on one side and dilatation/compression (bulk viscoelasticity) on the other.

Bulk



$$F_x = D \cdot x \rightarrow \sigma_x = K \cdot 3 \cdot \varepsilon_x$$

$$F_y = \eta \cdot \frac{a}{d} \cdot \frac{dy}{dt} \rightarrow \sigma_y = \eta_V \cdot \frac{d\varepsilon_y}{dt}$$

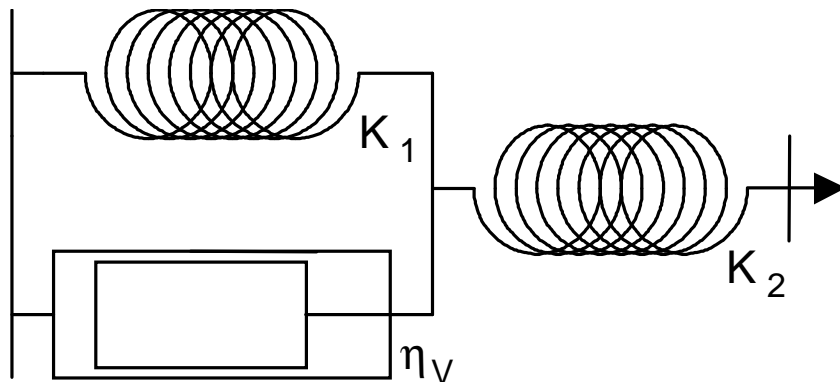
The factor „3“ enters to allow for the fact that the bulk modulus does not relate the stress to the strain in one dimension but to the one in three dimensions, i.e. the volume change.

For the simple Kelvin-Voigt-model one gets assuming constant stress:

$$\varepsilon = \frac{\sigma_0}{3K} \cdot \left(1 - e^{-t \cdot (3K/\eta_V)}\right) = \left(\frac{1}{3K} \cdot \left(1 - e^{-t \cdot (3K/\eta_V)}\right)\right) \cdot \sigma_0 =: J(t) \cdot \sigma_0$$

1st Realistic Picture of Bulk Relaxation: Kelvin-Voigt + Spring

The simple Kelvin-Voigt-model, however, does not contain all features to be observed during the volume change of an inorganic glass being under pressure. It contains the delayed strain occurring at high temperatures, but not the instantaneous effect which is observable at all temperatures. The simplest suited model is the series of a Kelvin-Voigt-model and a spring. (Compared to the Burger-model, the second dashpot is missing which takes into account that in contrast to shear, creep is not possible.)



Compare the damper in a car!

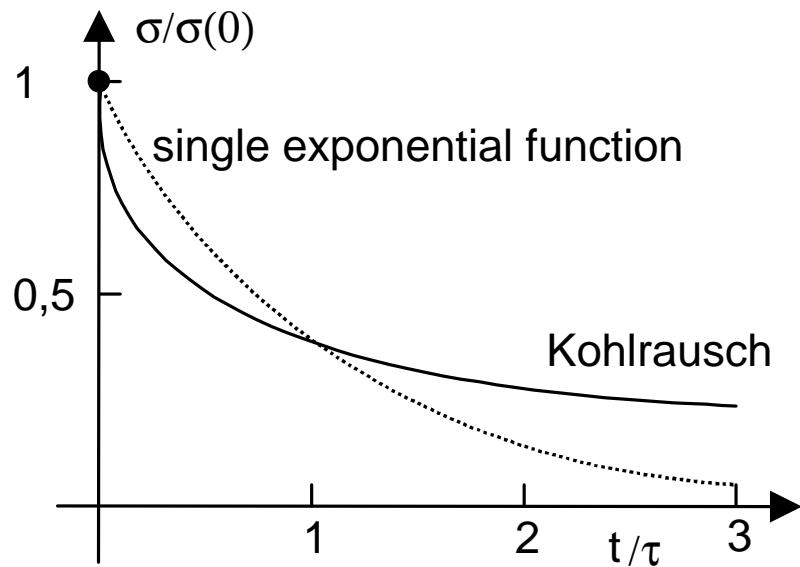
For constant stress, one gets:

$$\varepsilon = \left[\left(\frac{1}{3K_2} \right) + \left(\frac{1}{3K_1} \cdot \left(1 - e^{-t \cdot (3K_1/\eta_V)} \right) \right) \right] \cdot \sigma_0 =: J(t) \cdot \sigma_0$$

For constant strain, one gets:

$$\sigma(t) = \frac{3 \cdot K_1 \cdot K_2}{K_1 + K_2} \cdot \varepsilon_0 + \frac{3 \cdot K_2 \cdot K_2}{K_1 + K_2} \cdot \varepsilon_0 \cdot e^{-t \cdot 3 \cdot (K_1 + K_2) / \eta_V}$$

Kohlrausch-Kinetics for Bulk Relaxation



Again, an even better coincidence of theoretical and experimental data than with the Burger model is obtained, if the single exponential function from above is replaced with a stretched exponential or Kohlrausch(-Williams-Watts)-function.

