

**I.
Fragility and its Relation to Other Glass Properties**

**II.
Networks**

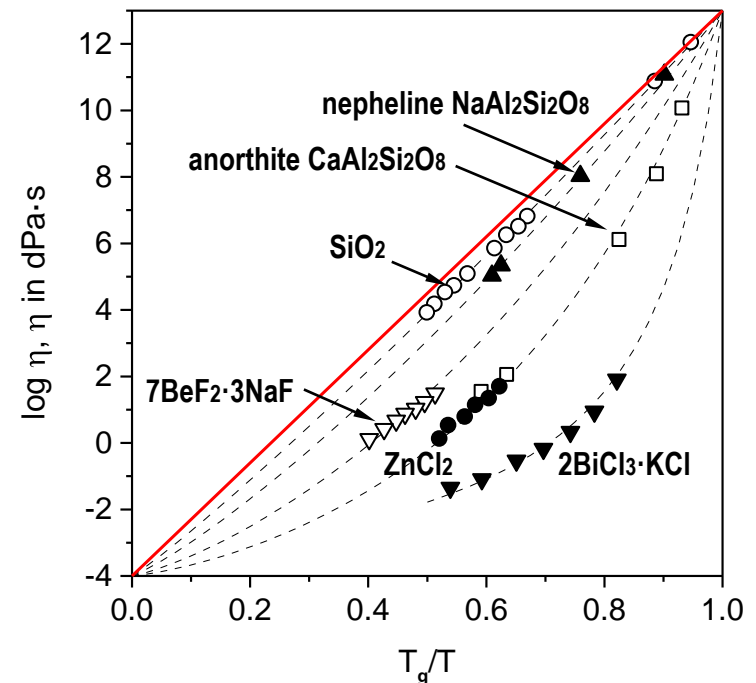
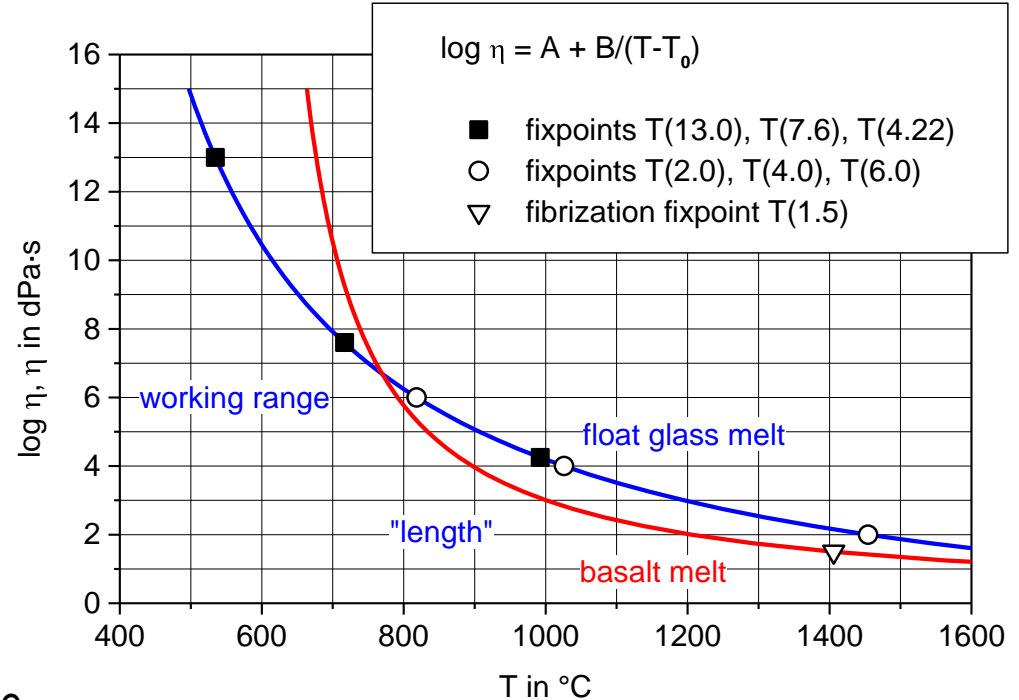
Reinhard Conradt

April 6 & 8, 2010

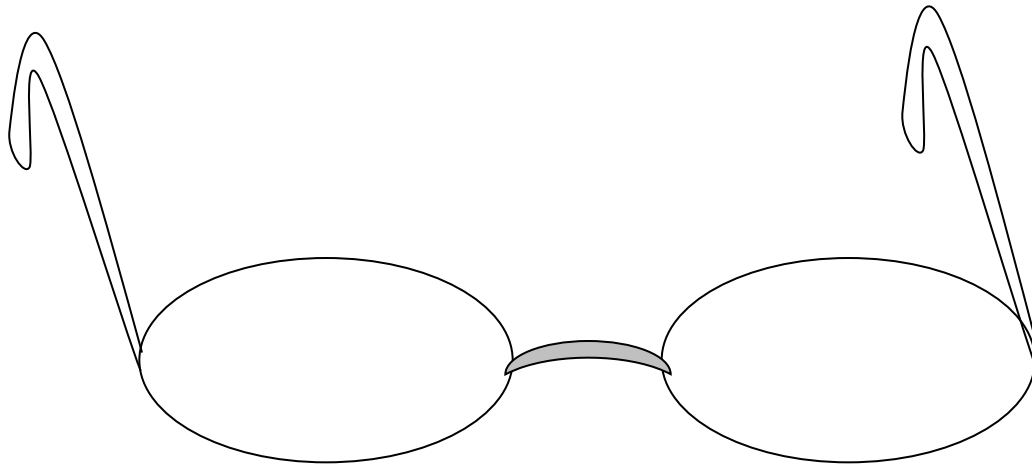
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What is „fragility“?

- Technologists call it the „length“ of a glass; it determines the working range.
- Fragility is a measure of deviation from an exponential T dependence of relaxation times and transport properties (deviation from Arrhenius behavior).
- Fragility may be quantified, e.g., by the slope for $T_g / T \rightarrow 1$ in the so-called Angell plot.
- The non-exponential T dependence is commonly represented by the KWW eq. $\tau \propto \exp B \cdot (T_g/T)^n$, by the VFT eq. $\tau \propto \exp B/(T - T_0)$, etc.
- While the high-T branch reflects a universal feature of liquids, the fragility is likely to reflect the thermodynamics of the liquids.

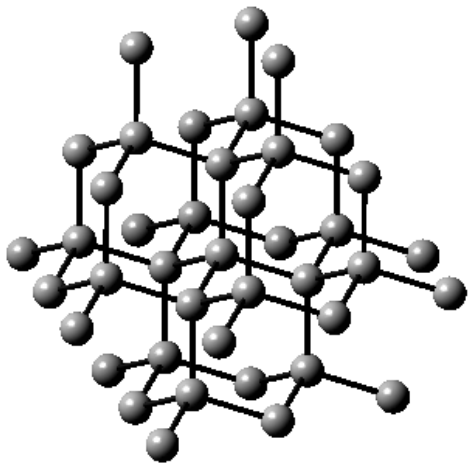


the glassy „state“



structure

dynamics



topology hkl

networks

bonds &
coordination

field strength

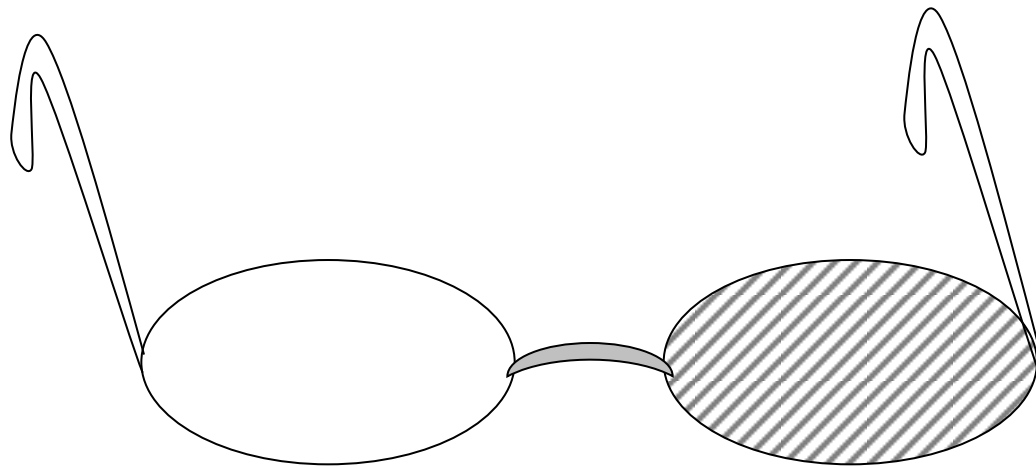
energy kT

$c_p(T)$, $v(T)$

stoichiometry &
constitution

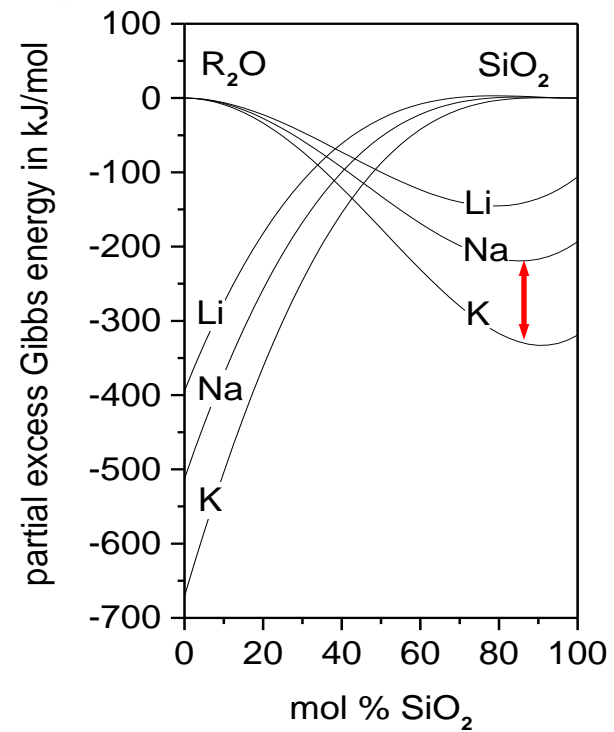
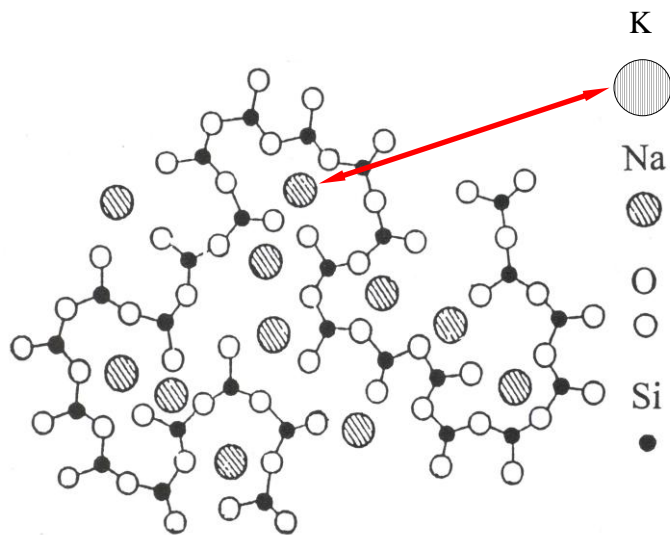
phase diagrams

ΔH , ΔG , K_p , k^+ , k^-



structure

dynamics



matter in internal equilibrium:

state A: $T^A, P^A, n_j^A \rightarrow$ state B: T^B, P^B, n_j^B ; n_j stands for $\vec{n} = \begin{pmatrix} n_{j=1} \\ \vdots \\ n_{j=N} \end{pmatrix}$ (chemical composition)

- a state depends on variables T, P, n_j only; it is „path independent“; let $Z = Z(T, P, n_j)$ be a property (a „function of state“), then

$$\int_A^B dZ = Z(B) - Z(A); \quad \oint dZ = 0$$

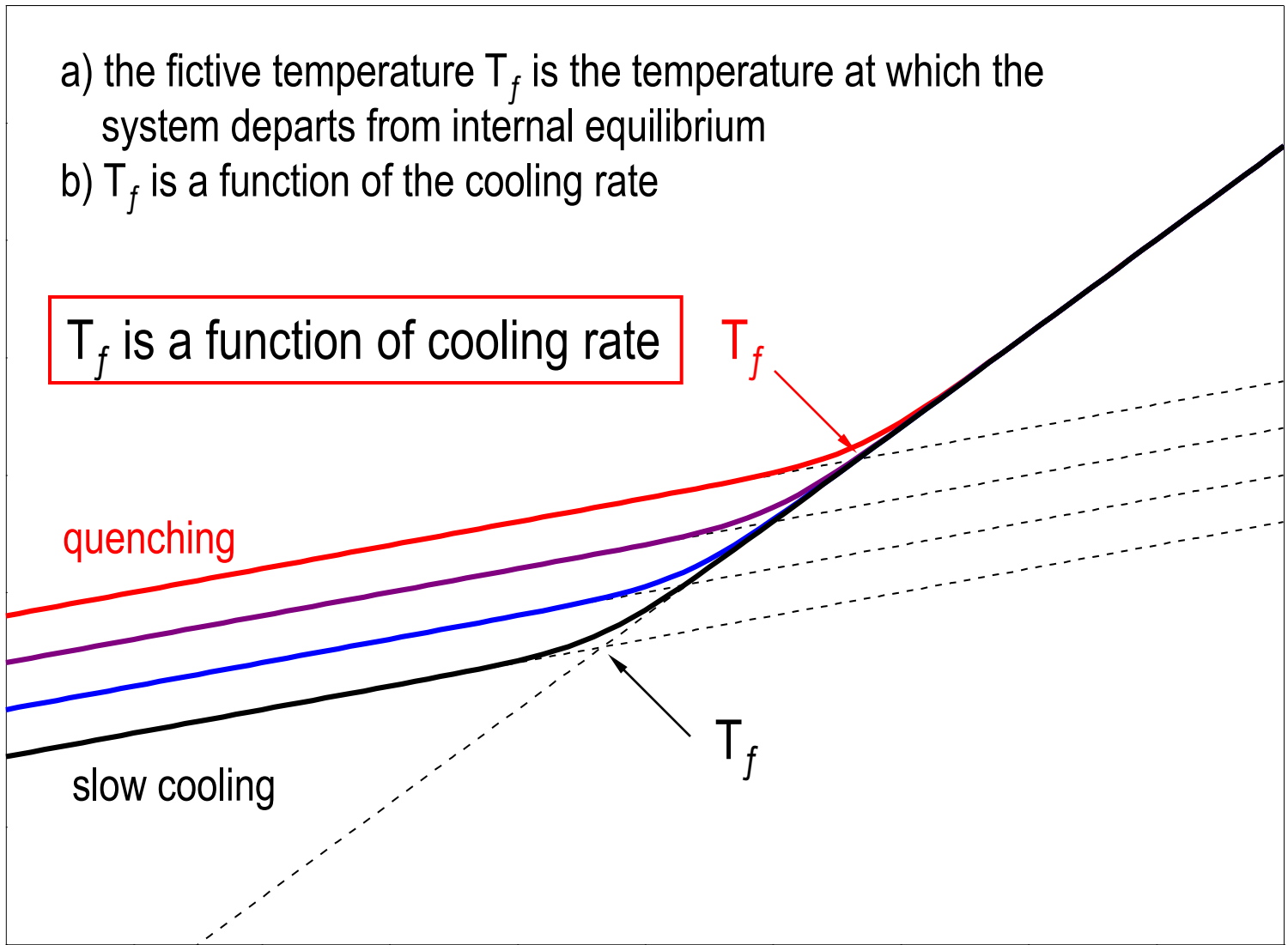
matter at T^B, P^B, n_j^B not in internal equilibrium:

- Z is no longer a function of (T^B, P^B, n_j^B) alone; it is path dependent;
- a path may be defined as
 - $T^A \rightarrow T^B, P = \text{const.}, n_j = \text{const.}$ (e.g., a cooling process)
 - $P^A \rightarrow P^B, T = \text{const.}, n_j = \text{const.}$ (e.g., an impact process)
 - $n_j^A \rightarrow n_j^B, T = \text{const.}, P = \text{const.}$ (e.g., a polymerization process)

or as combinations of T, P, n_j changes;

- the value of T, P, n_j at which the system departs from internal equilibrium is termed fictive temperature T_f , fictive pressure P_f , or fictive composition n_{jf} ;
- Z is a function of $(T^B, T_f^B; P^B, P_f^B; n_j^B, n_{jf}^B)$
- T_f^B, P_f^B, n_{jf}^B depend on the process time

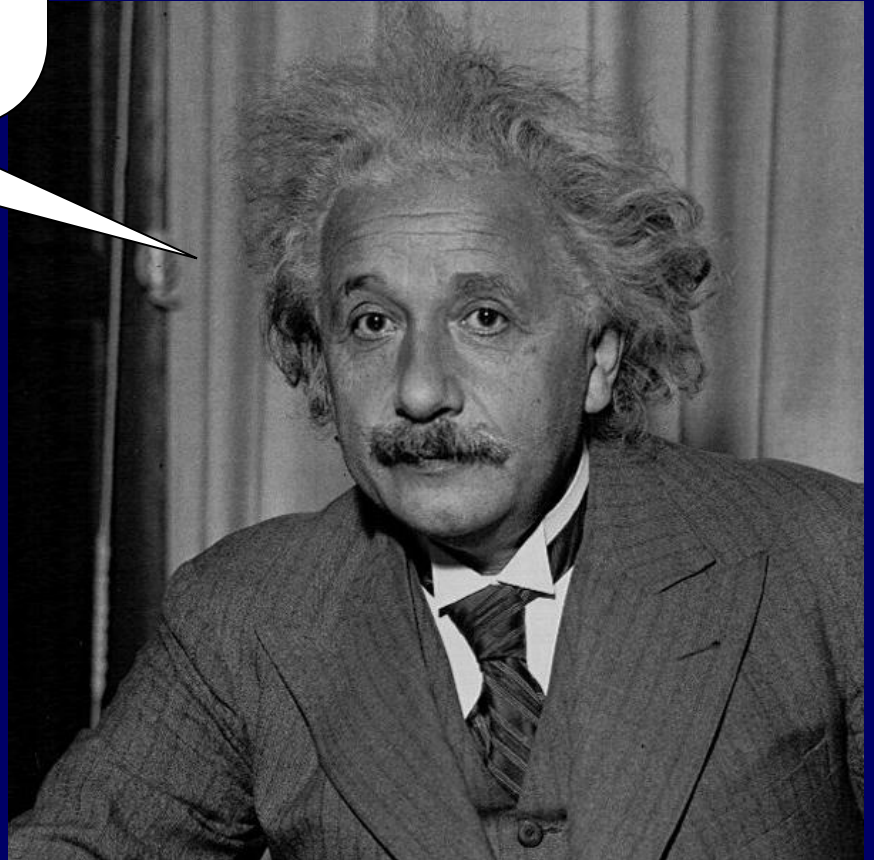
property p (e.g., molar volume V_m or enthalpy H)