For constant stress one gets:

$$\varepsilon = \left[\frac{1}{3K_0} + \frac{1}{3} \left(\frac{1}{K_\infty} - \frac{1}{K_0} \right) \cdot \left(1 - e^{-(t/\tau_d)^{b_d}} \right) \right] \cdot \sigma_0 =: J(t) \cdot \sigma_0$$

$J(t)$ is the time-dependent compliance; τ_d and b_d are retardation parameters.

For constant strain one gets

$$\sigma(t) = \left[3 \cdot K_\infty + (3 \cdot K_0 - 3 \cdot K_\infty) \cdot e^{-(t/\tau_x)^{b_x}} \right] \cdot \varepsilon_0 =: K(t) \cdot \varepsilon_0$$

$K(t)$ is the time-dependent bulk modulus; τ_x and b_x are relaxation parameters.

Boltzmann's superposition principle

In general it is assumed that the stress effects arising from strain contributions imposed at different times overlay without interfering and vice versa:

$$\begin{aligned}\sigma(t) &= \sum \Delta\varepsilon(\text{at the time } t') \cdot \left[3 \cdot K_\infty + (3 \cdot K_0 - 3 \cdot K_\infty) \cdot e^{-((t-t')/\tau_x)^{b_x}} \right] \\ &\approx 3 \cdot K_\infty \cdot \varepsilon(t) + (3 \cdot K_0 - 3 \cdot K_\infty) \cdot \int_0^t e^{-((t-t')/\tau_x)^{b_x}} \cdot \frac{d\varepsilon(t')}{dt'} \cdot dt'\end{aligned}$$

One may write introducing relaxation function Ψ

$$\sigma(t) = 3 \cdot K_\infty \cdot \varepsilon(t) + (3 \cdot K_0 - 3 \cdot K_\infty) \cdot \int_0^t \Psi\left(\frac{t-t'}{\tau_x}\right) \cdot \frac{d\varepsilon(t')}{dt'} \cdot dt'$$

and retardation function Φ

$$\varepsilon(t) = \frac{1}{3K_\infty} \cdot \sigma(t) - \frac{1}{3} \left(\frac{1}{K_\infty} - \frac{1}{K_0} \right) \cdot \int_0^t \Phi\left(\frac{t-t'}{\tau_d}\right) \cdot \frac{d\sigma(t')}{dt'} \cdot dt'$$

Relaxation Retardation

Laplace-Transform yields: $L(\sigma) = 3 \cdot K_\infty \cdot L(\varepsilon) + (3 \cdot K_0 - 3 \cdot K_\infty) \cdot L(\Psi) \cdot s \cdot L(\varepsilon)$

and $L(\varepsilon) = \frac{1}{3K_\infty} \cdot L(\sigma) - \frac{1}{3} \left(\frac{1}{K_\infty} - \frac{1}{K_0} \right) L(\Phi) \cdot s \cdot L(\sigma)$

This allows the mutual conversion of relaxation and retardation parameters.

Prony Series

For computational purposes, the Kohlrausch-function may again be represented by a number of single exponentials (Prony-series):

$$e^{-(t/\tau)^b} = \sum_i v_i \cdot e^{-t/\tau_i} \quad , \quad \sum_i v_i = 1$$

Temperature dependence of the relaxation and retardation times

Again, an Arrhenius ansatz is made: $\tau = \tau_0 \cdot e^{\frac{H}{R \cdot T}}$

Some exercises:

1. Consider the combination of a Kelvin-Voigt model and a spring as 1st realistic representation of bulk viscoelasticity. Consider constant strain. What is $\sigma(0)$?
2. Consider the relaxation function in case of single-step strain, $\varepsilon(t) = \varepsilon_0 \cdot \Theta(t)$, Θ : Heaviside function. Calculate $\sigma(t)$ from the below formula (with the script of Prof. Cox, it is very simple!)

$$\sigma(t) = 3 \cdot K_{\infty} \cdot \varepsilon(t) + (3 \cdot K_0 - 3 \cdot K_{\infty}) \cdot \int_0^t \Psi\left(\frac{t-t'}{\tau_x}\right) \cdot \frac{d\varepsilon(t')}{dt'} \cdot dt'$$

3. Consider the Arrhenius law. With $H = 8380$ J/mole and the gas constant $R = 8.38$ J/(mole K), which temperature step is necessary to double relaxation time τ (start temperature: 600°C)?

Bulk relaxation experiments by University of Erlangen-Nürnberg, Clemson-University, Bayrisches Geoinstitut, and SCHOTT AG I



Measurement of the Dimensional Relaxation Effect in Optical Glasses

diploma thesis

Submitted by: Christian Bienert
 Major: materials science
 Referent: Prof. Dr. Rudolf Weißmann
 Advisors: Dr. Ulrich Fotheringham
 Prof. Kathleen Richardson
 Time frame: from October, 3rd 2005
 to July, 3rd 2006 (9 months)



Friedrich-Alexander University Erlangen-Nuremberg
 Department of Materials Science III Glass and Ceramics
 Martensstr. 5, 91058 Erlangen, ww3@ww.uni-erlangen.de

DGG, FA I

Würzburg, 17.03.2009

Optische und volumetrische Messungen an Gigapascal-verdichteten optischen Gläsern

U. Fotheringham¹, O. Sohr¹, P. Fischer¹, G. Westenberger¹, D. Frost²,
 C. Bienert³, R. Weißmann³, K. Richardson⁴

¹SCHOTT AG, Mainz, ²Bayrisches Geoinstitut, Universität Bayreuth, ³Universität Erlangen-Nürnberg, ⁴Clemson University

SCHOTT
 glass made of ideas

Bulk relaxation experiments by Erlangen-Nürnberg, Clemson, Bayrisches Geoinstitut, SCHOTT II

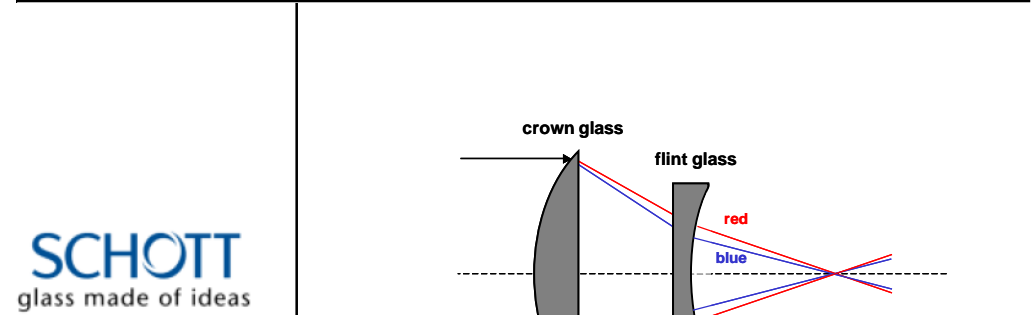
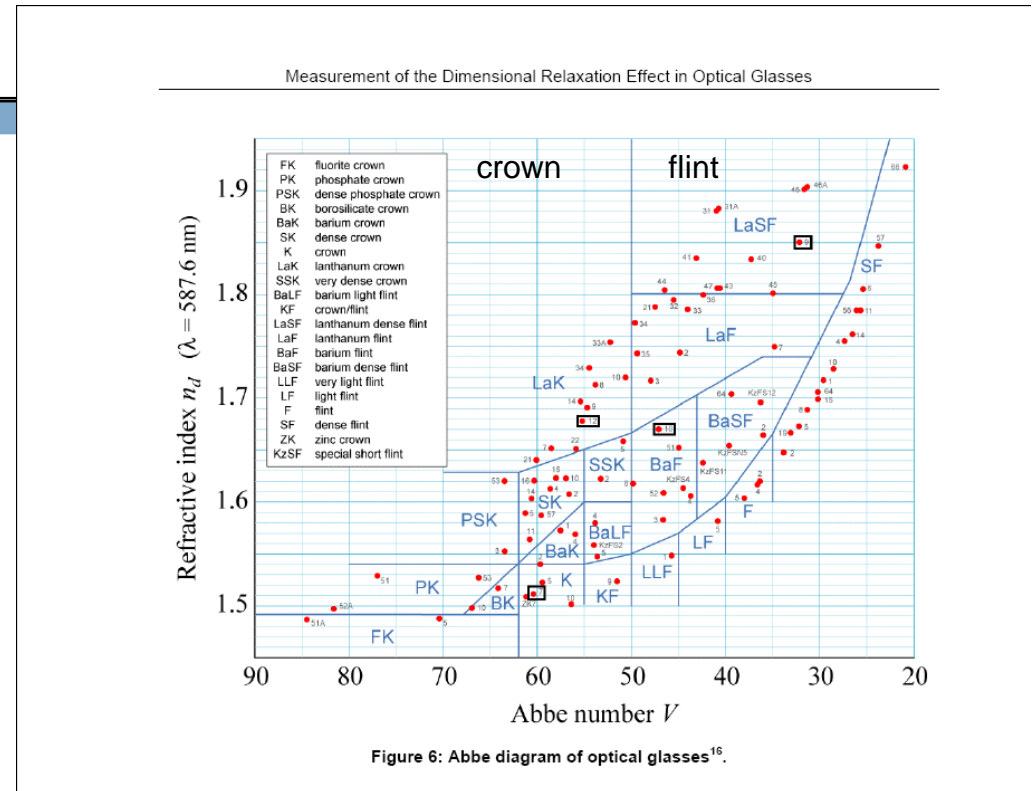
Glasses