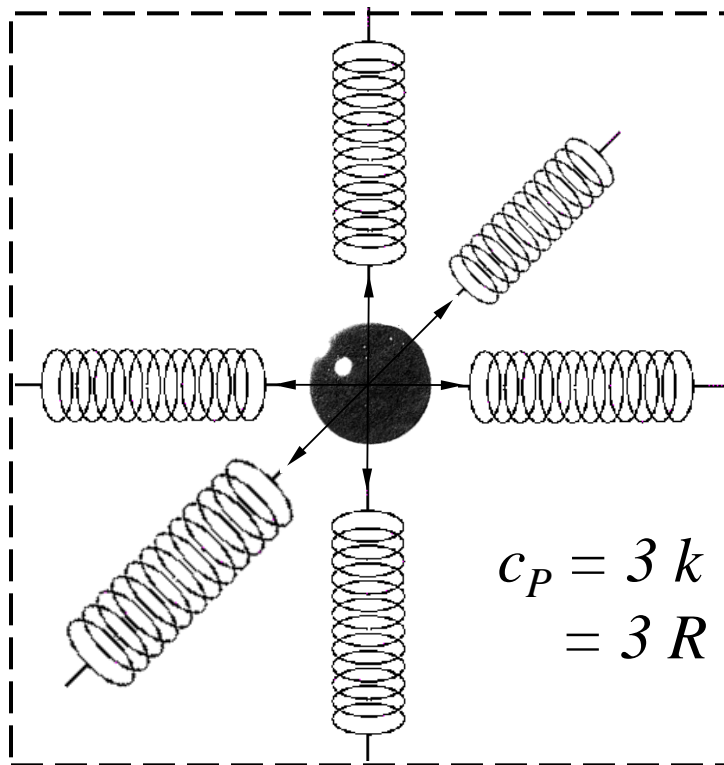
temperature T

heat capacities

**If you have only one
chance to perform a
measurement on a
material,
then measure c_p !**

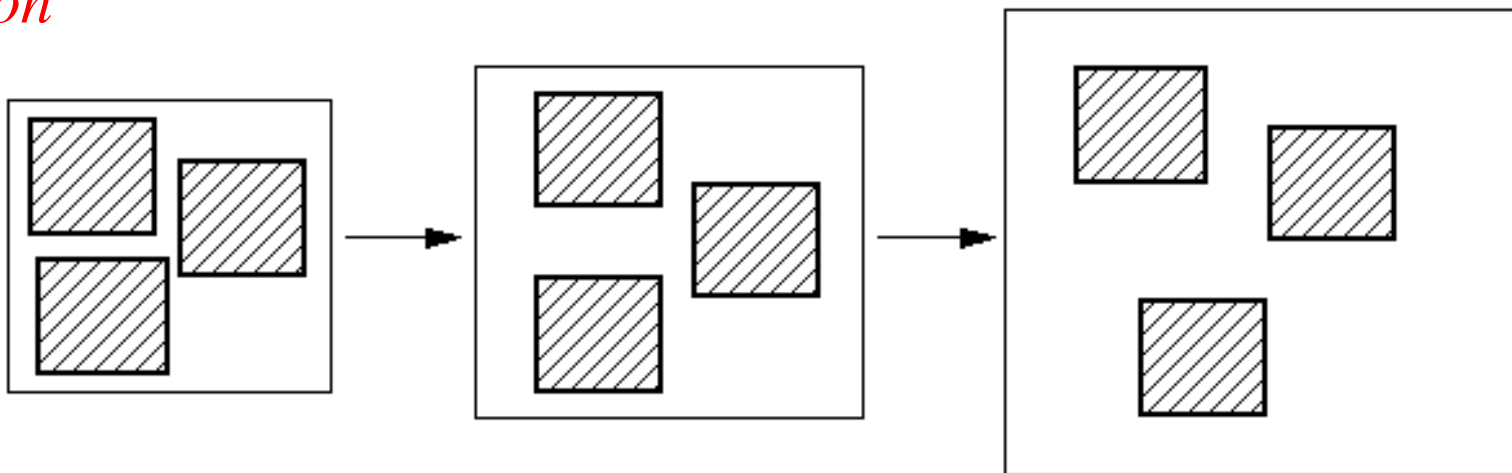


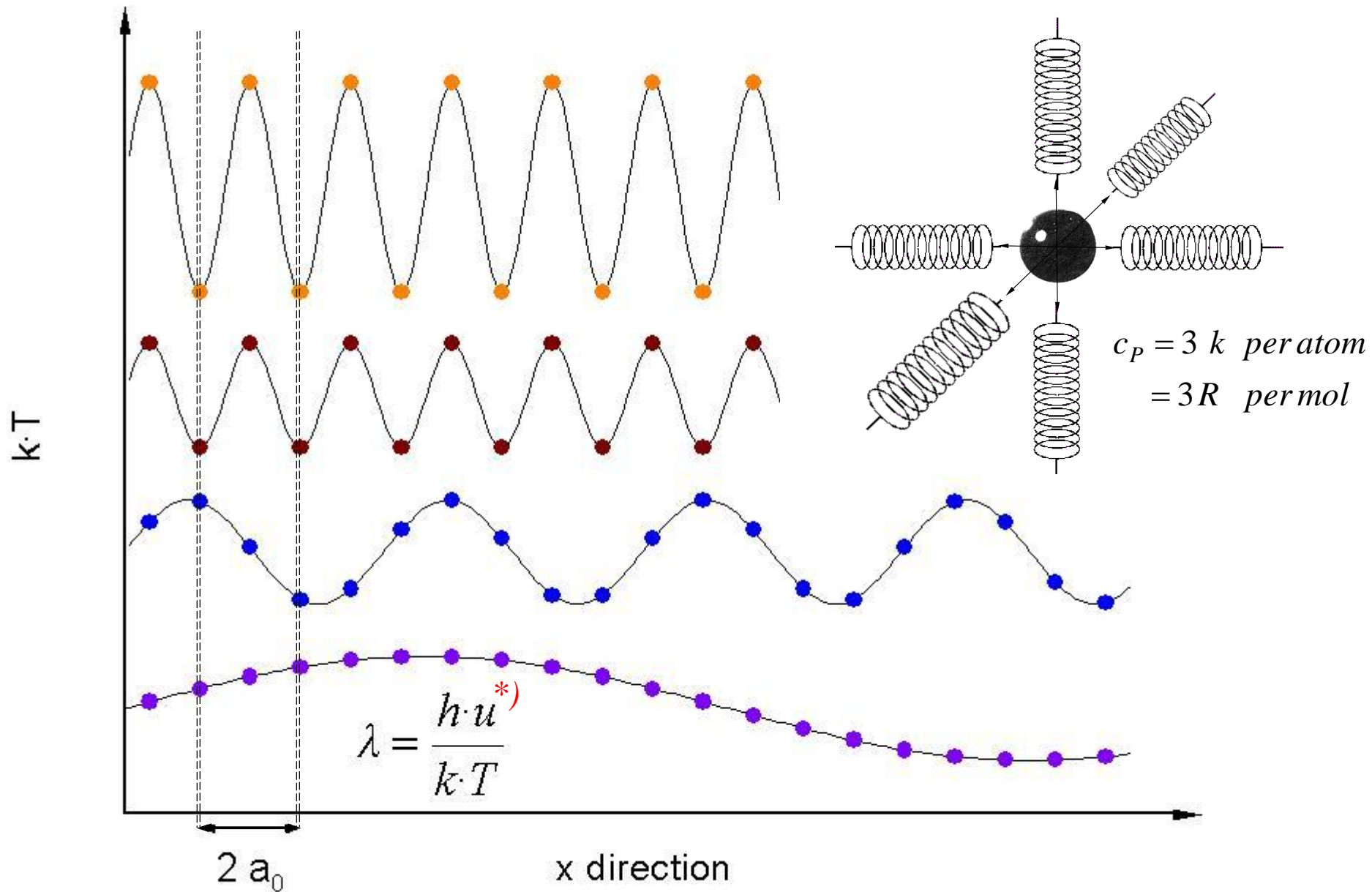
vibration →



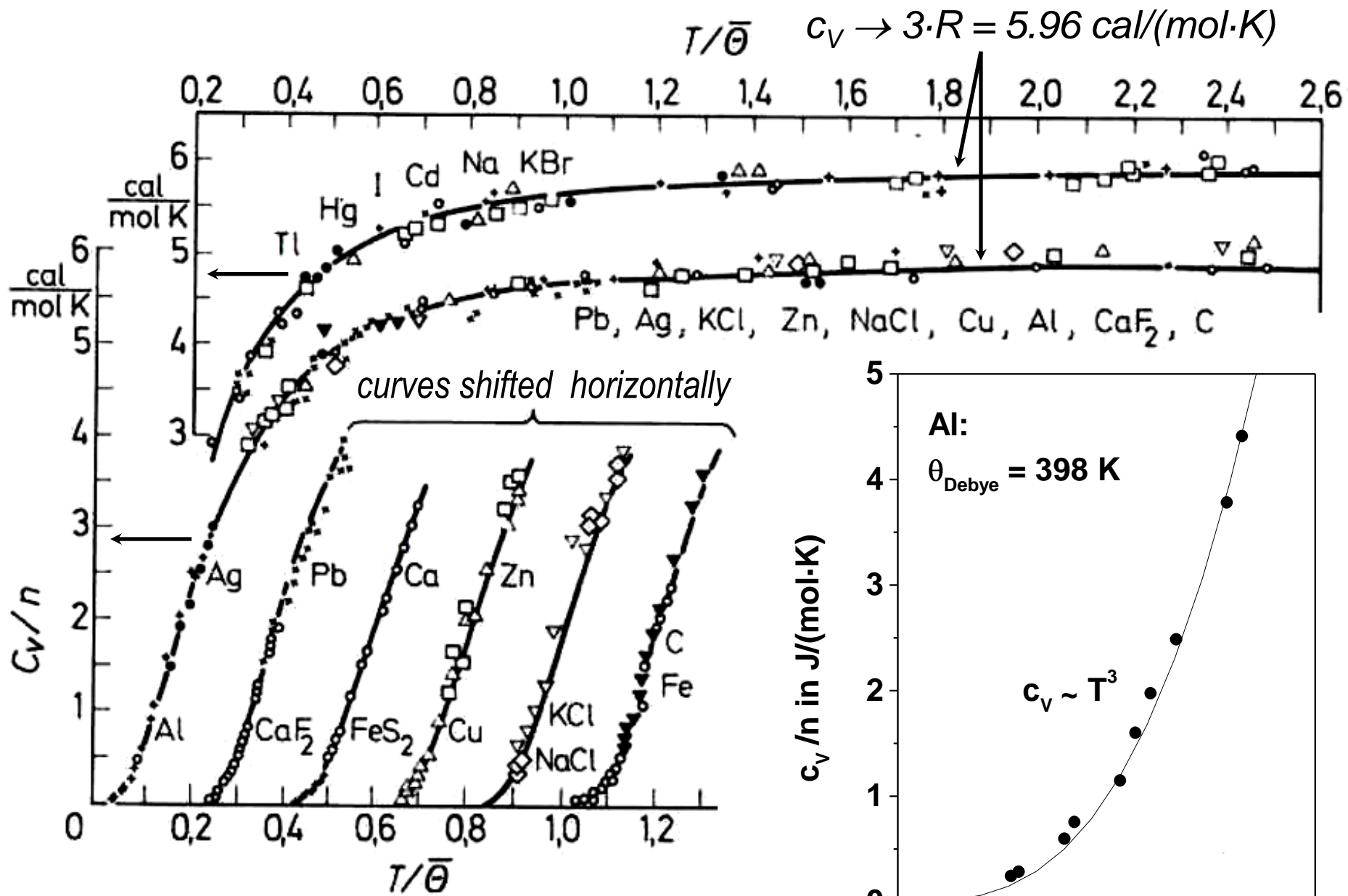
$$c_P = 3k \text{ per atom} \\ = 3R \text{ per mol}$$

configuration

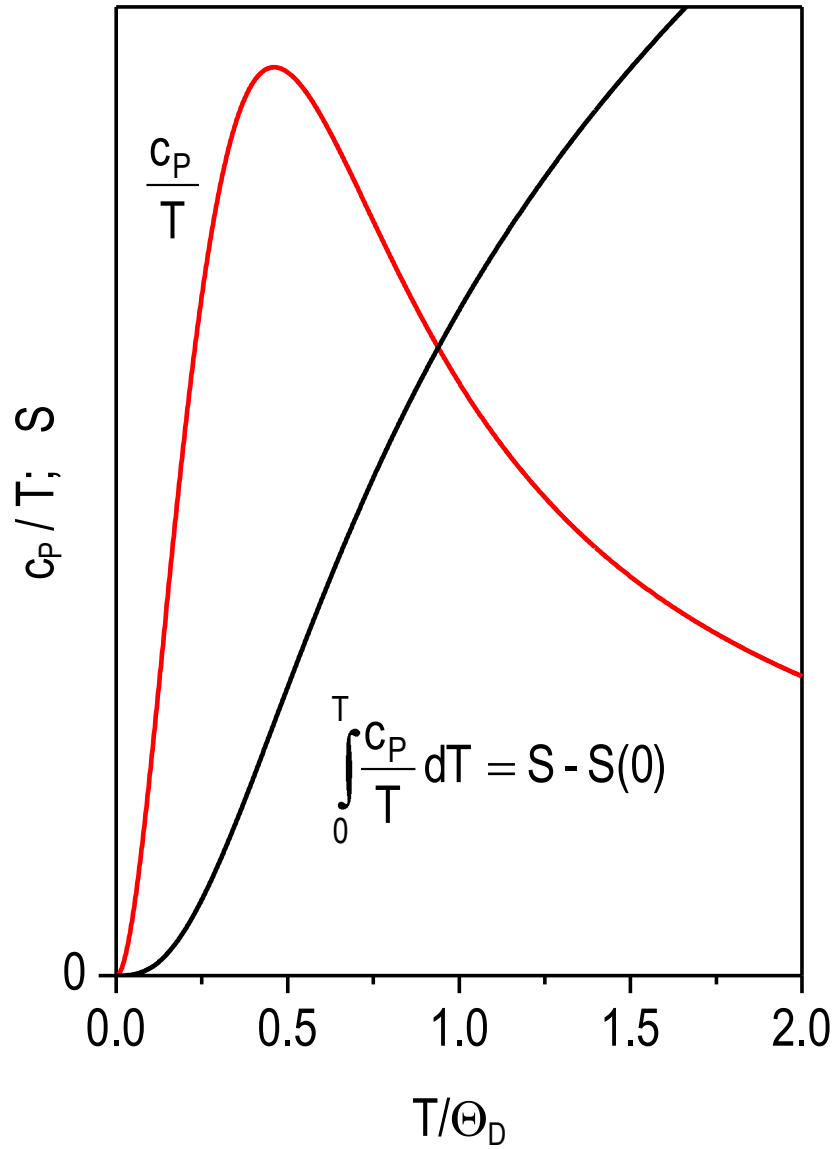
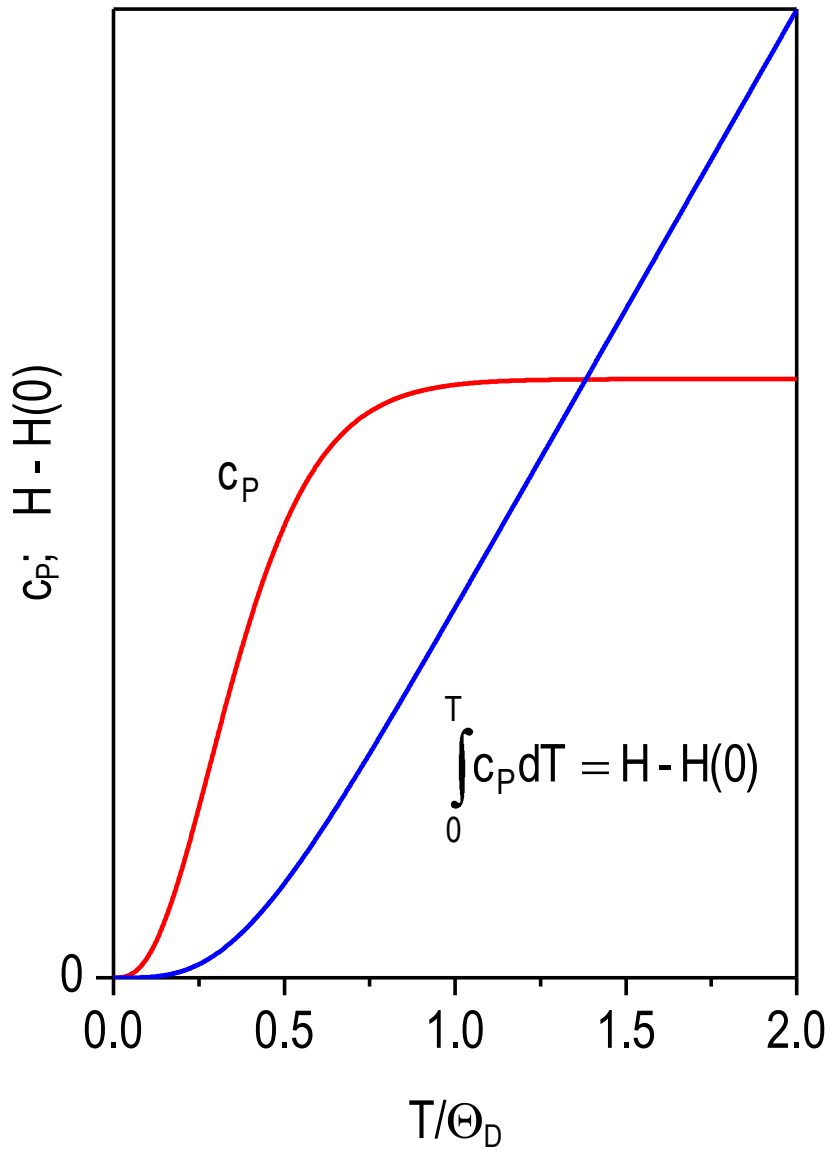




*) $u = \text{sound velocity}$



Hellwege 1976



Example: N₂(g) at boiling point

$$S(T) = \int_0^T \frac{c_p}{T} dT + \sum_i \frac{\Delta H_i}{T_i}, \quad T_i < T$$

T in K	comment	ΔS in J/(mol·K)
0 – 10	Debye integral	1.916
10 – 35.61	integration of c_p/T	25.246
at 35.61	polym. phase transition	6.427
35.61 – 63.14	integration of c_p/T	23.384
at 63.14	melting point T^m	11.418
63.14 – 77.32	integration of c_p/T	11.414
at 77.32	boiling point T^b	72.120
	sum for N ₂ (g) at T^b	151.925
	calc. from partition function	152.236

Giauque & Clayton (1933), compiled by Wilks (1963)

other examples: Ar, O₂, Cl₂, HCl, CH₄, C₂H₄, ...

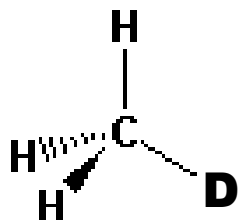
Examples: CO and CH₃D

gas	S(T ^b) in J/(mol·K)		
	S _{calc}	S _{exp}	S _{calc} - S _{exp}
CO at boiling point T ^b	160.33	155.64	4.69
CH ₃ D at boiling point T ^b	165.23	153.64	11.59



a frozen-in random linear orientation would yield

$$\Delta S = -R \ln(1/2) = 5.76 \text{ J/(mol}\cdot\text{K)}$$

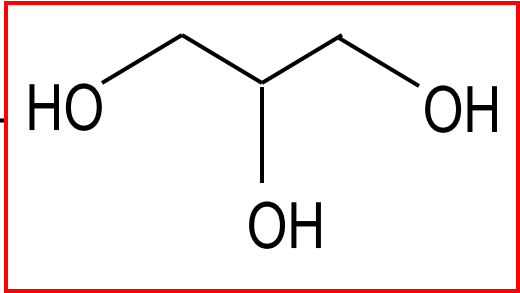


a frozen-in random position of D in the molecule would yield

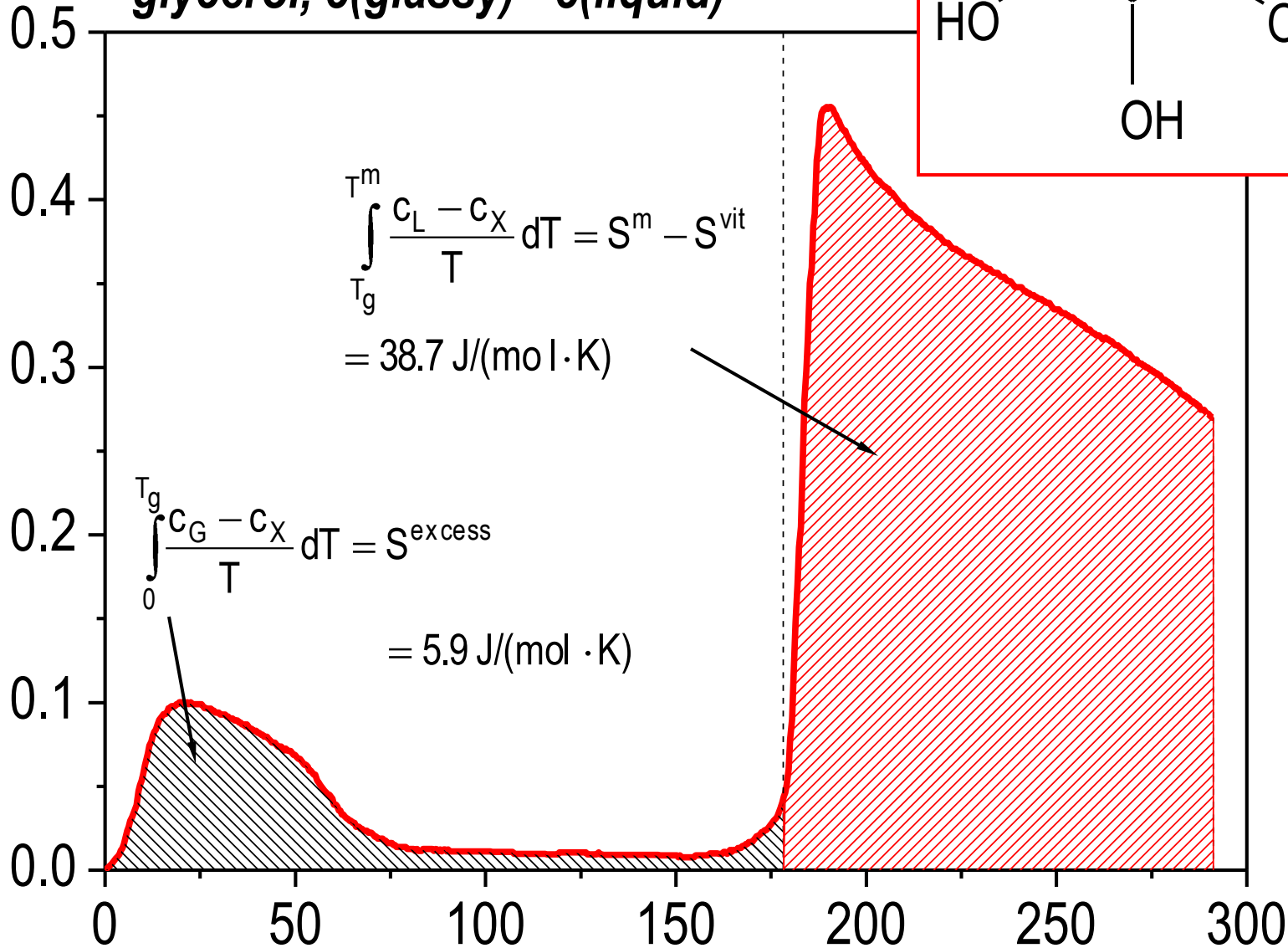
$$\Delta S = R \ln(1/4) = 11.53 \text{ J/(mol}\cdot\text{K)}$$

CO: Clayton & Giauque (1932); CH₃D: Clusius et.al. (1937);
compiled by Wilks (1963); other examples: NO, N₂O, H₂O, D₂O, ...

glycerol, $c(\text{glassy}) - c(\text{liquid})$



$\Delta c/T$ in $\text{J}/(\text{mol}\cdot\text{K}^2)$



T in K

Simon & Lange 1926

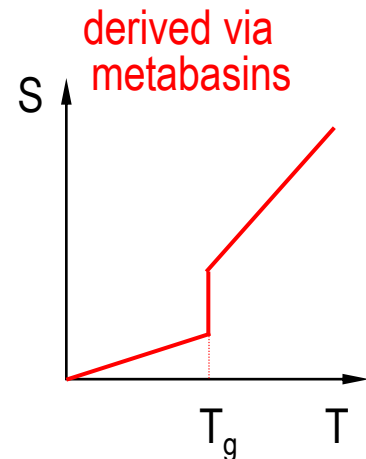
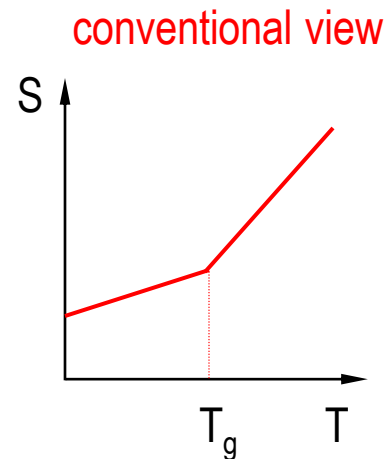
If c_p data are available for a glass and its isochemical crystal from very low T up to T^m , then the following relations hold:

$$\int_0^{T^m + \delta T} \frac{c_{G,L}}{T} dT = S_L(T^m) - S_G(0) \quad \text{for the glass,}$$

$$\int_0^{T^m + \delta T} \frac{c_X}{T} dT = S^m + S_X(T^m) = S_L(T^m) \quad \text{for the crystal.}$$