DGG, FA I

Würzburg, 17.03.2009

Gläser II: Thermomechanische Eigenschaften, Zusammensetzung

Sample	K 7	N-BaF 10	N-LaK 12	N-LaSF 9
ρ [g/cm ³]	2.53	3.75	4.10	4.41
T_g [°C]	513	660	614	683
$T_{10 \times 13,0}$ [°C]	528	652	615	700
Aluminum Oxide		1-10		
Antimony Trioxide		< 1	< 1	< 1
Arsenic Trioxide	< 1			< 1
Barium Oxide		40-50	40-50	20-30
Boron Oxide	1-10	1-10	10-20	1-10
Lanthanum Oxide			10-20	20-30
Lead Oxide	< 1			
Niobium Pentoxide				1-10
Potassium Oxide	10-20			
Silica	60-70	30-40	10-20	10-20
Sodium Oxide	1-10			1-10
Titanium Oxide	< 1	1-10	< 1	10-20
Zinc Oxide	1-10	1-10		1-10
Zirconium Oxide		1-10	< 1	1-10



SCHOTT
glass made of ideas

Bulk relaxation experiments by Erlangen-Nürnberg, Clemson, Bayrisches Geoinstitut, SCHOTT III

Multi-Anvil-Device experiments:

Pressure in the Gigapascal range

DGG, FA I

Würzburg, 17.03.2009

Multi-Anvil-Device I: Aufbau

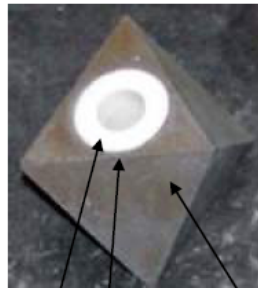


3,8mm
3,8mm



Samples are wrapped in Iridium foil

Einpacken in Ir-Folie



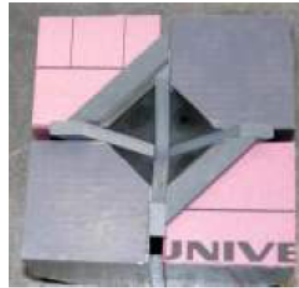
ZrO₂

MgO-Oktaeder mit Heizelement

Position des Thermoelementes

Thermocouple

Heating element in the octa-eder



WC-Kuben; eine Ecke abgesägt; durch dickes Papier getrennt

Tungsten carbide cubes with one corner cut off



Hydraulische Presse, z.B. Sumitomo 1200 (400 bar Öldruck entsprechen 5 GPa)

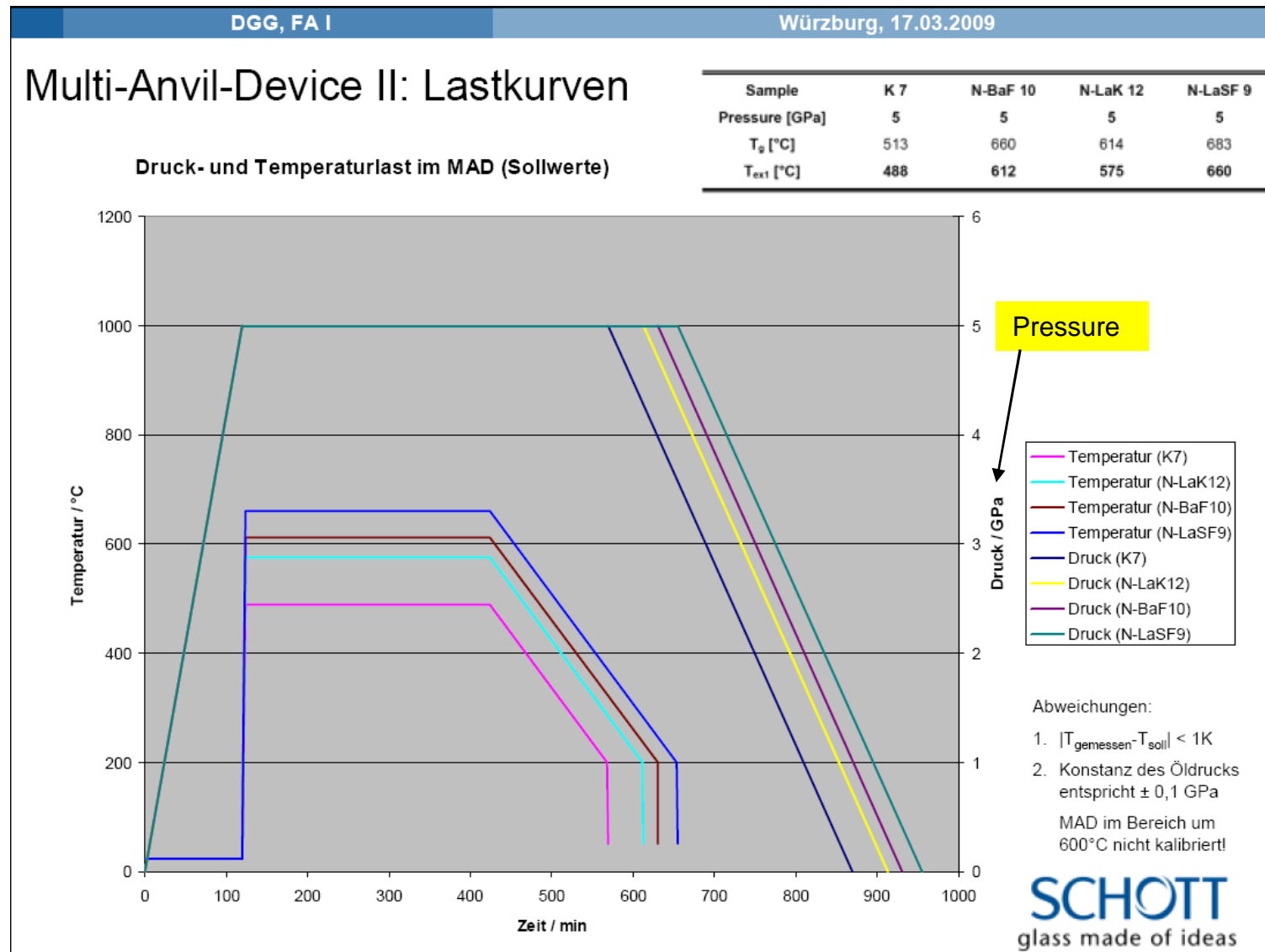
Hydraulic press

SCHOTT
glass made of ideas

Bulk relaxation experiments by Erlangen-Nürnberg, Clemson, Bayrisches Geoinstitut, SCHOTT IV

Multi-Anvil-
Device
experiments:

Temperature
and
pressure
loads



Bulk relaxation experiments by Erlangen-Nürnberg, Clemson, Bayrisches Geoinstitut, SCHOTT V

Multi-Anvil-
Device
experiments:

Density
measurements:

when and how

DGG, FA I

Würzburg, 17.03.2009

Dichte I: Messungen

wann:

nach Probenpräparation (5×)



1st measurement after
sample preparation

nach Verdichtung (5×)

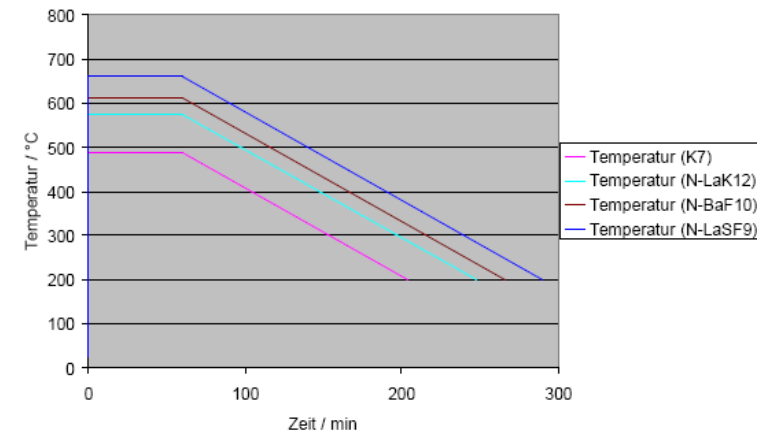


2nd measurement after
densification

3rd measurement after relaxation
in forced convection furnace

nach Entspannungsintervall im Umluftofen (10×)

Temperaturführung bei Entspannung (Beispieldauer 1h, Sollwerte)



wie:

Messung nach Archimedes mit wäßriger Lösung von Nekal BX trocken (1g/l; BASF; 0.9977190 g/cm^3 bei 25° ; „Nekal“ kommt von „netzt kalt“ und ist ein Meilenstein der Tensidentwicklung)

Density measurement after Archimedes in
Nekal solution (Nekal: tenside from BASF)

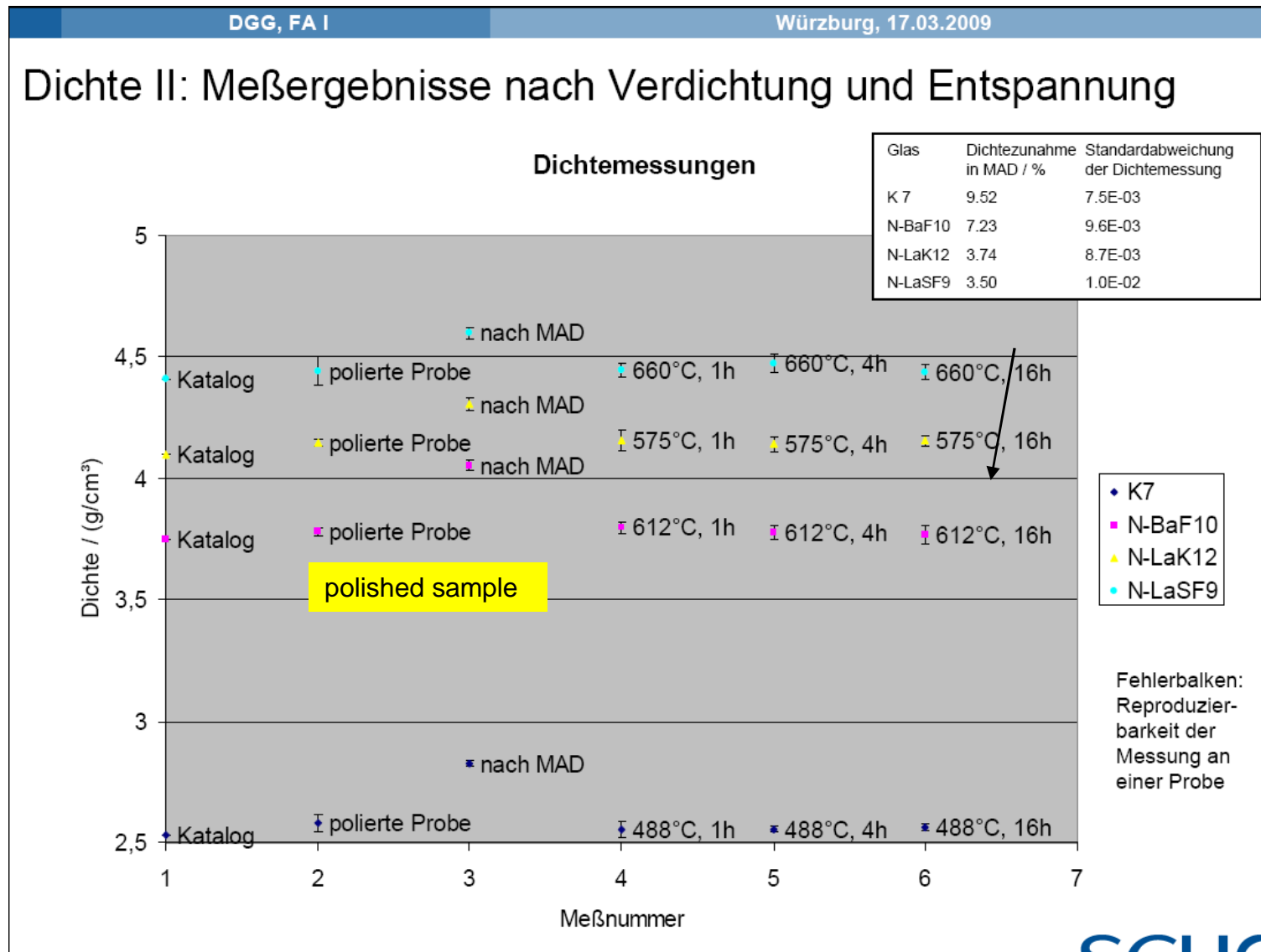
SCHOTT

Bulk relaxation experiments by Erlangen-Nürnberg, Clemson, Bayrisches Geoinstitut, SCHOTT VI

Multi-Anvil-
Device
experiments:

Density
measurements:

results



Bulk relaxation experiments by Erlangen-Nürnberg, Clemson, Bayrisches Geoinstitut, SCHOTT VII

Hot Isostatic
Press
experiments:

Samples and
device

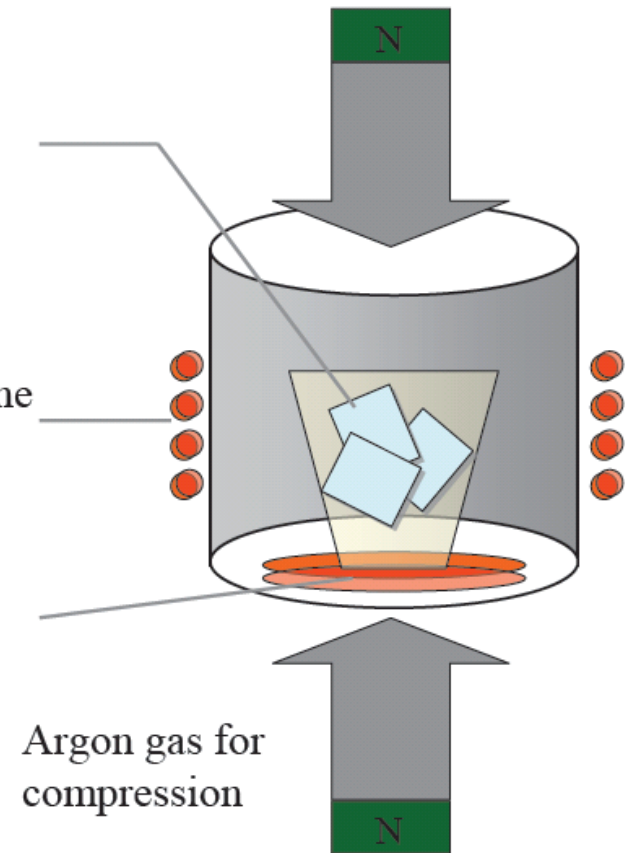


polished sample

thermocouple in
compression
chamber

thermocouple at the
lateral heater

thermocouple at
the bottom heater



Experimental conditions for the sample glasses in HIP

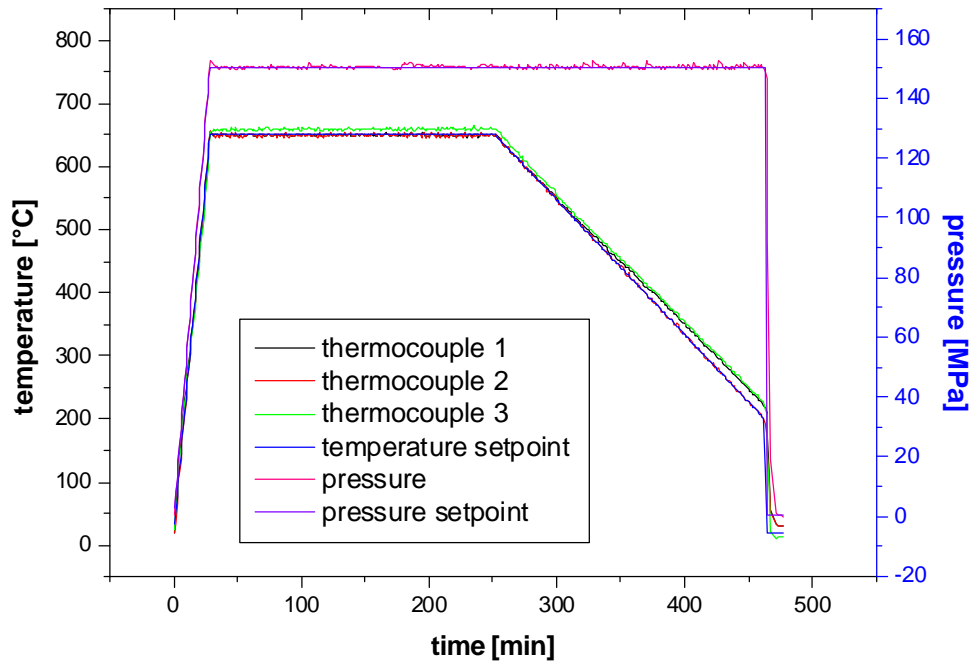
Sample	K 7	N-BaF 10	N-LaK 12	N-LaSF 9
Pressure [MPa]	150	150	150	150
T_g [°C]	513	660	614	683
T_{ex} [°C]	488	612	575	660

Bulk relaxation experiments by Erlangen-Nürnberg, Clemson, Bayrisches Geoinstitut, SCHOTT VIII