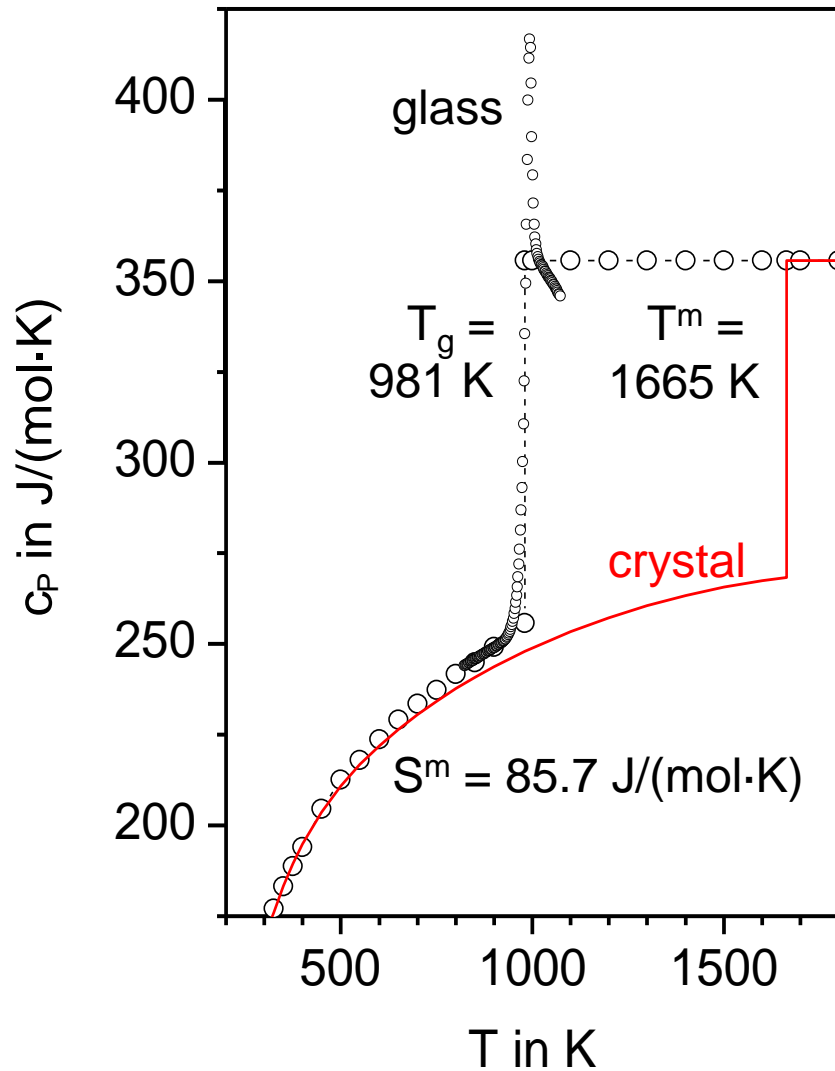
At $T = T^m + \delta T$, both crystal and glass assume the state of eq. liquid. Thus

$$S_G(0) = S_L(T^m) - \int_0^{T^m} \frac{c_{G,L}}{T} dT$$

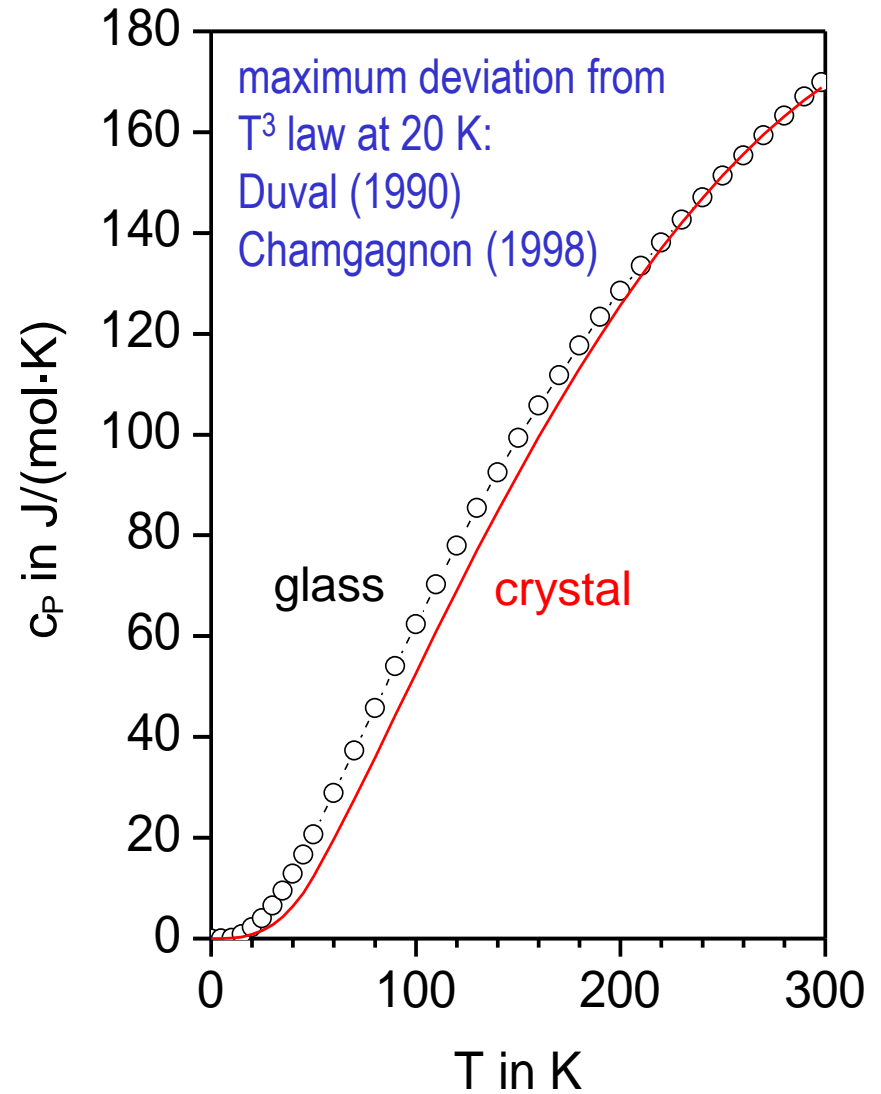


Based on calorimetric data, $S_{\text{glass}}(0)$ is found to be > 0 . This is the basis of the conventional view of entropy of glasses.

example: $\text{CaO} \cdot \text{MgO} \cdot 2\text{SiO}_2$ (diopside)



Robie et al. (1978); Martens et al. (1987)
Richet & Bottinga (1995); Courtial et al. (2000)



Krupka, Robie, Hemingway (1985)
Richet, Robie, Heminway (1986)

diopside $\text{CaMg}[\text{Si}_2\text{O}_6]$

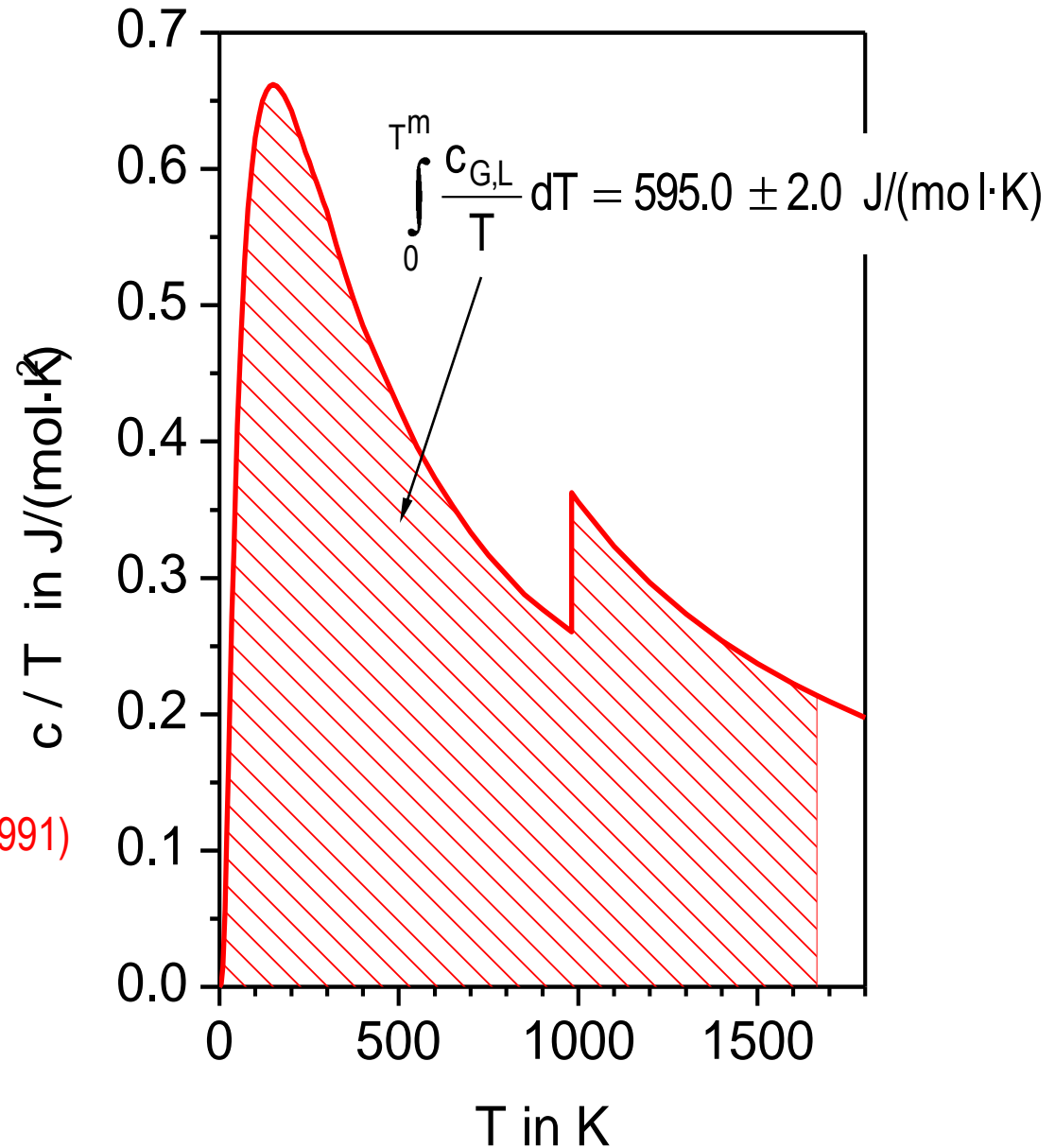
$$S_G(0) = S_L(T^m) - \int_0^{T^m} \frac{c_{G,L}}{T} dT$$

$$S_L(T^m) = 619.8 \pm 1.0 \text{ J}/(\text{mol}\cdot\text{K})^*)$$

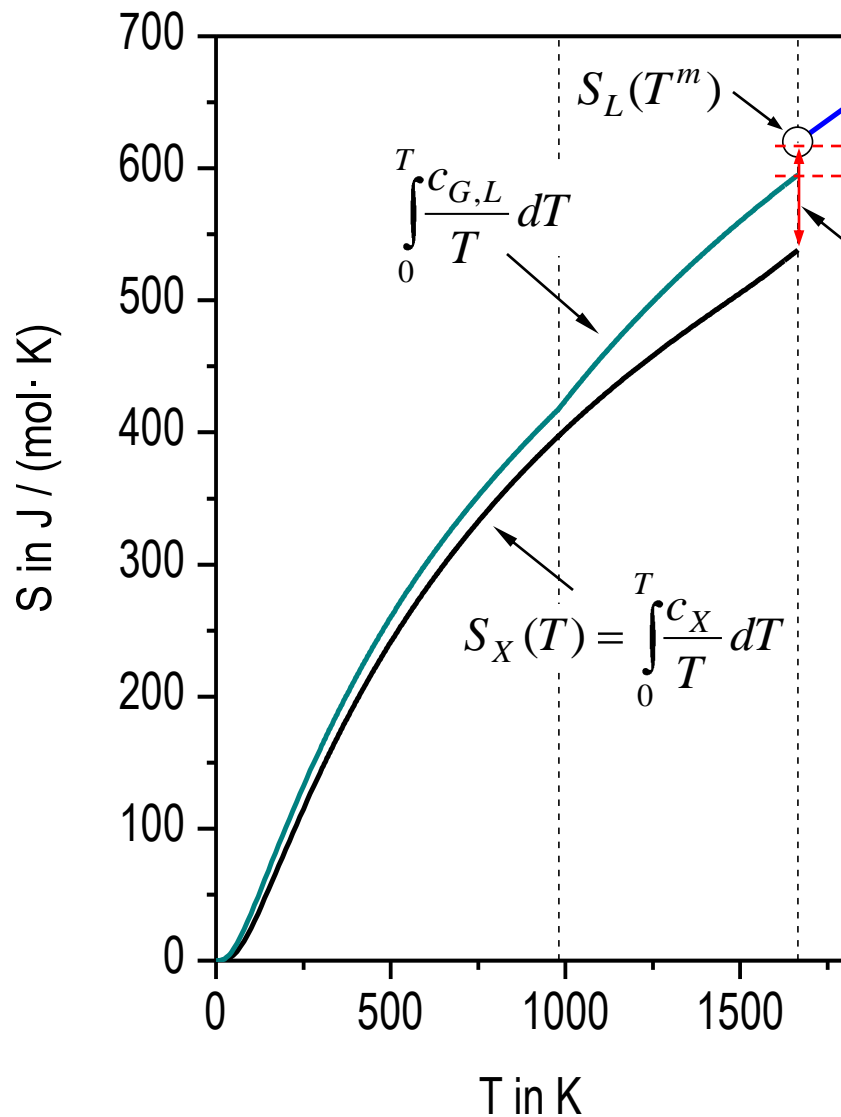
\Rightarrow

$$S_G(0) = 24.8 \pm 3.0 \text{ J}/(\text{mol}\cdot\text{K})$$

*) Knacke, Kubachewski, Hesselmann (1991)



diopside $\text{CaMg}[\text{Si}_2\text{O}_6]$



$$S_L(T^m) - \int_0^{T^m} \frac{c_{G,L}}{T} dT = \Delta$$

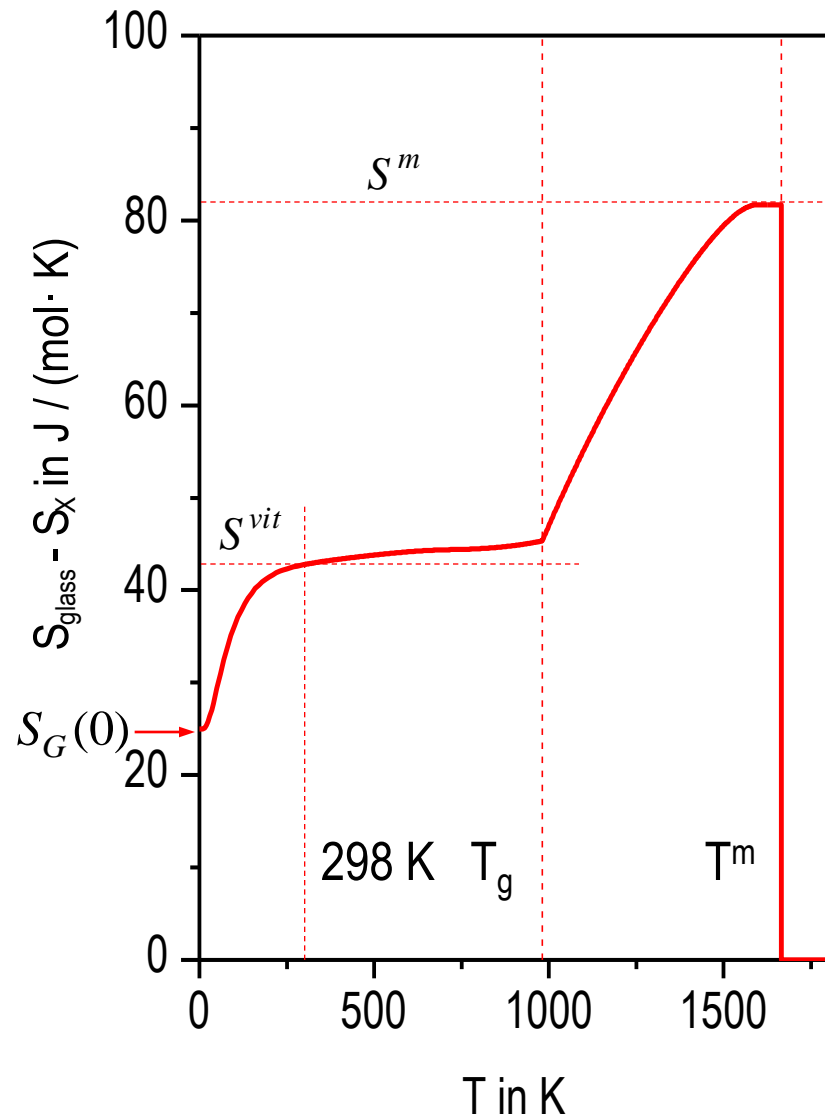
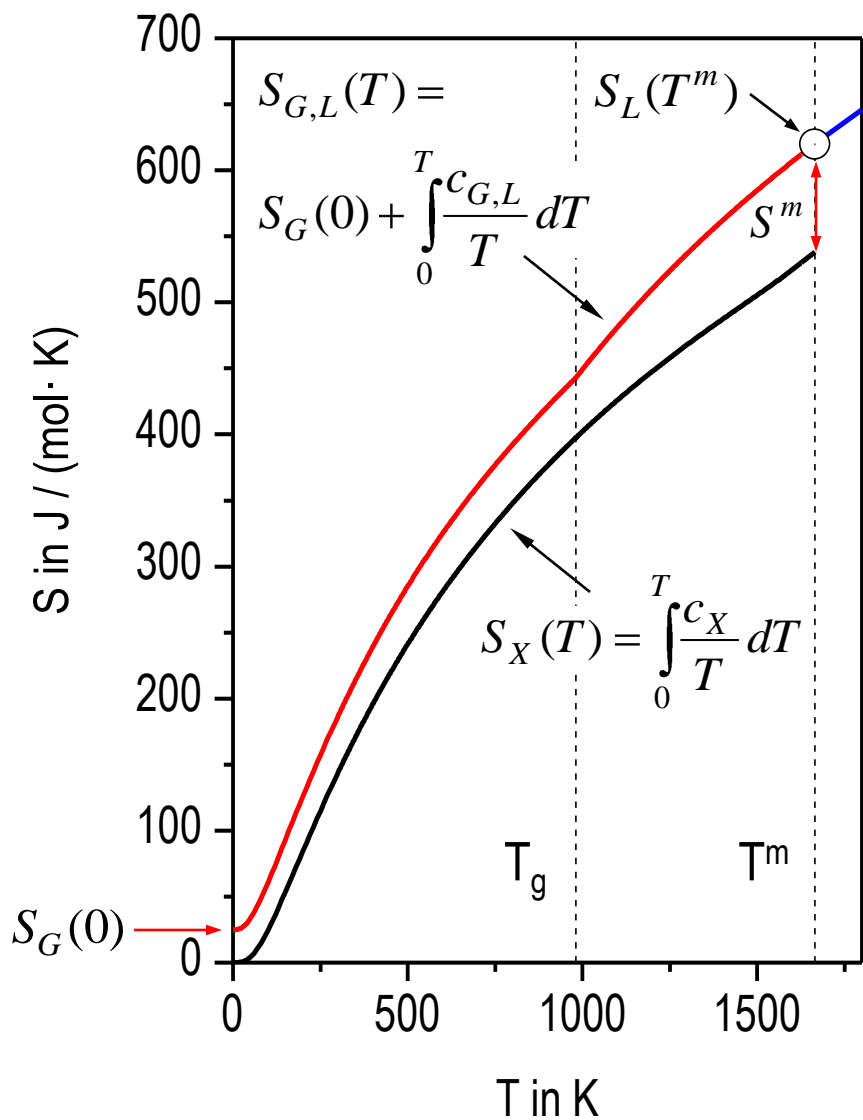
At T^m , both X and G,L must assume the entropy value of the eq. liquid:

$$S_L(T^m) =$$

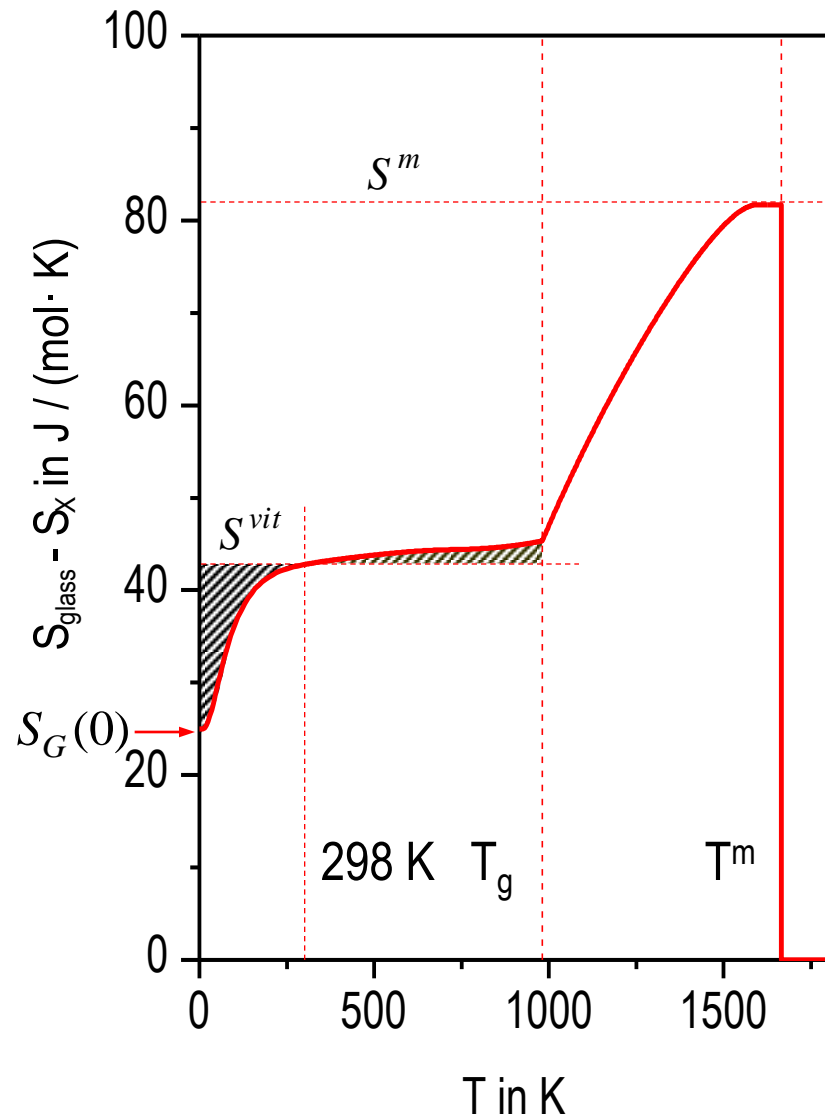
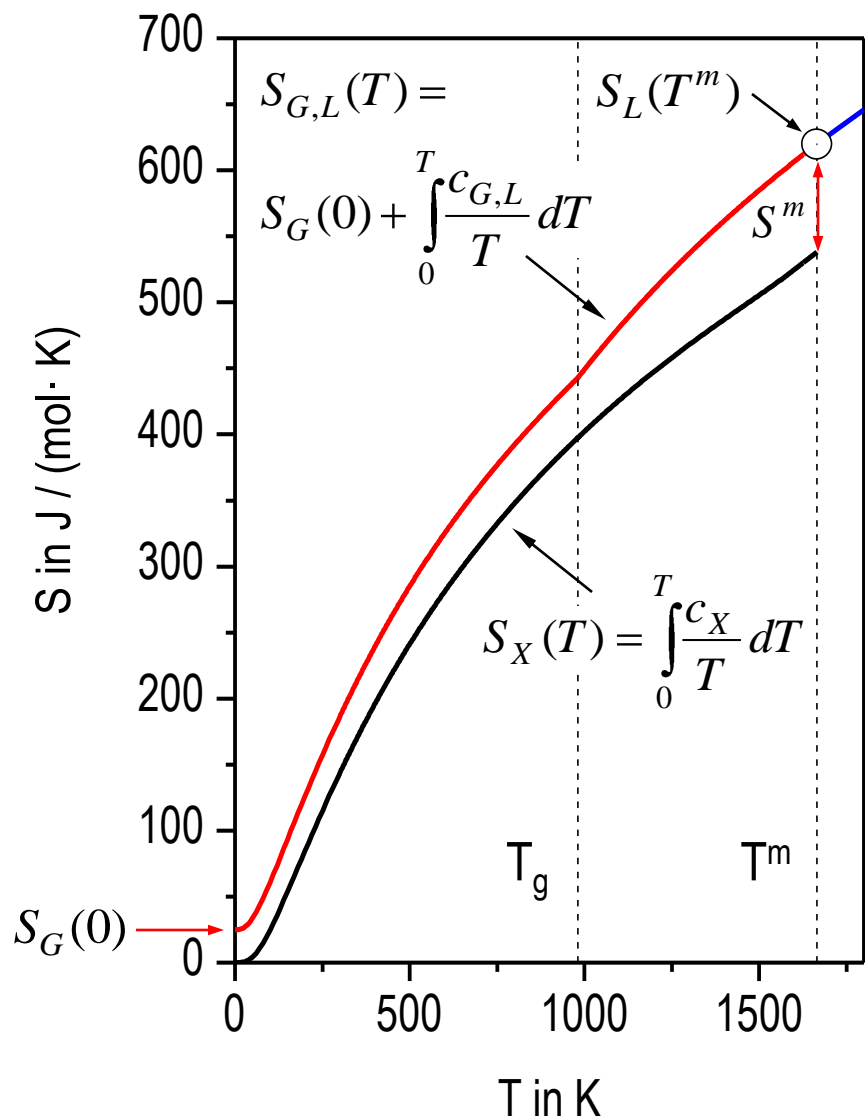
$$S^m + \lim_{T \rightarrow T^m} \left(\underbrace{\int_0^T \frac{c_X}{T} dT}_{S_X(T)} \right) = \lim_{T \rightarrow T^m} \left(\underbrace{\int_0^T \frac{c_{G,L}}{T} dT}_{S_{G,L}(T)} \right)$$

Unless we assume a calorimetrically inaccessible entropy change at $T < T^m$, we have to shift the green curve by Δ towards the state of the eq. liquid.

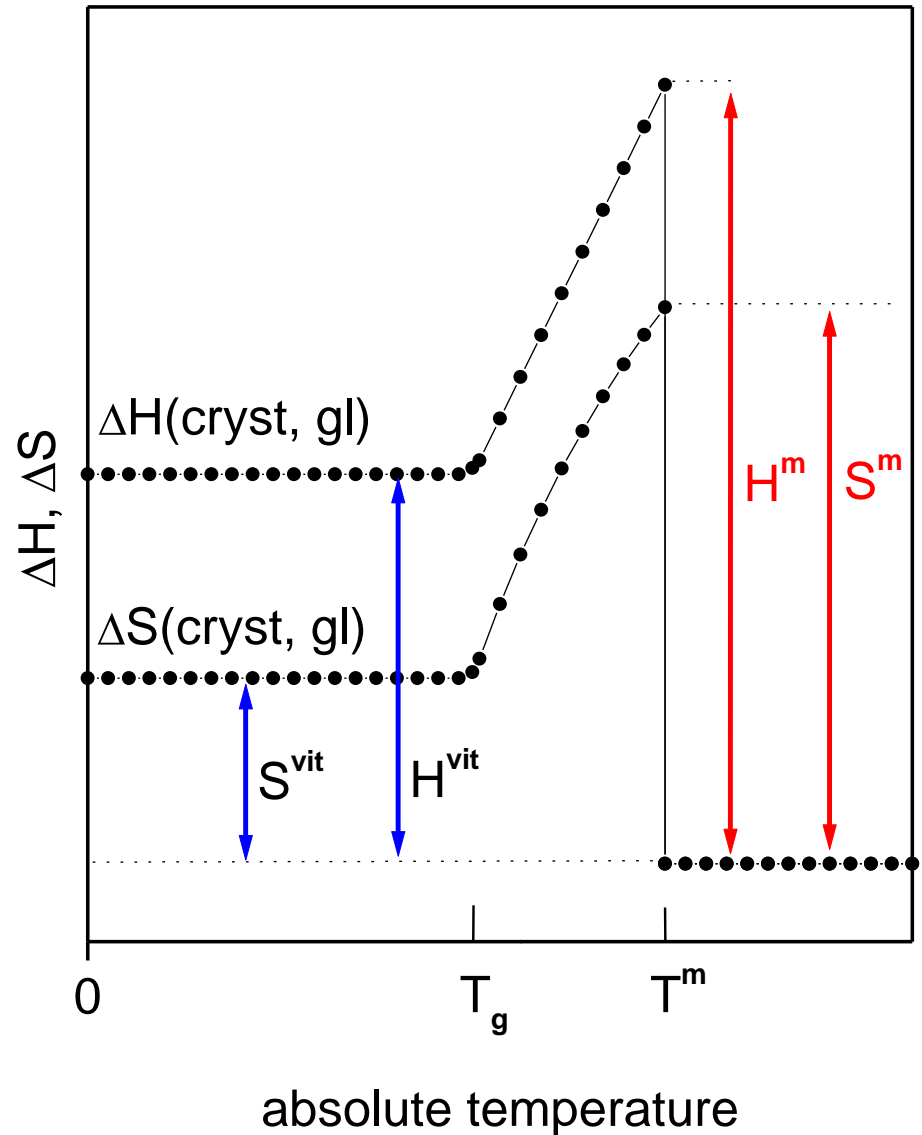
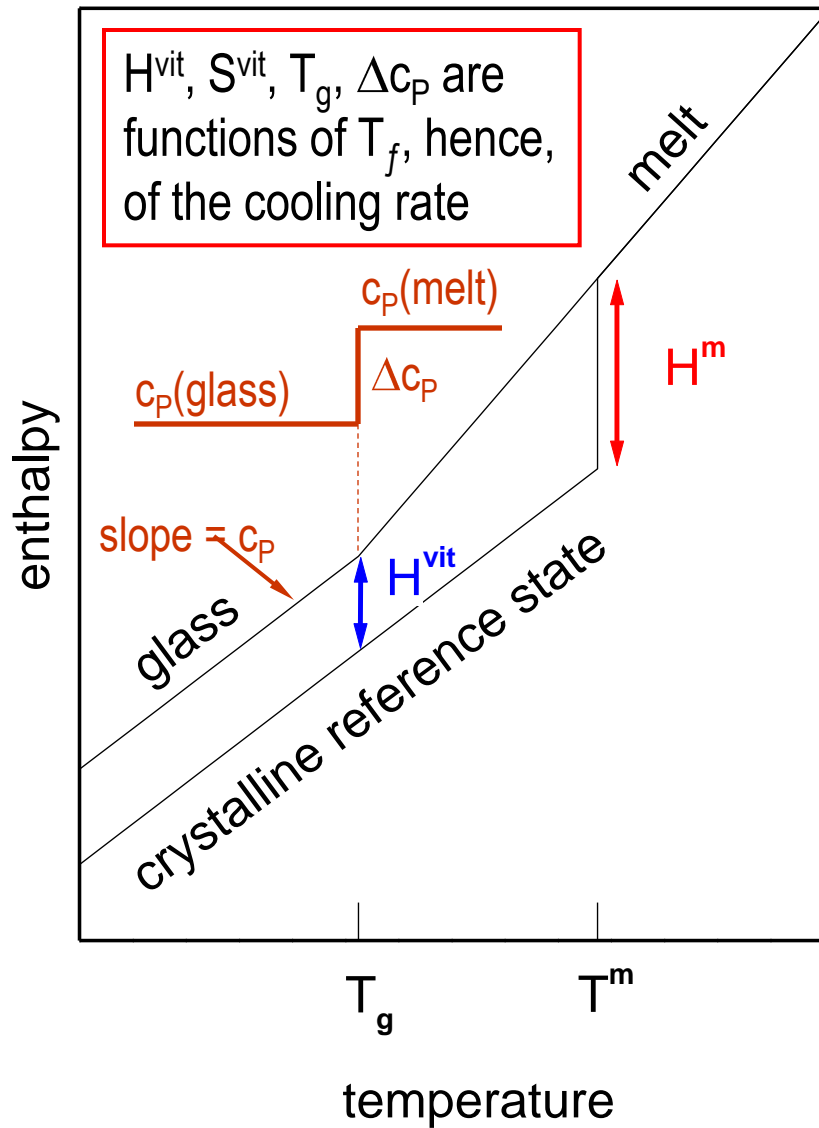
diopside $\text{CaMg}[\text{Si}_2\text{O}_6]$

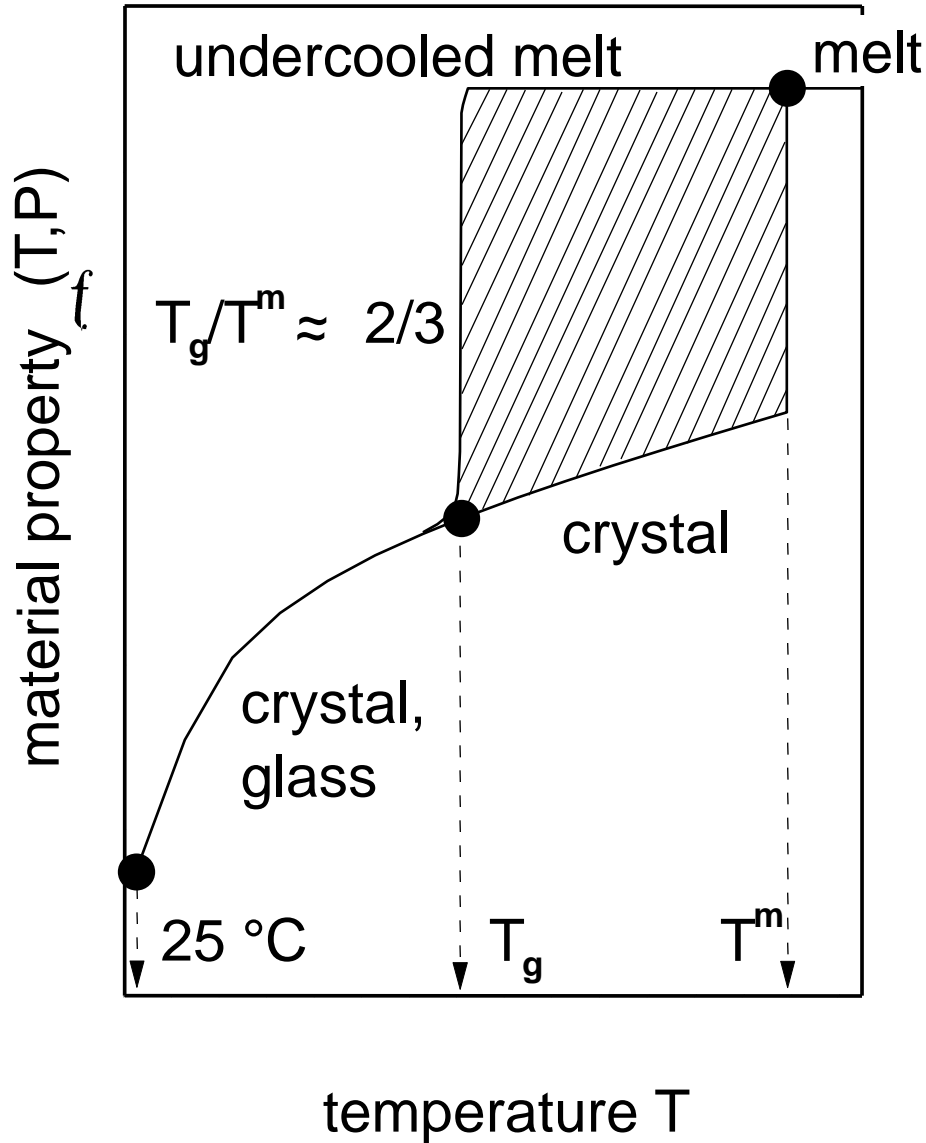


diopside $\text{CaMg}[\text{Si}_2\text{O}_6]$



glasses, melts, and the crystalline reference system





$$\mathbf{G}_{T,P} = \begin{pmatrix} -c_P/T & +V \cdot \alpha_V \\ +V \cdot \alpha_V & -V \cdot \kappa \end{pmatrix}$$

$$c_P = -T \cdot \left(\frac{\partial^2 G}{\partial T^2} \right)_P$$

$$\kappa = -\frac{1}{V} \cdot \left(\frac{\partial^2 G}{\partial P^2} \right)_T$$

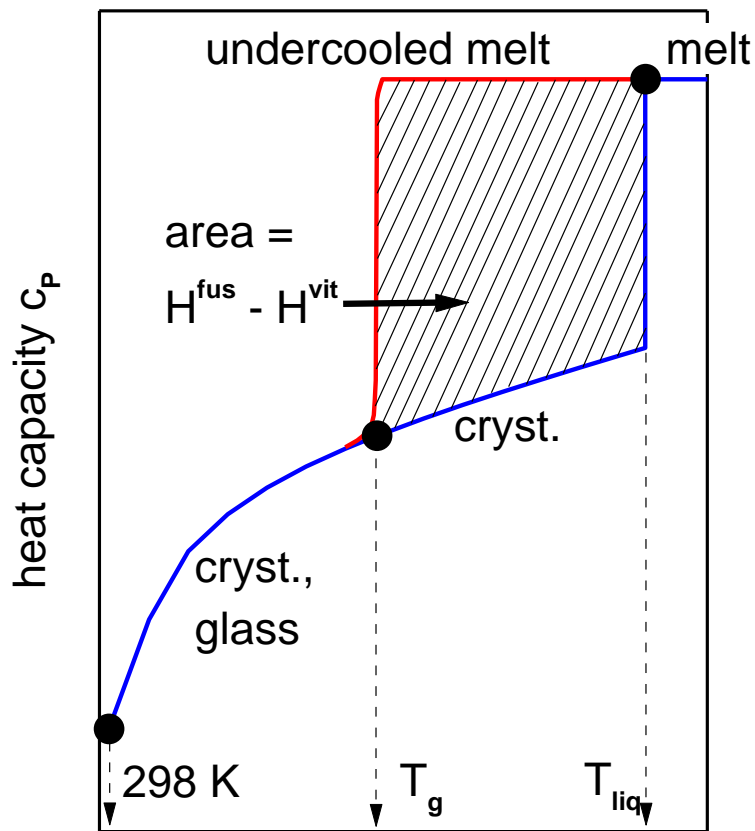
$$\alpha_V = +\frac{1}{V} \cdot \left(\frac{\partial^2 G}{\partial T \partial P} \right)$$

for phases in internal equilibrium:

$$\det(G)_{T,P} = c_P \cdot \kappa \cdot \rho - (\alpha_V)^2 \cdot T > 0$$

for the glass transition:

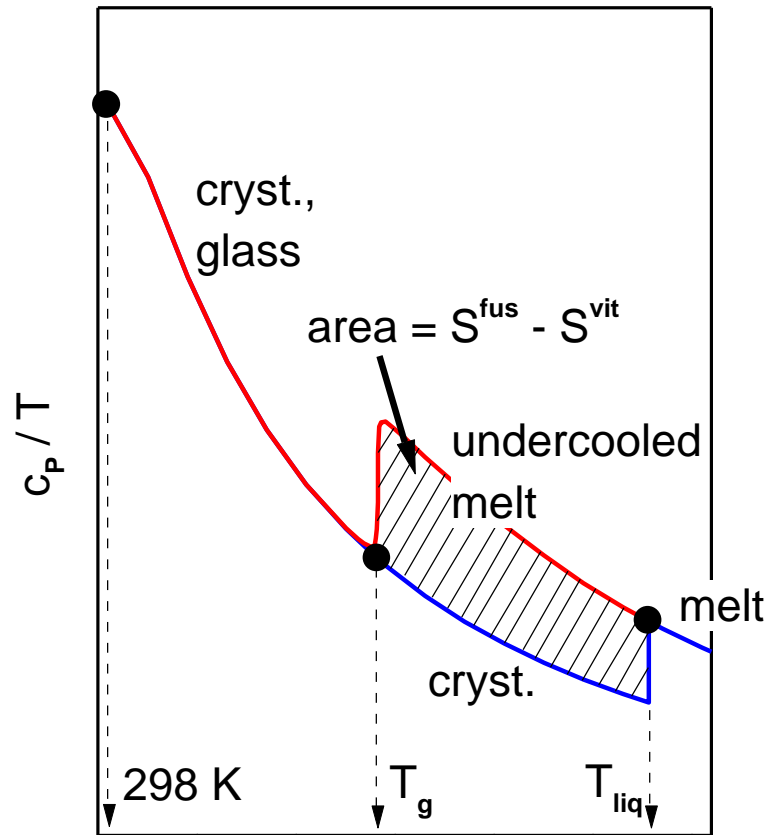
$$\Delta c_P \cdot \Delta \kappa \cdot \rho - (\Delta \alpha_V)^2 \cdot T > 0$$



temperature T

$$H_C(T) = \int_{T_g}^T \Delta c_P dT$$

$$H^{\text{fus}} \approx H^{\text{vit}} + \Delta c_P \cdot (T_{\text{liq}} - T_g)$$



temperature T

$$S_C(T) = \int_{T_g}^T \frac{\Delta c_P}{T} dT$$

$$S^{\text{fus}} \approx S^{\text{vit}} + \Delta c_P \cdot \ln \frac{T_{\text{liq}}}{T_g}$$

- One-component glasses and glass melts are described in relation to a crystalline reference state by the following data set:
 $H^{\text{fus}}, S^{\text{fus}}, T_{\text{liq}}; H^{\text{vit}}, S^{\text{vit}}, \Delta c_p, T_g$
- The set is redundant: Only 4 out of 7 quantities are independent.
- $H^{\text{fus}}, S^{\text{fus}}, T_{\text{liq}}$ are path-independent equilibrium properties, while $H^{\text{vit}}, S^{\text{vit}}, \Delta c_p, T_g$ are path-dependent; they depend on the direction of the path in the (T, P, n_j) space and on the time required to move from state I to state II.

A data set for a one-component glass-forming system thus reads:

H° = standard enthalpy (298 K, 1bar)
 for the Xtal stable at $T = T_g$,

S° = standard entropy (s.a.)

$c_p(T) = A + B \cdot T + C/T^2$

H^{fus} = enthalpy of fusion

S^{fus} = entropy of fusion

T_{liq} = liquidus temperature

$S^{\text{fus}} = H^{\text{fus}} / T_{\text{liq}}$
 two independent quantities

H^{vit} = vitrification enthalpy

S^{vit} = vitrification entropy (zero Kelvin entropy)

T_g = glass transition temperature

Δc_p = jump of the heat capacity at T_g

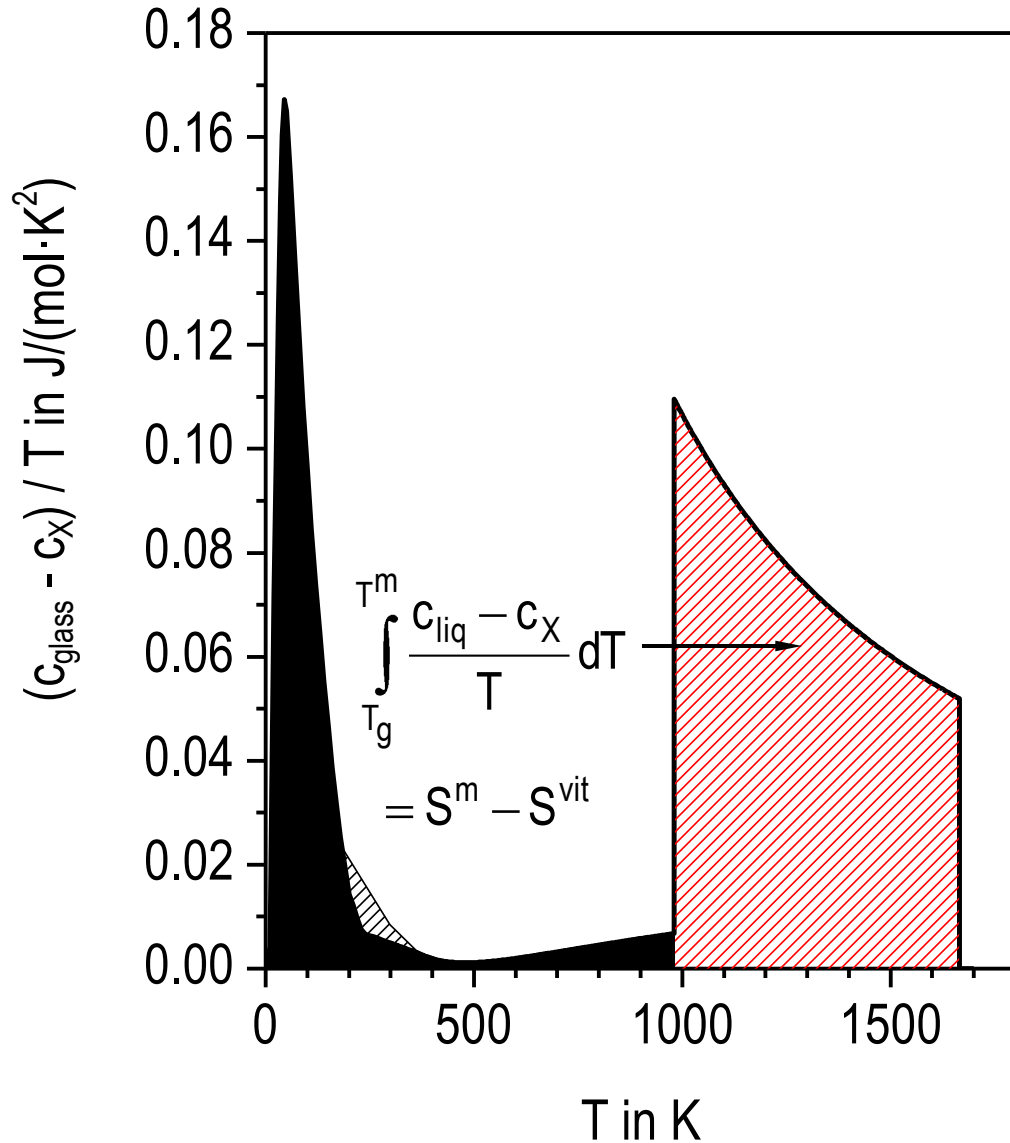
$H^{\text{vit}} \approx H^{\text{fus}} - \Delta c_p \cdot (T_{\text{liq}} - T_g)$

$S^{\text{vit}} \approx S^{\text{fus}} - \Delta c_p \cdot \ln(T_{\text{liq}}/T_g)$

two independent quantities;
 $H^{\text{vit}}, S^{\text{vit}}, T_g, \Delta c_p$ are functions of T_f

short excursion: viscosity

diopside $\text{CaMg}[\text{Si}_2\text{O}_6]$



configurational entropy:

$$S_C(T) = S^{\text{vit}} + \int_{T_g}^T \frac{C_{\text{liq}}(T) - C_X(T)}{T} dT$$

$$\approx S^{\text{vit}} + \Delta c \cdot \ln \frac{T}{T_g}$$

see, e.g., the Adam-Gibbs equation:

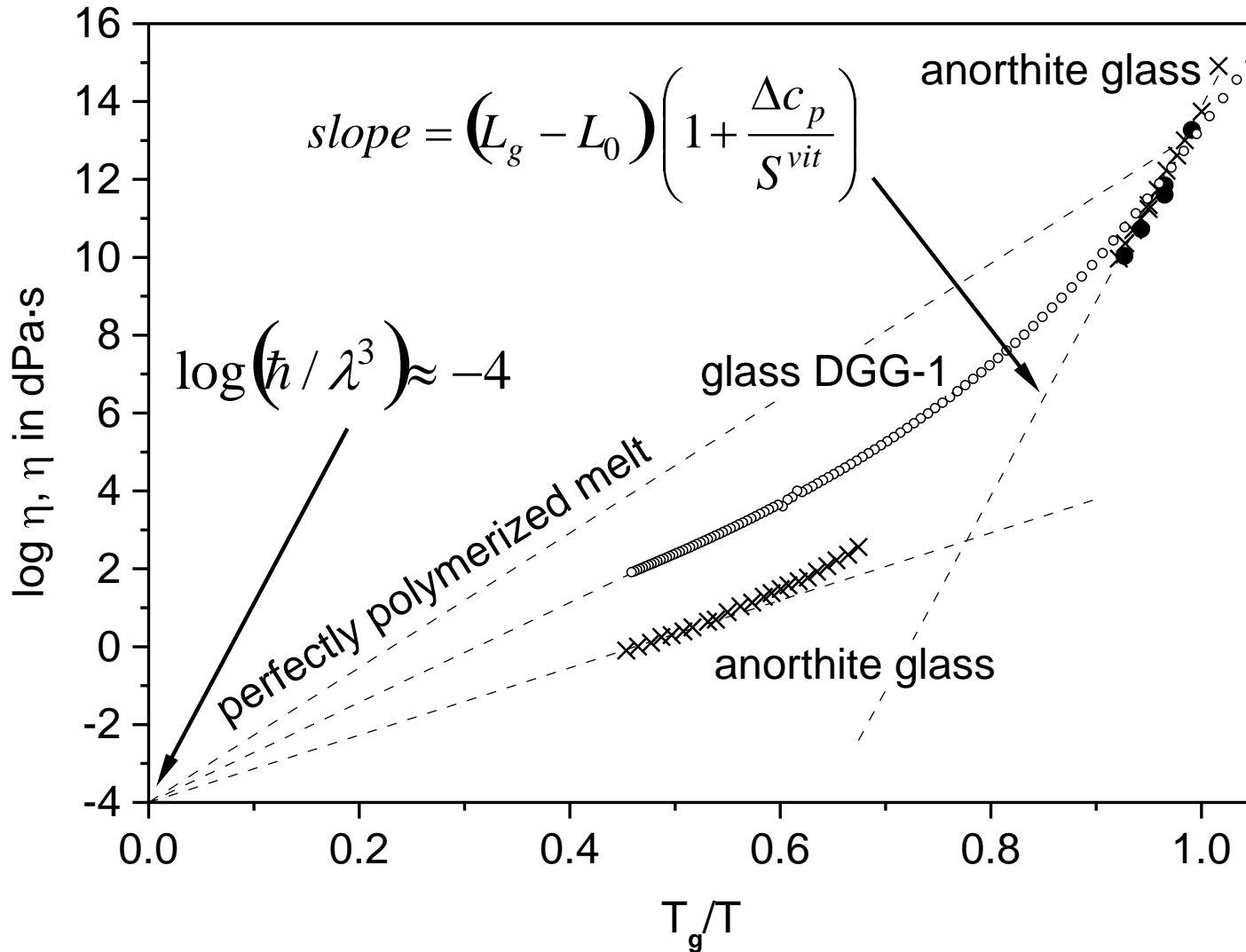
$$L = L_\infty + \frac{C}{T \cdot S_C(T)}; \quad L \equiv \log \eta$$

$$= L_\infty + \frac{\left(L_g - L_\infty \right) \frac{T_g}{T}}{1 + \frac{\Delta c}{S^{\text{vit}}} \cdot \ln \frac{T}{T_g}}$$

viscosity slope for $T \rightarrow T_g$:

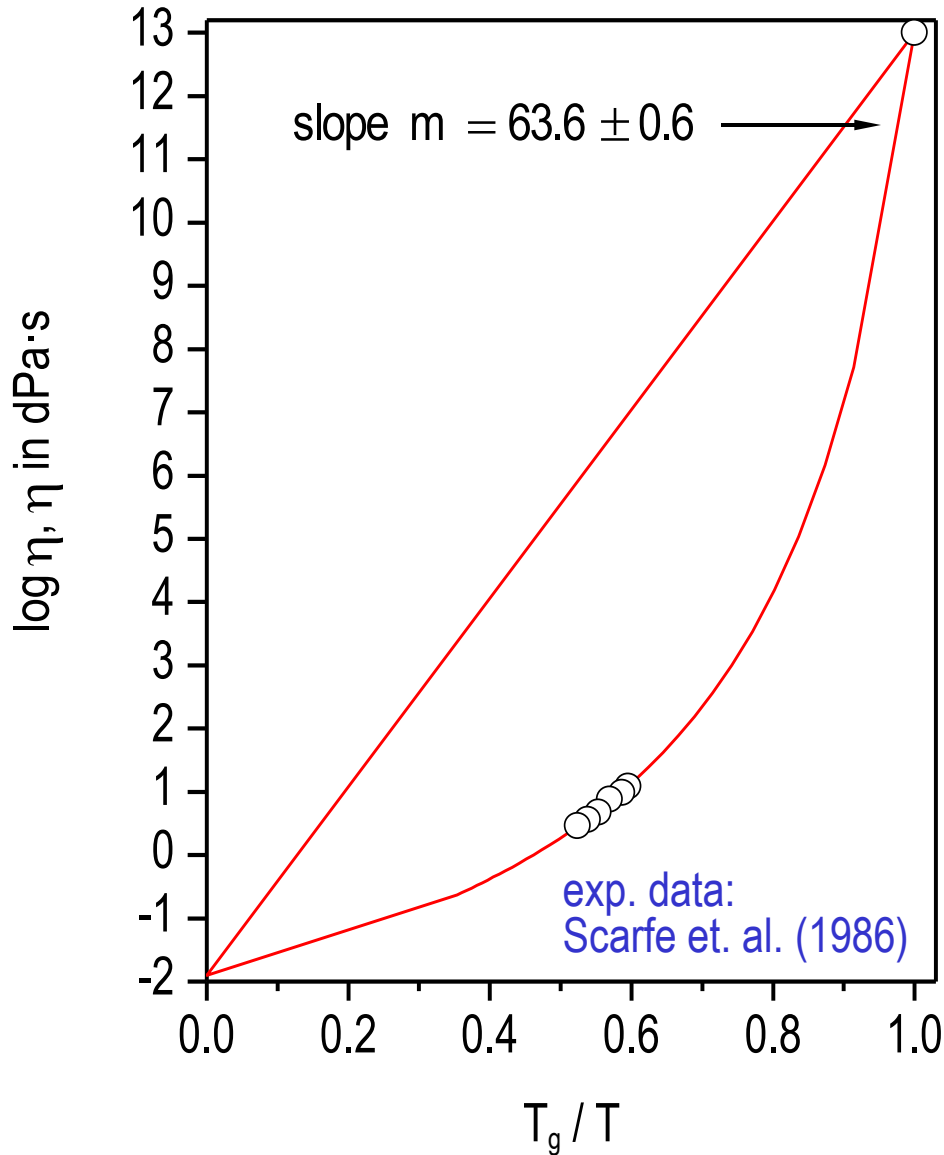
$$\left. \frac{\partial L}{\partial (T_g / T)} \right|_{T=T_g} = \left(L_g - L_\infty \right) \left(1 + \frac{\Delta c}{S^{\text{vit}}} \right)$$

Angell's plot of viscosity



high-T limit of diffusion: $D_k^* = \frac{\lambda^2}{3} \cdot \frac{kT}{h}$ Stokes-Einstein: $D_k^* = \frac{kT}{6\pi \cdot \eta \cdot \lambda} \Rightarrow \eta_\infty = \hbar / \lambda^3$

diopside $\text{CaMg}[\text{Si}_2\text{O}_6]$



$$L = \log \eta = A + \frac{B}{T - T_0}$$

$$A = -1.899$$

$$T_g = 708 \text{ } ^\circ\text{C}$$

$$B = 2493.4$$

$$T_0 = 540.6$$

$$\pm \delta \log \eta = 0.022$$

$$\left. \frac{\partial \log \eta}{\partial T_g / T} \right|_{T=T_g} = \frac{T_g \cdot B}{(T - T_0)^2} = m$$

$$\Rightarrow m = 63.3 \pm 0.6$$

$$m = \underbrace{L_g - L_0}_{14.9} \left(1 + \frac{\Delta c_{G,L}}{S_G(0)} \right)$$

$$c_{G,L}(T_g) = 254.6 \text{ J}/(\text{mol}\cdot\text{K})$$

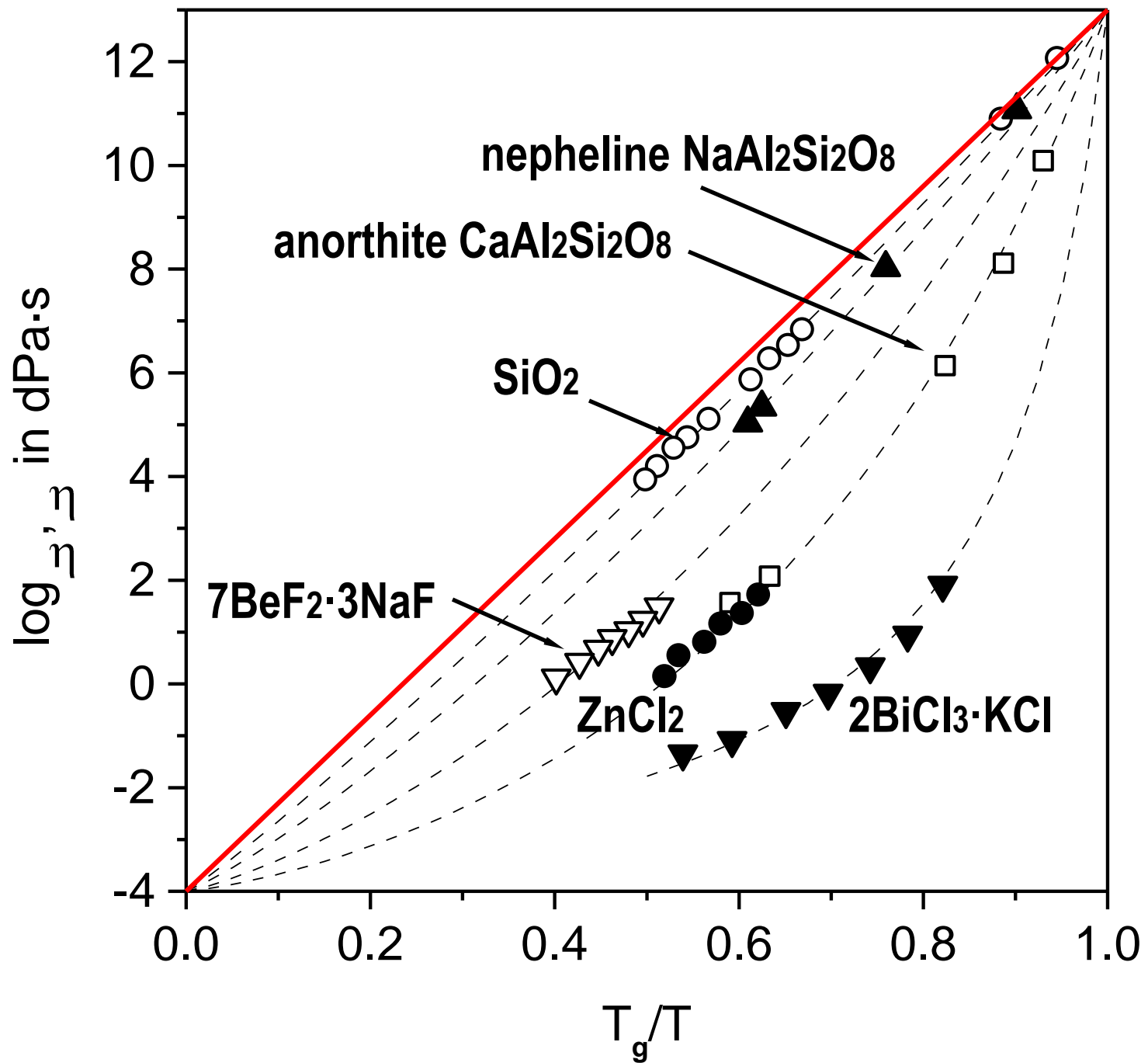
$$c_X(T_g) = 246.5 \text{ J}/(\text{mol}\cdot\text{K})$$

$$3 \cdot N \cdot R = 249.0 \text{ J}/(\text{mol}\cdot\text{K})$$

$$\Delta c_{G,L} = 80.0 \text{ J}/(\text{mol}\cdot\text{K})$$

$$S_G(0) = 24.8 \text{ J}/(\text{mol}\cdot\text{K})$$

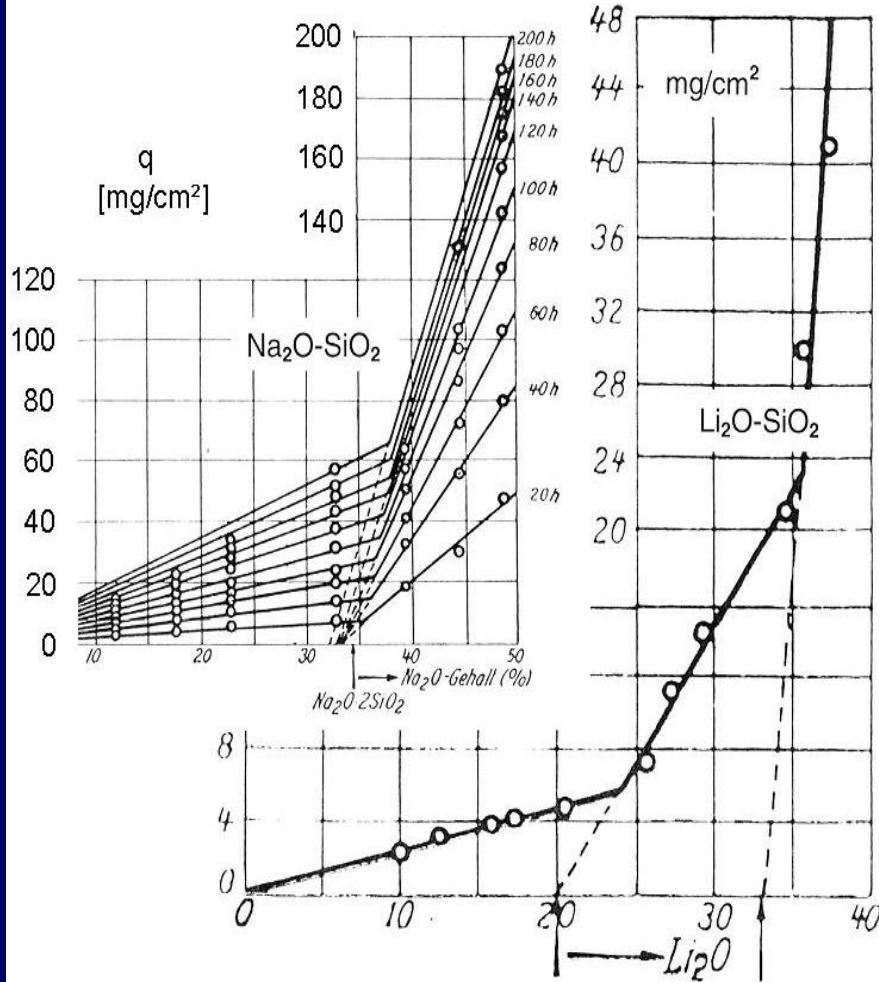
$$\Rightarrow m = 63.0$$



the crystalline reference system

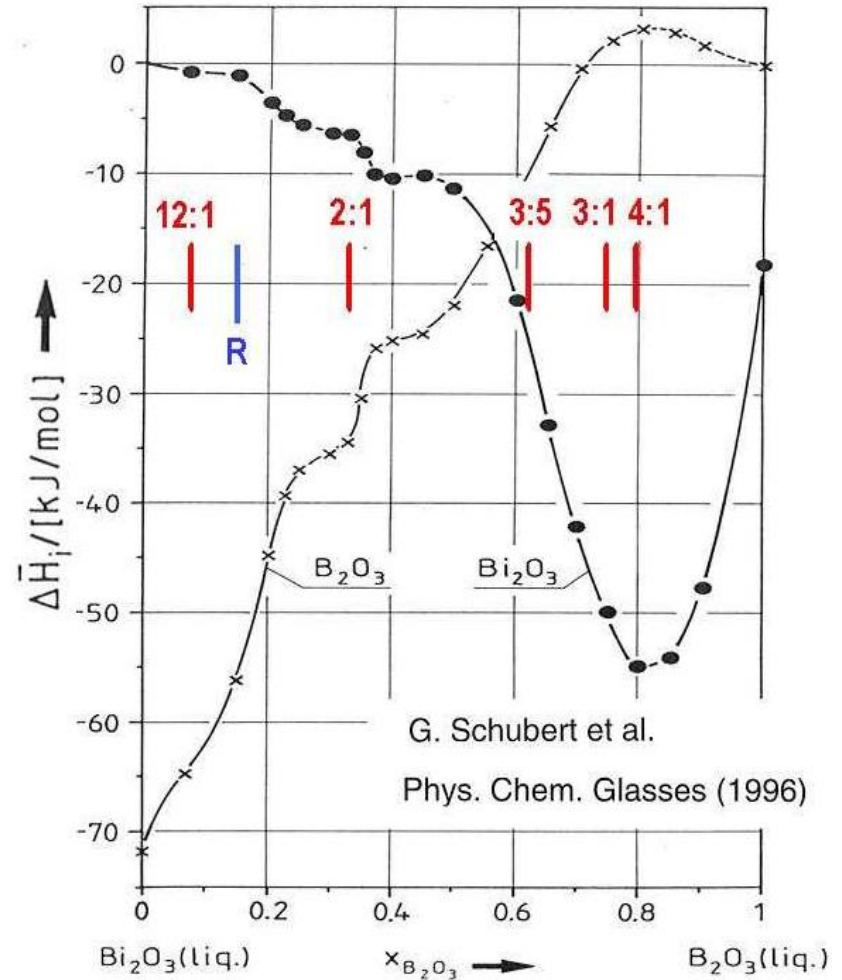
reflections of solid state phase relations clearly seen in the properties of high-T liquids:

evaporation loss by transpiration method



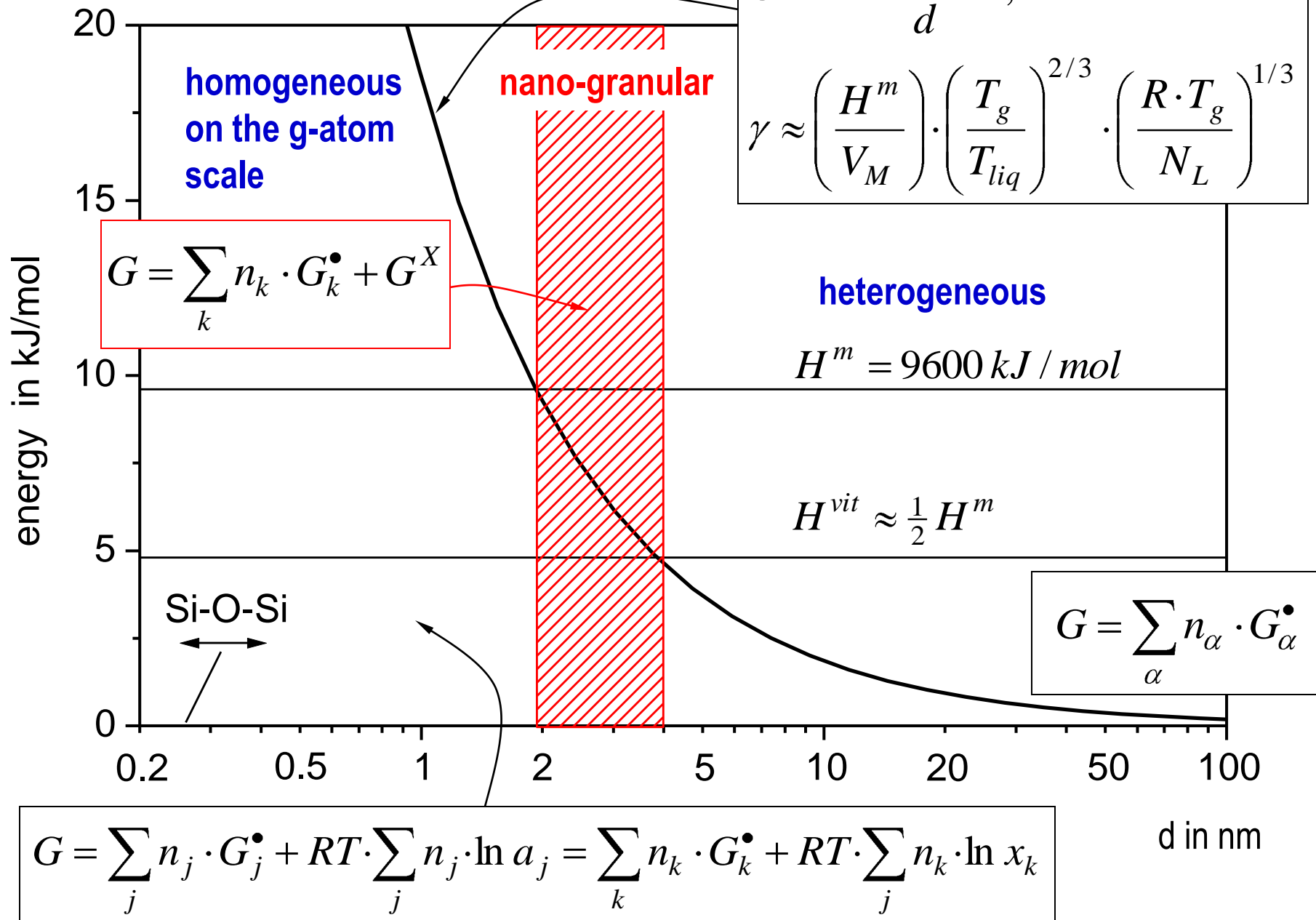
W.E.S. Turner, 1934

solution calorimetry



see also: $\text{Na}_2\text{O}-\text{MoO}_3$, $-\text{WO}_3$, $-\text{V}_2\text{O}_5$; $\text{PbO}-\text{B}_2\text{O}_3$
Kaufmann, 2000

glasses, melts, and M.R.O.