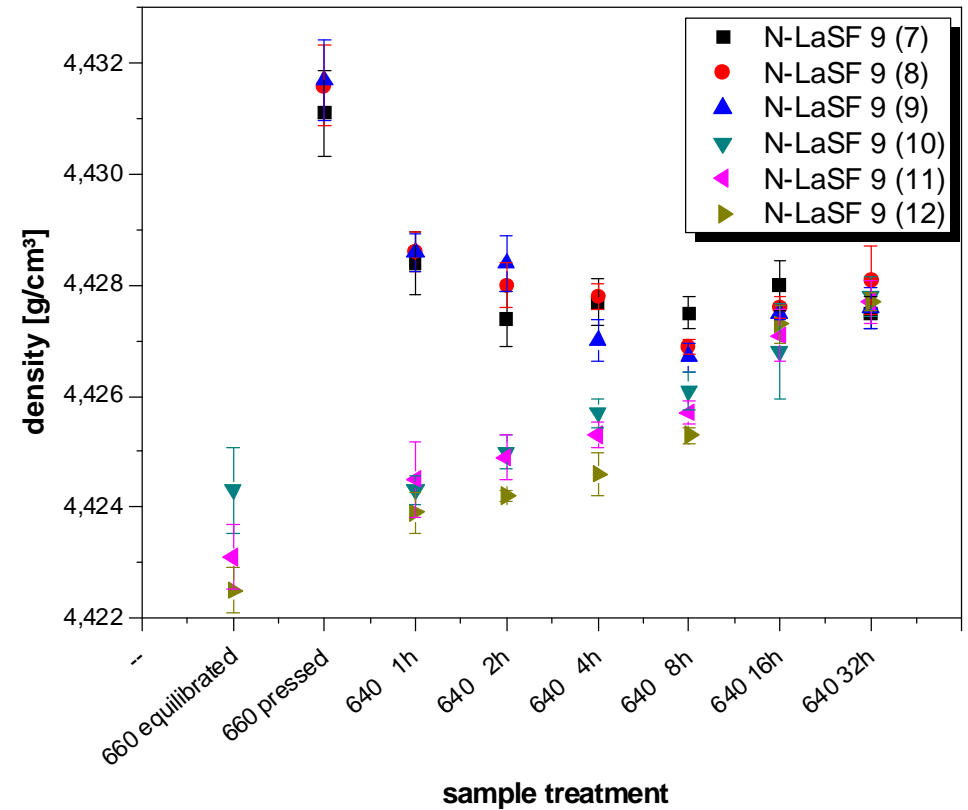
Hot Isostatic Press experiments:

- temperature and pressure load for N-LaSF9

- time-dependent compliance of N-LaSF9



Recorded data of HIP compression experiment for N-LaSF 9.



Density evolution over treatment time with change in fictive temperature T_f of N-LaSF 9 (HIP) (compressed: 7-9; reference: 10-12).

Bulk relaxation experiments by Erlangen-Nürnberg, Clemson, Bayrisches Geoinstitut, SCHOTT IX

Results

Table 10: Table of stand. deviation in [%] of density increase obtained in this study.

Glass	Pressure device	Increase in density [%]	Mean stand. error of density measurements	Mean stand. error in [%] of increase in density
K 7	HIP	0.34	1.4E-04	4
	MAD	9.52	7.5E-03	7.9
N-BaF 10	HIP	0.30	2.4E-04	7.9
	MAD	7.23	9.6E-03	13
N-LaK 12	HIP	0.16	2.5E-04	15
	MAD	3.74	8.7E-03	23
N-LaSF 9	HIP	0.26	2.3E-04	8.9
	MAD	3.50	1.0E-02	29

Table 9: Comprehensive table of bulk moduli and compliances obtained in this study.

	K 7		N-BaF 10		N-LaK 12		N-LaSF 9	
	HIP	MAD	HIP	MAD	HIP	MAD	HIP	MAD
K_0 [GPa]	40		65		68		85	
K_∞ [GPa]	21	23	28	33	39	45	35	53
K_∞ / K_0	0.53	0.57	0.43	0.52	0.57	0.66	0.41	0.63
$J(0)$ [Pa ⁻¹]	8.3E-12		5.1E-12		4.9E-12		3.9 E-12	
$J(\infty)$ [Pa ⁻¹]	1.6E-11	1.5E-11	1.2E-11	1.0E-11	8.5E-12	7.4E-12	9.6E-12	6.2E-12
$J(\infty) / J(0)$	1.9	1.8	2.3	1.9	1.8	1.5	2.5	1.6
$3J(0)$ [Pa ⁻¹]	2.5E-11		1.5E-11		1.5E-11		1.2E-11	
$3J(\infty)$ [Pa ⁻¹]	4.7E-11	4.4E-11	3.6E-11	3.0E-11	2.6E-11	2.2E-11	2.9E-11	1.9E-11