$$H_{glass}^{\circ} = \sum_k n_k \cdot (H_k^{\circ} + H_k^{vit})$$

$$H_{1673,liq}^{\circ} = \sum_k n_k \cdot H_{1673,liq,k}^{\circ}$$

$$S_{glass}^{\circ} = \sum_k n_k \cdot (S_k^{\circ} + S_k^{vit})$$

$$S_{1673,liq}^{\circ} = \sum_k n_k \cdot S_{1673,liq,k}^{\circ}$$

$$C_{P,liq} = \sum_k n_k \cdot C_{P,liq,k}$$

$$H_{T,liq} = H_{1673,liq}^{\circ} + C_{P,liq} \cdot (T - 1673)$$

$$S_{T,liq} = S_{1673,liq}^{\circ} + C_{P,liq} \cdot \ln(T / 1673)$$

$$\Delta H_{T,liq} = H_{T,liq} - H_{glass}^{\circ}$$

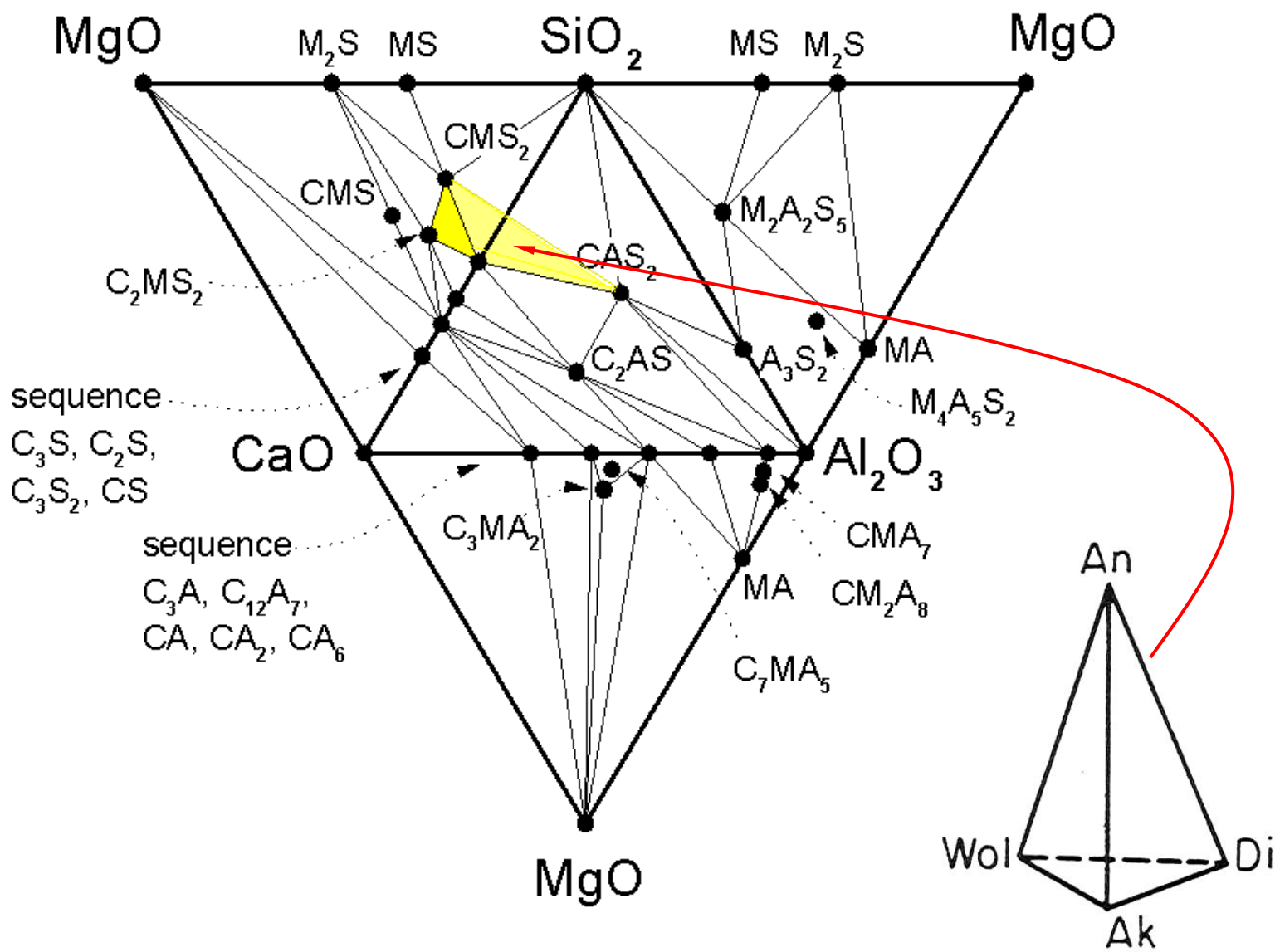
$$\ln a_j = \sum_k \left[A_{jk} \cdot \frac{G_k^{\circ}}{RT} \right] - \frac{G_j^{\circ}}{RT}$$

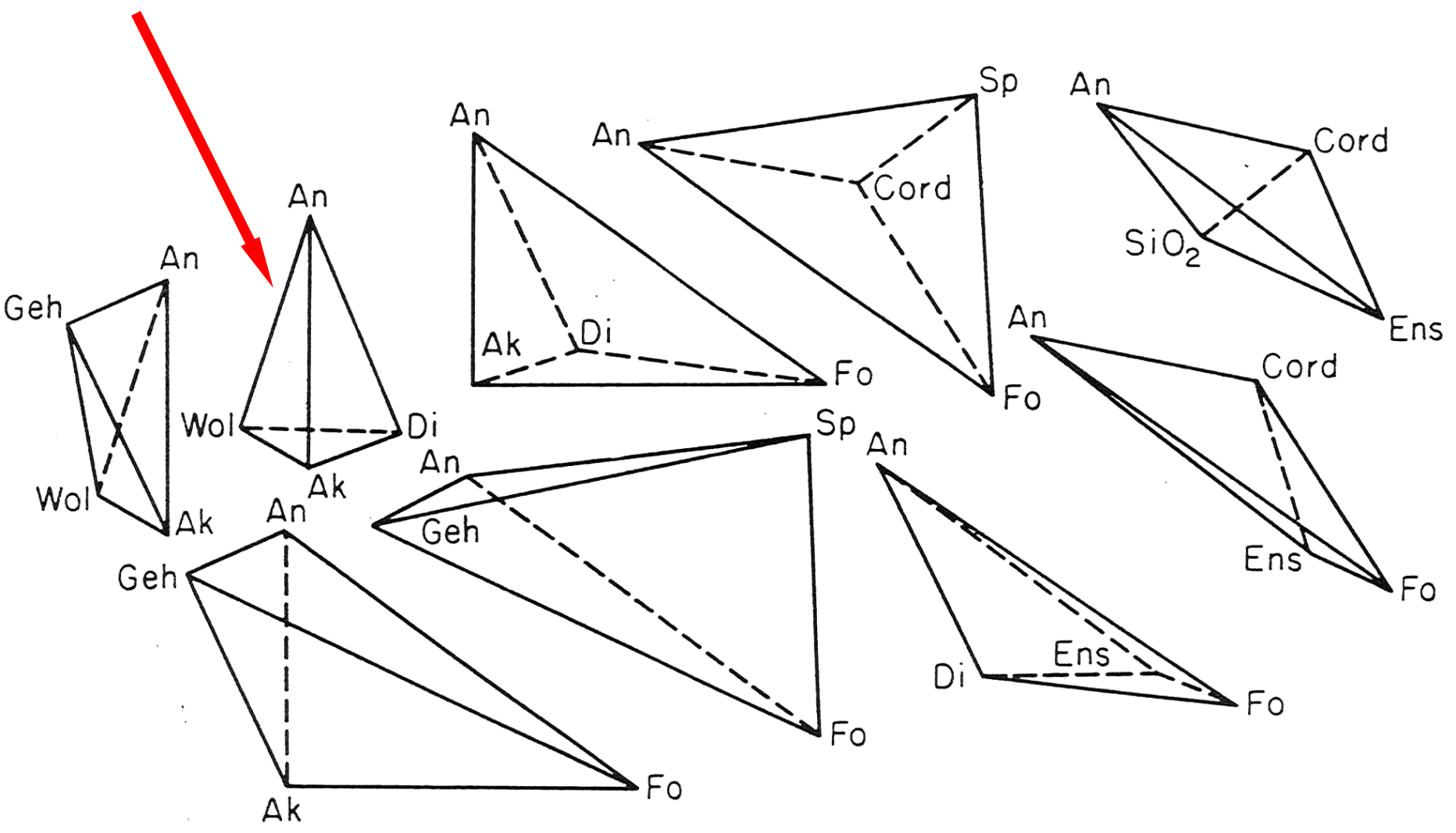
k	-H°	S°	H ^{vit}	S ^{vit}	-H _{1673,liq}	S _{1673,liq}	C _{P,liq}
P ₂ O ₅ ·3CaO	4117.1	236.0	135.1	51.5	3417.1	898.7	324.3
P ₂ O ₅	1492.0	114.4	18.2	9.5	1138.5	586.6	181.6
Fe ₂ O ₃	823.4	87.4	45.2	17.2	550.2	370.3	142.3
FeO·Fe ₂ O ₃	1108.8	151.0	82.8	31.4	677.8	579.9	213.4
FeO·SiO ₂	1196.2	92.8	36.7	13.8	962.3	342.7	139.7
2FeO·SiO ₂	1471.1	145.2	55.2	20.5	1118.8	512.1	240.6
MnO·SiO ₂	1320.9	102.5	40.2	15.1	1085.3	345.2	151.5
2ZnO·SiO ₂	1643.1	131.4	82.4	31.4	1261.1	494.5	174.5
ZrO ₂ ·SiO ₂	2034.7	84.5	86.6	32.6	1686.2	381.2	149.4
CaO·TiO ₂	1660.6	93.7	67.4	25.5	1365.7	360.2	124.7
BaO·Al ₂ O ₃ ·2SiO ₂	4222.1	236.8	130.5	95.4	3454.3	1198.3	473.2
BaO·2SiO ₂	2553.1	154.0	81.6	26.8	2171.1	533.5	241.4
BaO·SiO ₂	1618.0	104.6	56.5	41.0	1349.8	361.1	146.4
Li ₂ O·Al ₂ O ₃ ·4SiO ₂	6036.7	308.8	184.1	12.1	5235.4	1173.2	498.7
Li ₂ O·SiO ₂	1648.5	79.9	16.7	6.3	1416.7	339.7	167.4
K ₂ O·Al ₂ O ₃ ·6SiO ₂	7914.0	439.3	106.3	29.3	6924.9	1559.4	765.7
K ₂ O·Al ₂ O ₃ ·2SiO ₂	4217.1	266.1	80.4	22.1	3903.7	666.5	517.6
K ₂ O·4SiO ₂	4315.8	265.7	26.4	21.3	3697.8	983.7	410.0
K ₂ O·2SiO ₂	2508.7	190.6	12.6	23.9	2153.1	595.4	275.3
Na ₂ O·Al ₂ O ₃ ·6SiO ₂	7841.2	420.1	125.0	28.4	6870.1	1512.5	648.1
Na ₂ O·Al ₂ O ₃ ·2SiO ₂	4163.5	248.5	92.0	27.9	3614.1	856.9	423.8

**kJ/mol;
J/(mol·K)**

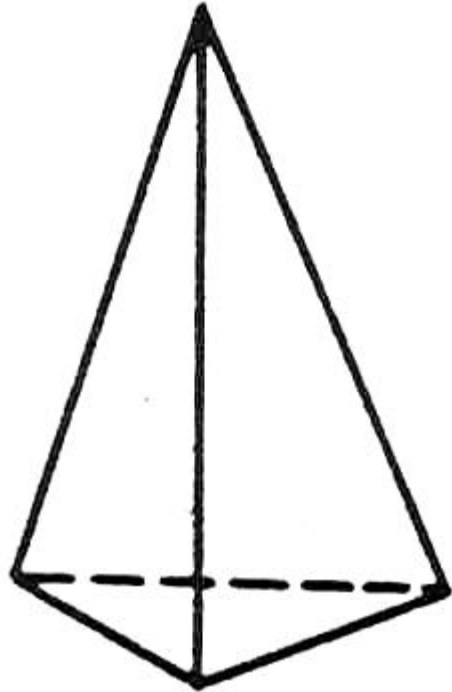
k	-H°	S°	H ^{vit}	S ^{vit}	-H _{1673,liq}	S _{1673,liq}	C _{p,liq}
B ₂ O ₃	1273.5	54.0	18.2	11.3	1088.7	271.1	129.7
Na ₂ O·B ₂ O ₃ ·4SiO ₂	5710.9	270.0	42.7	21.1	4988.0	1090.2	637.6
Na ₂ O·4B ₂ O ₃	5902.8	276.1	58.3	40.1	4986.7	1275.5	704.2
Na ₂ O·2B ₂ O ₃	3284.9	189.5	48.8	26.6	2735.9	780.3	444.8
Na ₂ O·B ₂ O ₃	1958.1	147.1	43.6	19.5	1585.7	538.7	292.9
2MgO·2Al ₂ O ₃ ·5SiO ₂	9113.2	407.1	135.8	41.4	7994.8	1606.2	1031.8
MgO·SiO ₂	1548.5	67.8	46.6	13.6	1318.0	296.2	146.4
2MgO·SiO ₂	2176.9	95.4	61.4	11.0	1876.1	402.9	205.0
CaO·MgO·2SiO ₂	3202.4	143.1	92.3	25.7	2733.4	621.7	355.6
2CaO·MgO·2SiO ₂	3876.9	209.2	106.7	32.0	3319.2	775.3	426.8
CaO·Al ₂ O ₃ ·2SiO ₂	4223.7	202.5	103.0	37.7	3628.8	791.2	380.7
2CaO·Al ₂ O ₃ ·SiO ₂	3989.4	198.3	129.9	49.4	3374.0	787.8	299.2
3Al ₂ O ₃ ·2SiO ₂	6820.8	274.9	188.3	71.5	5816.2	1231.8	523.4
CaO·SiO ₂	1635.1	83.1	49.8	18.8	1382.0	329.7	146.4
2CaO·SiO ₂	2328.4	120.5	101.3	38.5	1868.2	509.2	174.5
Na ₂ O·2SiO ₂	2473.6	164.4	29.3	13.2	2102.5	588.7	261.1
Na ₂ O·SiO ₂	1563.1	113.8	37.7	9.8	1288.3	415.1	179.1
Na ₂ O·3CaO·6SiO ₂	8363.8	461.9	77.3	20.5	7372.6	1555.6	786.6
Na ₂ O·2CaO·3SiO ₂	4883.6	277.8	57.7	13.4	4240.9	990.4	470.3
2Na ₂ O·CaO·3SiO ₂	4763.0	309.6	87.0	22.6	4029.6	1107.9	501.2
SiO ₂	908.3	43.5	6.9	4.0	809.6	157.3	86.2

**kJ/mol;
J/(mol·K)**





anorthite $An = CaO \cdot Al_2O_3 \cdot 2SiO_2$



matrix calculation is not restricted to spaces with 4 components only!

wollastonite

diopside

$CaO \cdot SiO_2 =$

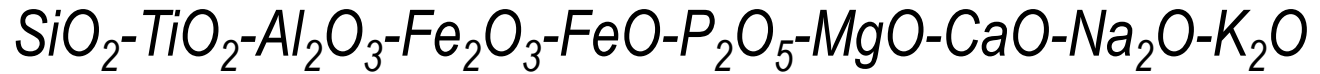
Wol

$Di = CaO \cdot MgO \cdot 2SiO_2$

akermanite $Ak = 2CaO \cdot MgO \cdot 2SiO_2$

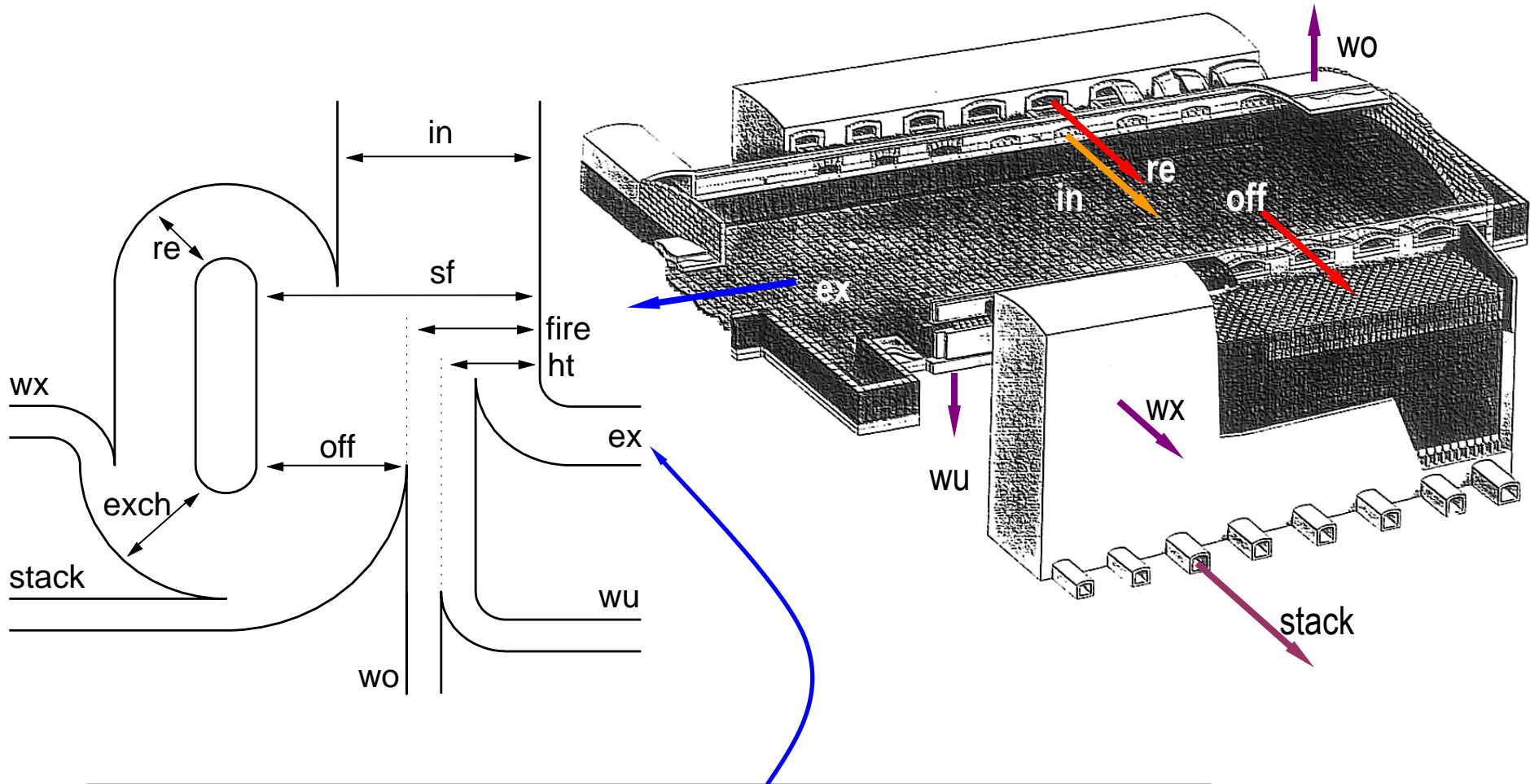
$$\begin{pmatrix} n(an) \\ n(wol) \\ n(ak) \\ n(di) \end{pmatrix} = \begin{pmatrix} 2 & 1 & 2 & 2 \\ 1 & 0 & 0 & 0 \\ 1 & 1 & 2 & 1 \\ 0 & 0 & 1 & 1 \end{pmatrix}^{-1} \cdot \begin{pmatrix} n(SiO_2) \\ n(Al_2O_3) \\ n(CaO) \\ n(MgO) \end{pmatrix}$$

prediction of G^f for 4 mineral fibre compositions



	G_{el}^f [kJ/mol]		G_{ox}^f [kJ/mol]	
	exp	calc	exp	calc
glass a	-852.0	-849.6	-12.9	-10.6
glass b	-865.0	-867.2	-35.4	-37.7
glass c	-880.8	-881.8	-34.0	-35.0
glass d	-885.4	-852.7	-44.1	-41.4

experiments by Richet et al.



raw materials (25 °C) → glass + batch gases (25 °C) ΔH_{chem}

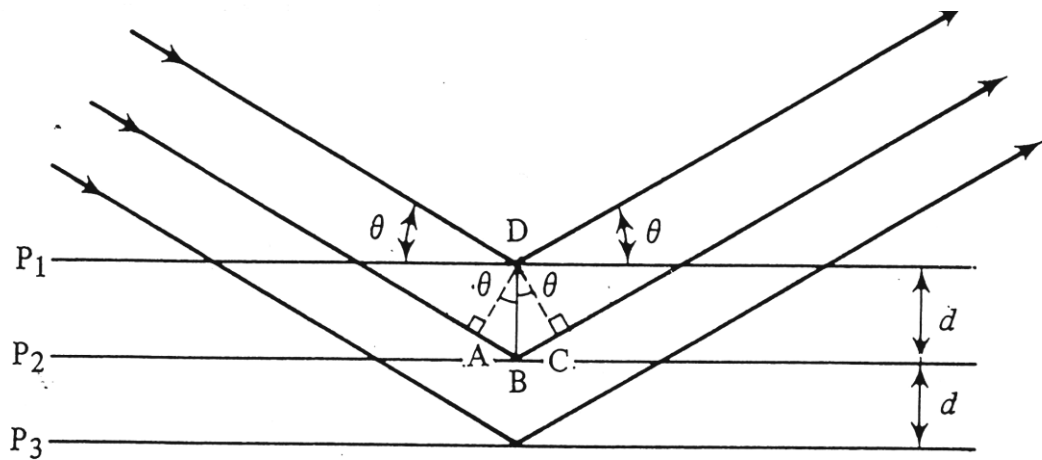
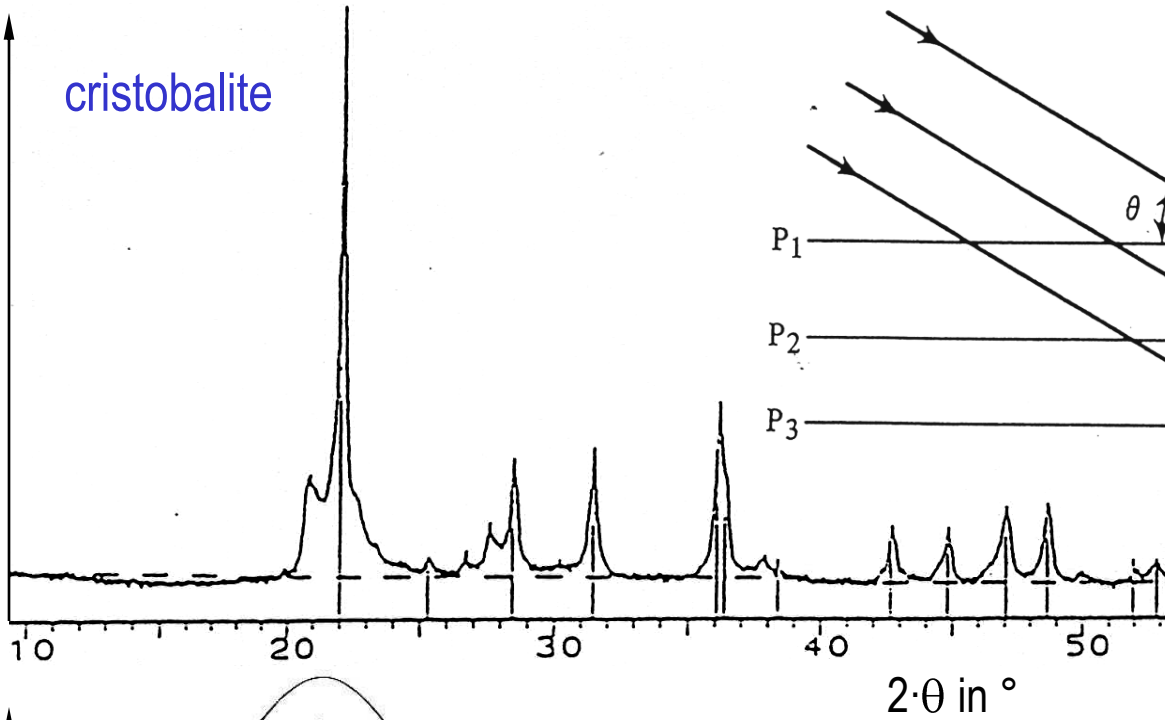
glass (25 °C) → glass melt ($T = T_{\text{ex}}$) $\Delta H_{T=T_{\text{ex}}}$

$$H_{\text{ex}} = (1 - y_C) \cdot \Delta H_{\text{chem}} + \Delta H_{T=T_{\text{ex}}}$$

structural aspects

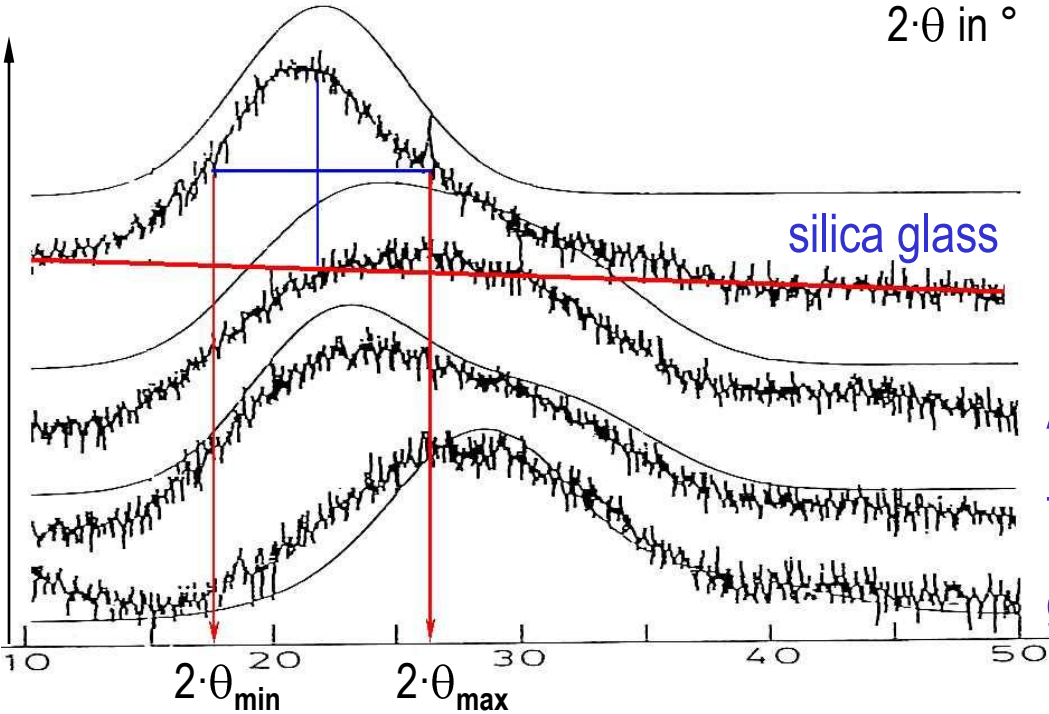
maybe Thursday

crystalite



$$d = \frac{X}{2 \cdot \sin(\frac{1}{2} \cdot 2 \cdot \theta)}$$

$X = 0.154... \text{ nm (Cu } K\alpha)$



peak broadening B

$$B = \frac{1}{2} \cdot (\theta_{\max} - 2 \cdot \theta_{\min}) \cdot \frac{\pi}{180}$$

A fibre glass

float glass

glassy basalt rock

Bragg's formula relates lattice spacing d and diffraction angle θ :

$$d = \frac{X}{2 \cdot \sin(\frac{1}{2} \cdot 2 \cdot \theta)} ; X = 0.154 \text{ nm (Cu } K\alpha)$$

broadening B of diffraction peaks:

$$B = \frac{1}{2} \cdot \left(\theta_{\max} - 2 \cdot \theta_{\min} \right) \cdot \frac{\pi}{180} ; (22 \pm 4)^\circ \Rightarrow B = 0.0698$$

interpretation: peak broadening $B \Leftrightarrow$ range b of translational symmetry

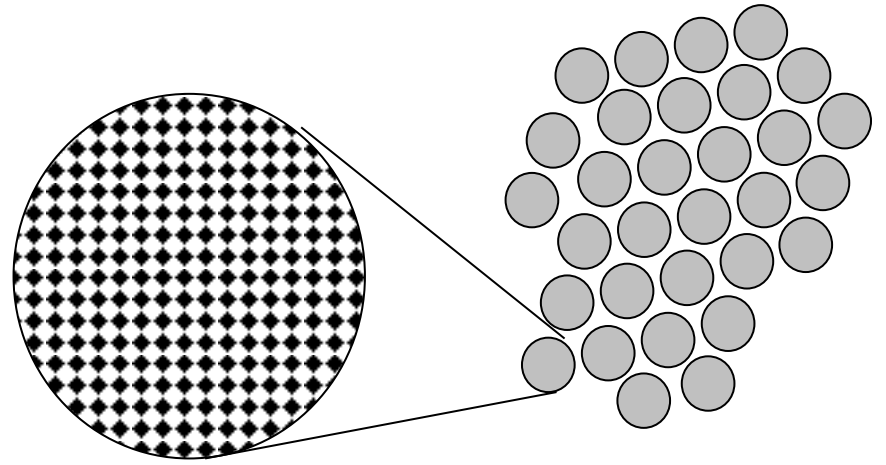
$$b = \frac{X}{B \cdot \cos(\frac{1}{2} \cdot 2 \cdot \theta)} ; (22 \pm 4)^\circ \Rightarrow b = 2 \text{ nm}$$

interpretation: peak broadening $B \Leftrightarrow$ fluctuation $\varepsilon = \delta d/d$ of lattice spacing

$$\varepsilon = \frac{\delta d}{d} = \frac{B}{4 \cdot \tan(\frac{1}{2} \cdot 2 \cdot \theta)} ; (22 \pm 4)^\circ \Rightarrow \varepsilon = 9 \%$$

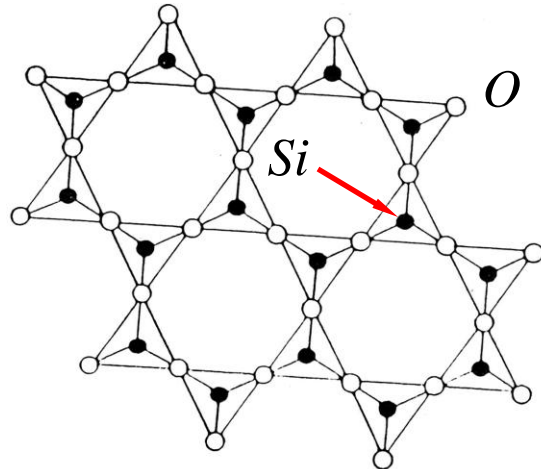
1.

*random packing
of length-correlated
clusters*

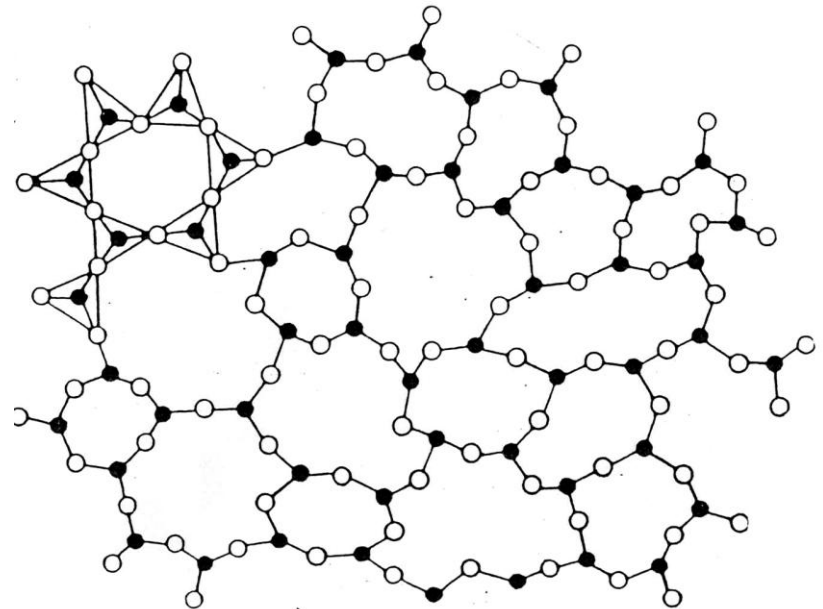


2.

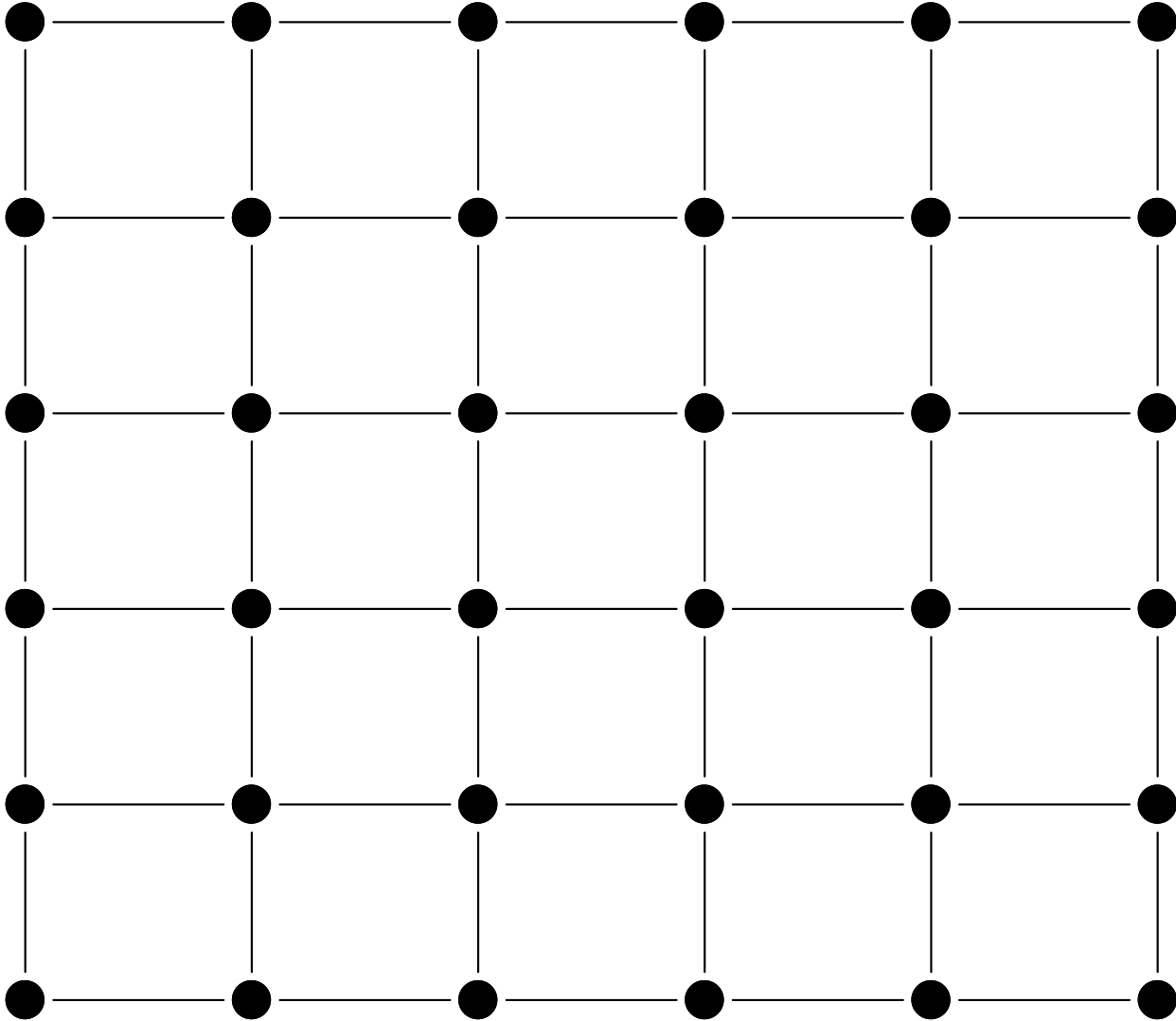
crystalline silica



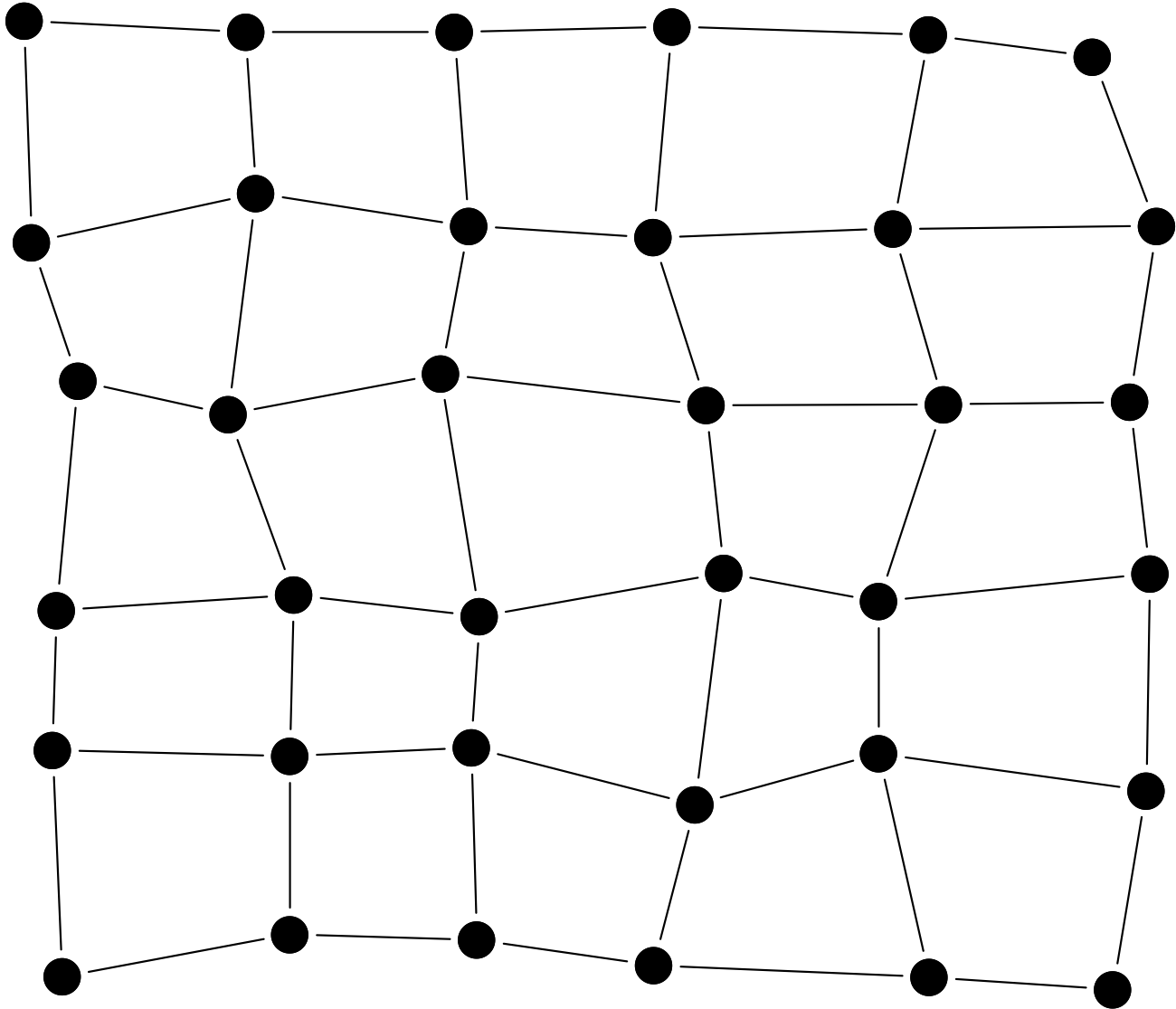
silica glass



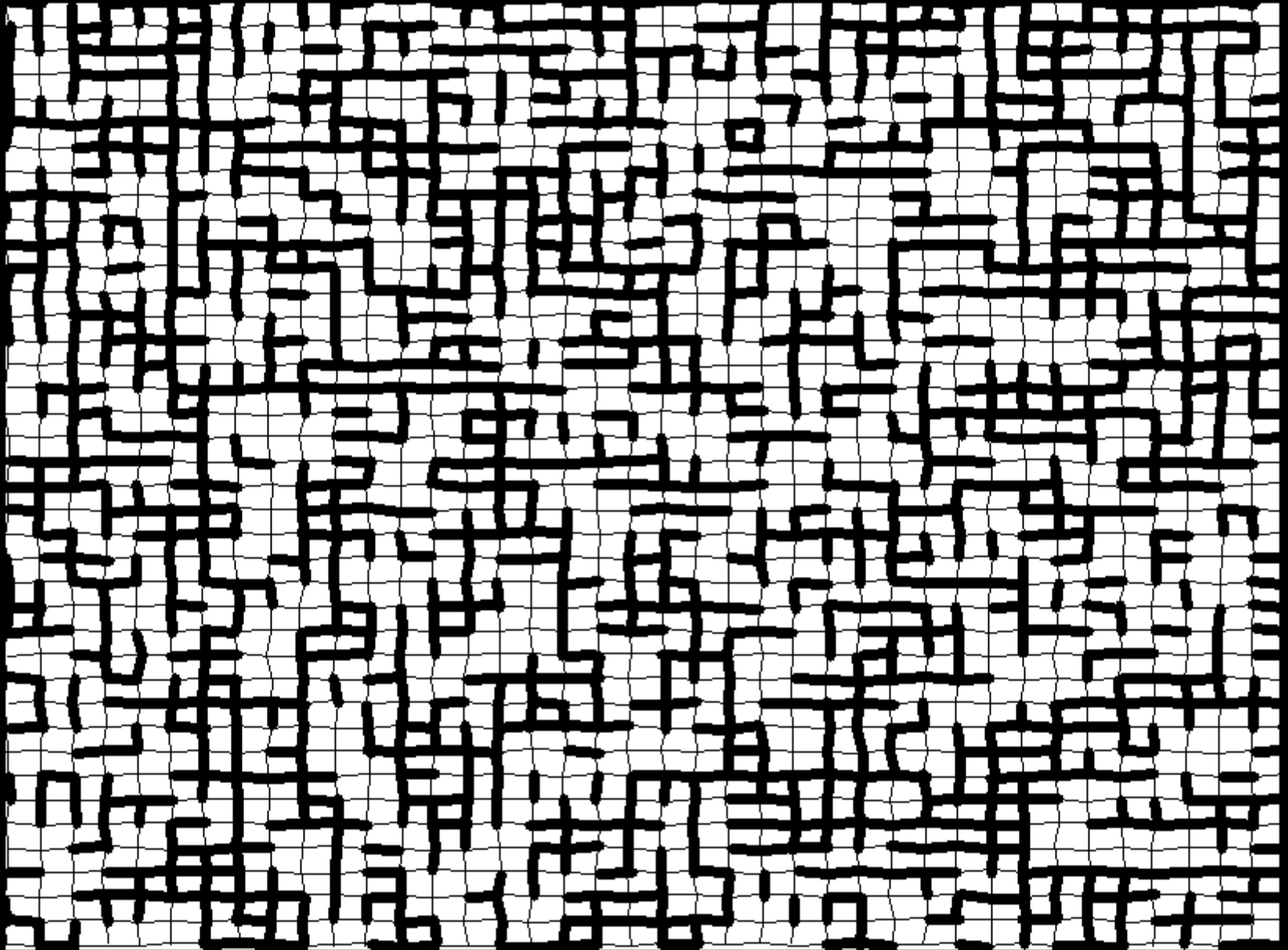
$$\varepsilon = \pm 0.00$$



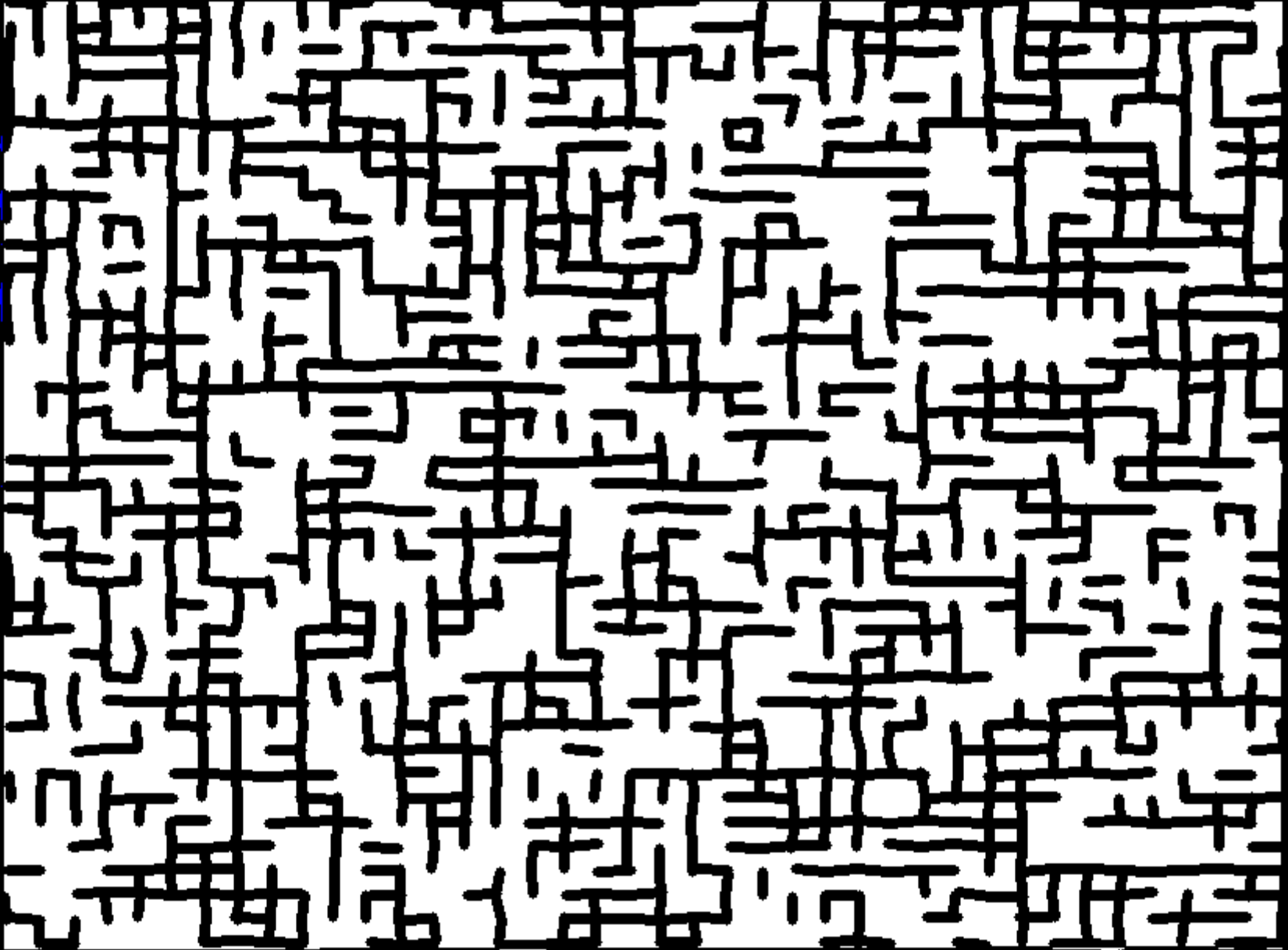
$$\varepsilon = \pm 0.09$$



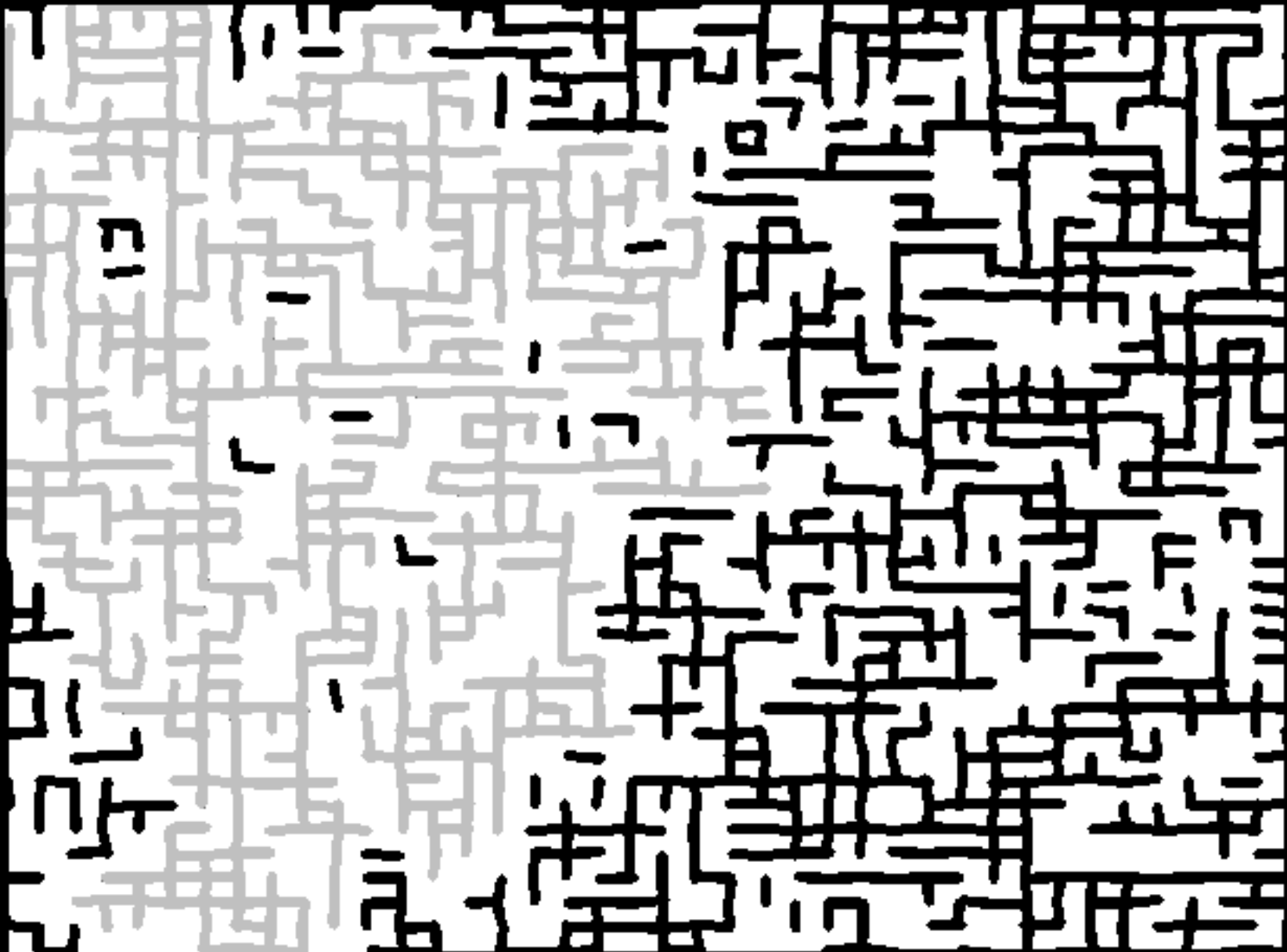
$\varepsilon = 0.09$; $d \leq \varepsilon \cdot \sqrt{2}$ is accepted as bond



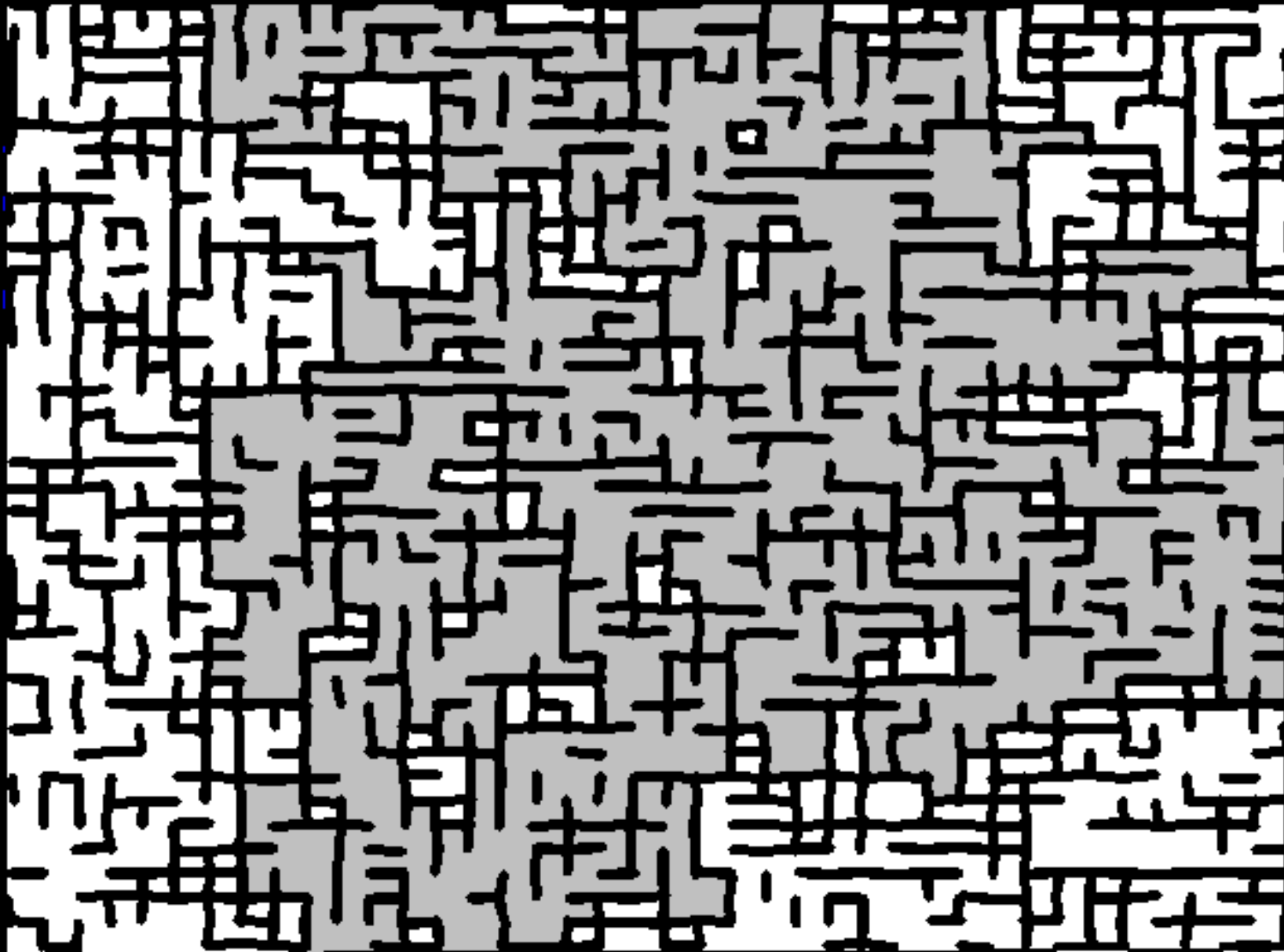
$\varepsilon = 0.09$; $d \leq \varepsilon \cdot \sqrt{2}$ is accepted as bond

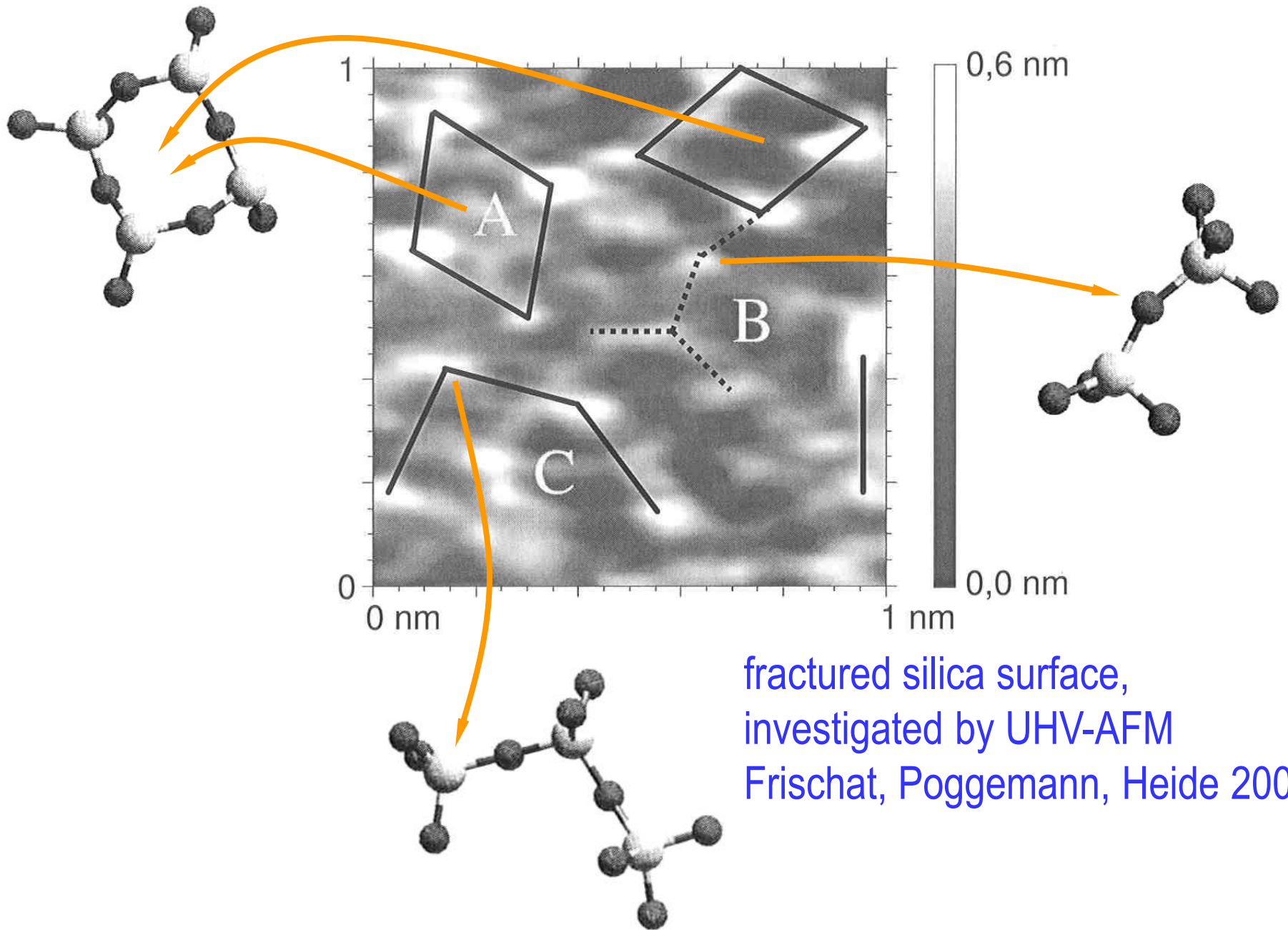


rigidity percolation; see —



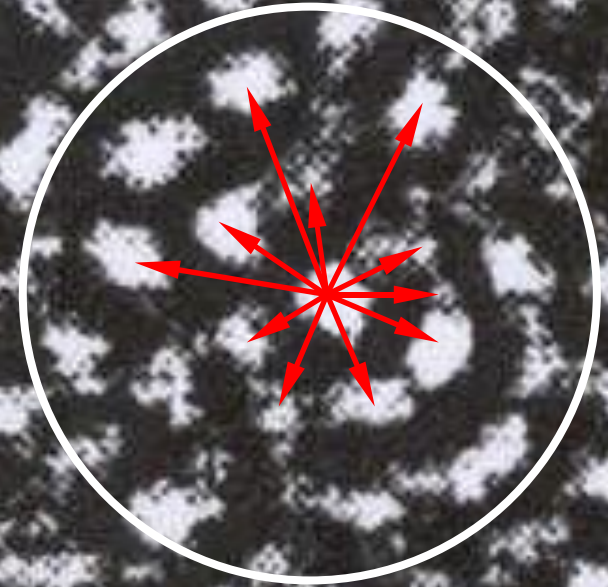
diffusion path percolation; see 



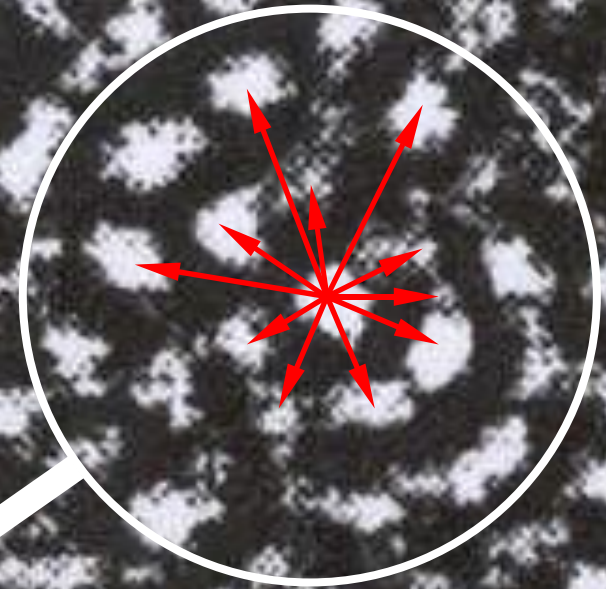
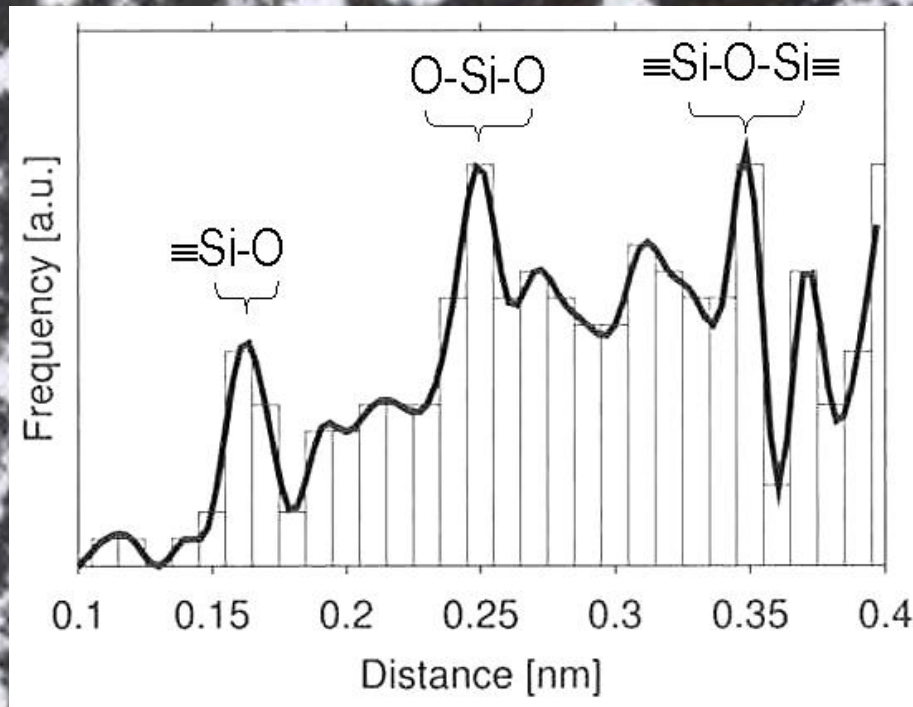


fractured silica surface,
investigated by UHV-AFM
Frischat, Poggemann, Heide 2003

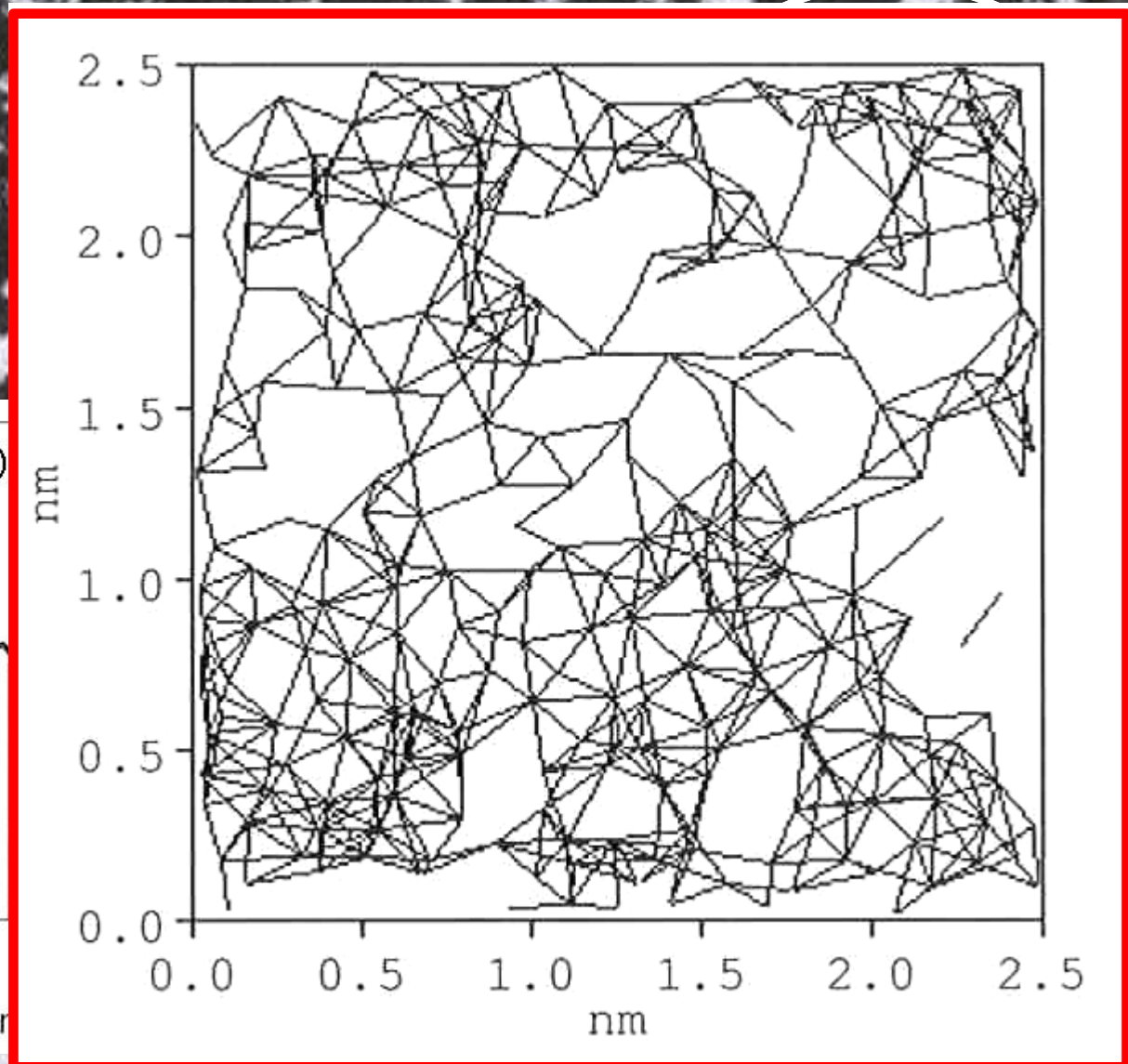
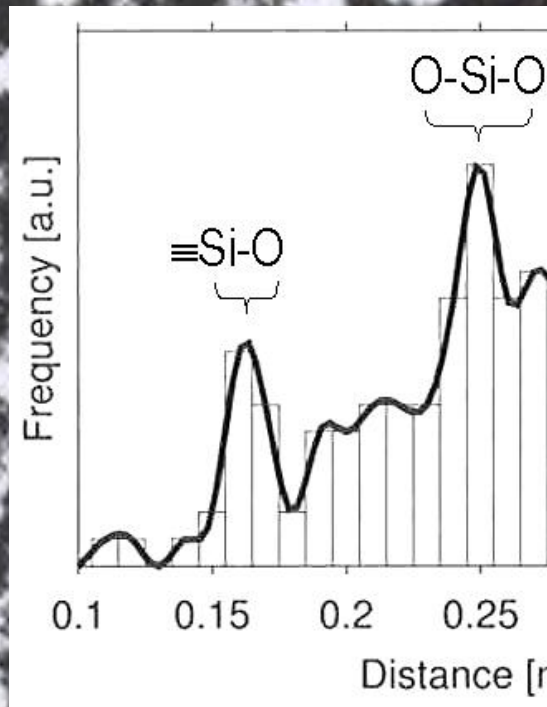
evaluate all distance correlations in the window 0.2 - 0.3 nm (O-Si-O distances)



evaluate all distance correlations in the window 0.2 - 0.3 nm (O-Si-O distances)



evaluate all distance correlations in the window 0.2 - 0.3 nm (O-Si-O distances)



QUIZ

1.

The c_P of a condensed phase is determined by

- the vibrational and rotational degrees of freedom of its molecular units.
- the vibrational, rotational, and configurational degrees of freedom of its molecular units.
- electrons may contribute to c_P .
- More than one answer is correct.

2.

In the glass transition, we observe

- a jump of the heat capacity and the thermal expansion coefficient to higher values.
- a jump of the compression modulus to lower values.
- Both answers are correct.

3.

The c_P of a glass at very low T

- shows an excess over the c_P of the isochemical crystal.
- steadily approaches the value of the isochemical crystal.
- may be larger or smaller than the c_P of the isochemical crystal, depending on the chemical composition.

4.

The properties of frozen-in phases

- are functions of state.
- are unpredictable.
- depend well-defined additional parameters.

QUIZ continued

5.

The c_p jump in the glass transition

- reflects the gain of additional degrees of freedoms.
- influences the fragility of a glass forming system.
- None of these answers is correct.

6.

Multi-component oxide melts above T^m

- are in perfect disorder.
- still contain quasi-crystalline entities (micro-crystals).
- reflect the constitutional relations found in the crystalline state.

7.

The entropy of condensed phases for $T \rightarrow 0$

- is always zero.
- is zero if the phases are in internal equilibrium.
- can be assessed unequivocally by calorimetry.
- The issue is a yet unresolved controversy of condensed matter physics.

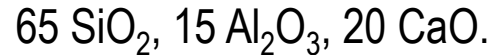
8.

Industrial glasses at 298 K obtained at normalized cooling rates

- cannot be described by thermochemical standard values.
- can be described by the thermochemical standard values of the corresponding crystalline state plus two additional quantities.
- can be described by the thermochemical standard values only if they are one-component systems.

Homework

Derive the enthalpy and entropy ($T = 298 \text{ K}$) of a glass of composition (in wt. %)



Compare the result to the values of the isochemical (polycrystalline) equilibrium state.

What is the Gibbs energy difference at 298 K?

With a value of $T_g = 760 \text{ }^\circ\text{C}$ and $H^{\text{fus}} = 475 \text{ kJ/kg}$ given, what estimate do you derive for Δc_p ?

What is the expected fragility slope?

**Thank you
for your attention**